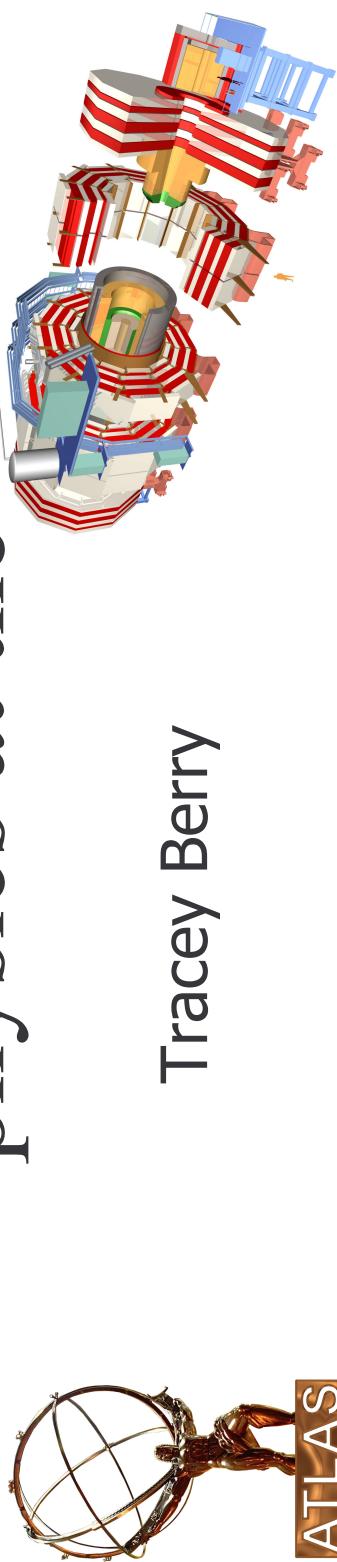


Prospects for the observation of new physics at the LHC



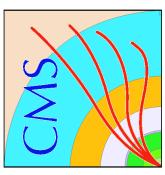
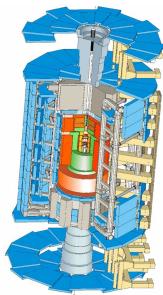
Tracey Berry

ATLAS

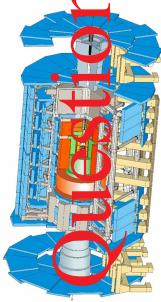
University of Royal Holloway,
London
Flavour in the era of the LHC,
CERN

27th March 2007

Overview



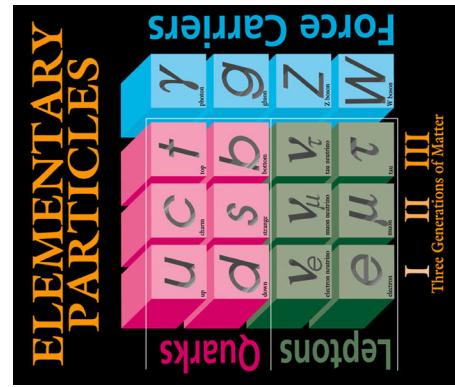
- Beyond the Standard Model
- Fundamental Symmetries
- Heavy Gauge Bosons W' and Z'
- Electroweak Symmetry Breaking
- SUSY, Little Higgs, Technicolor
- Leptons & Quarks, Other New Particles
- Left-right symmetry, E6 quarks
- Extra Dimensions
- Summary and Outlook

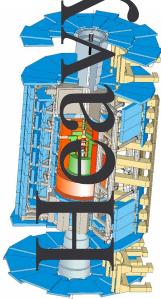


Questions remaining with the Standard Model

→ Reasons to search for new physics...

- Fundamental symmetries:
 - Are there more symmetries beyond $SU(3)_C \otimes SU(2)_L \otimes U(1)_Y$?
→ GUTs with larger symmetry group? Left-right symmetry?
 - ElectroWeak Symmetry Breaking (EWSB):
 - Unitarity violation in longitudinal WW scattering at high E
solution: Higgs boson or other new particle with mass < 1 TeV
 - If Higgs → *hierarchy problem*: fine tuning in rad corr to Higgs mass
solution: new physics at TeV scale (SUSY, Little Higgs, etc...)
 - If NO Higgs
solution: new strong interactions (Technicolor, etc...)
- Quark and lepton generations:
 - Why are there 3 generations? → Fermions composite?
 - Is there a lepto(n)-quark symmetry?
 - More than 3 generations of quarks & leptons?





Heavy Gauge Bosons

- Many extensions of the SM rely on larger symmetry groups (GUTs, string-inspired, left-right, little Higgs models, etc...) ➔ predict existence of new gauge bosons W' and Z' (or KK modes)

- Production: s-channel

- Clean decay channels:

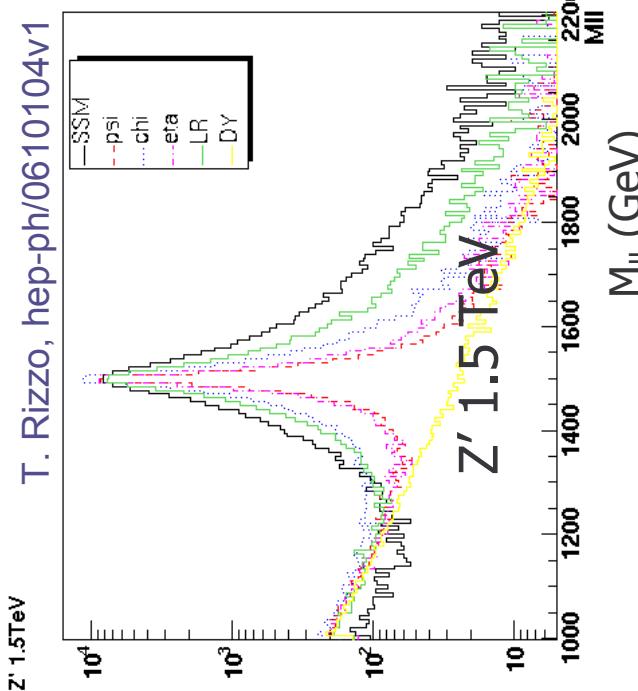
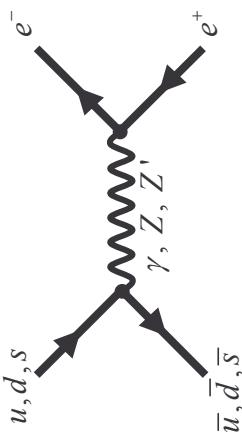
$$\begin{aligned} W' &\rightarrow e^+ \nu_e \text{ or } \mu^+ \nu_\mu \\ Z' &\rightarrow e^+ e^- \text{ or } \mu^+ \mu^- \end{aligned}$$

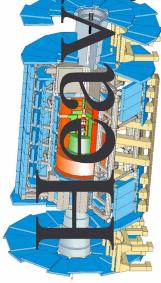
- Tevatron searches: M up to ~ 1 TeV

- Z' models considered:

- Sequential SM (SSM) with same Z' couplings to fermions as for Z

- Models based on different patterns of E6 symmetry breaking (ψ , χ and η)
- Left-right (LR) symmetry models

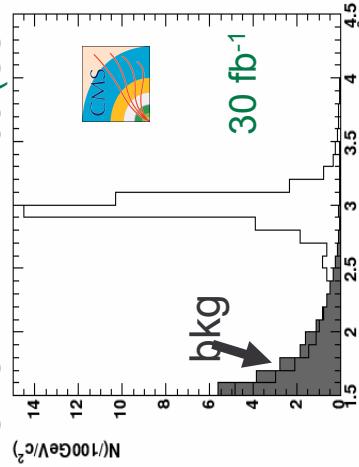




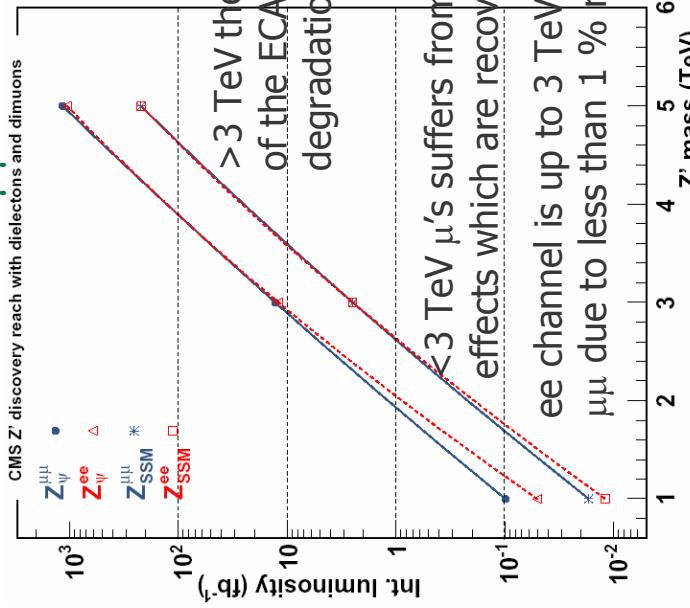
Heavy New Gauge Bosons Z'

- Selection: pairs of isolated e or μ
- Bkg: dominated by dileptons from Drell-Yan
- 5σ discovery up to ~ 5 TeV (model dependent) for both ATLAS and CMS

CMS PTDR $Z' \rightarrow ee$ (SSM)



Z' 5 σ reach ee vs. $\mu\mu$ channels



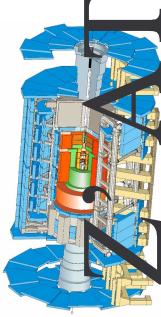
CMS PTDR 2006
 $Z' \rightarrow ee$



A very low luminosity, < 0.1 fb^{-1} , should be sufficient to discover Z' bosons at 1 TeV for all models and in a single channel (e/m)

Flavour Workshop, CERN

Tracey Berry (RHUL)

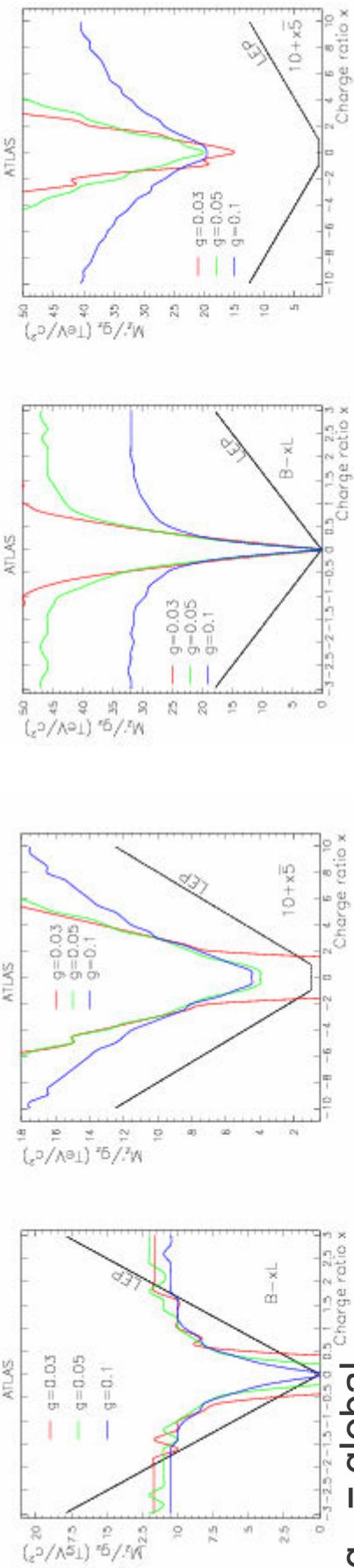


ATLAS reaches

ATLAS discovery reaches in CDDT models of Z'

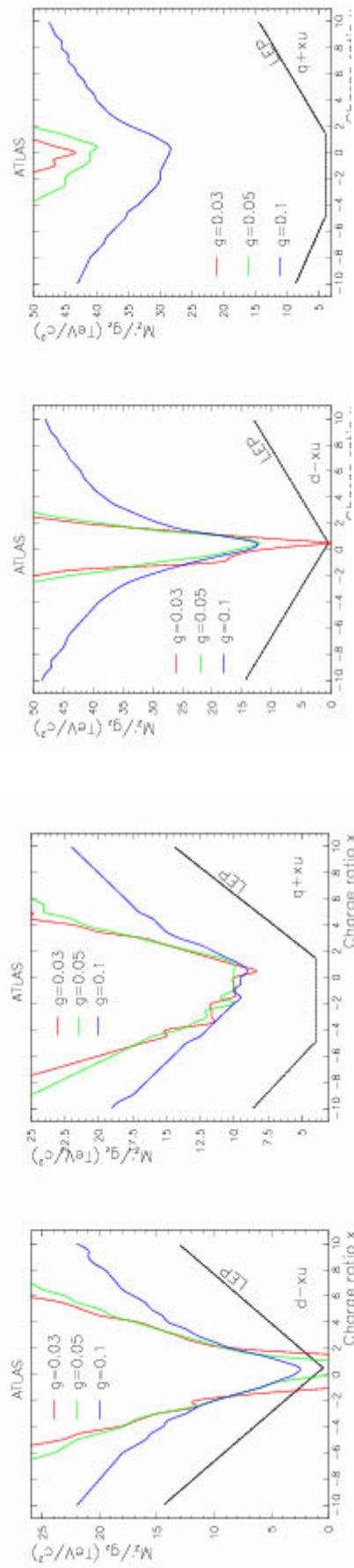
$L=400\text{pb}^{-1}$: first months of running at LHC

$L=100\text{fb}^{-1}$



g_Z = global
coupling strength

x = describes relative coupling strength to different fermions

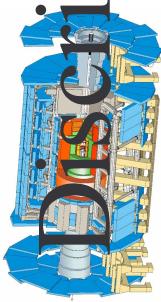


With only 400 pb^{-1} : can go beyond LEP exclusion limits and probe regions of parameter space not yet excluded by CDF.

Flavour Workshop, CERN

Tracey Berry (RHUL)

ATL-PHYS-PUB-2006-024



Discrimination between Z' models

- Models differ in the Z' couplings to fermions
- especially parity-violating couplings to leptons + couplings to initial u/d
- Distinguish by measuring:
 - $\sigma \times \Gamma$: Decay width (ee only due to worse $\mu\mu$ resolution)

ATLAS: $\sigma(p_T)/p_T \approx 0.7\% \text{ (e), } 10\% \text{ (\mu) at } p_T = 1 \text{ TeV}$

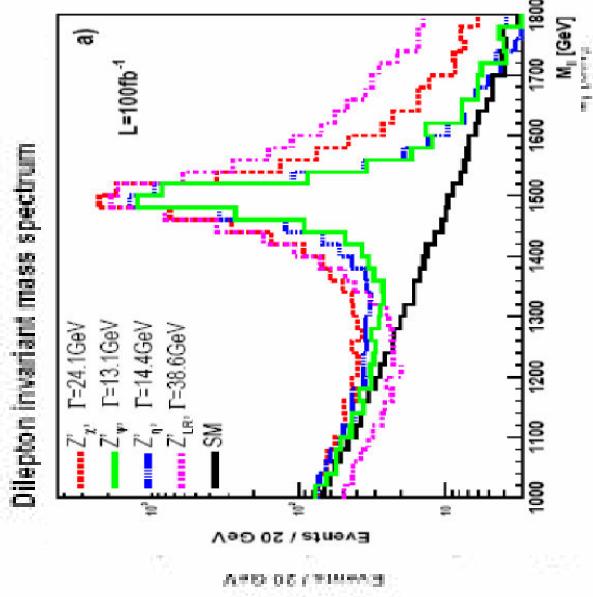
- Forward-backward asymmetry

- Z' rapidity

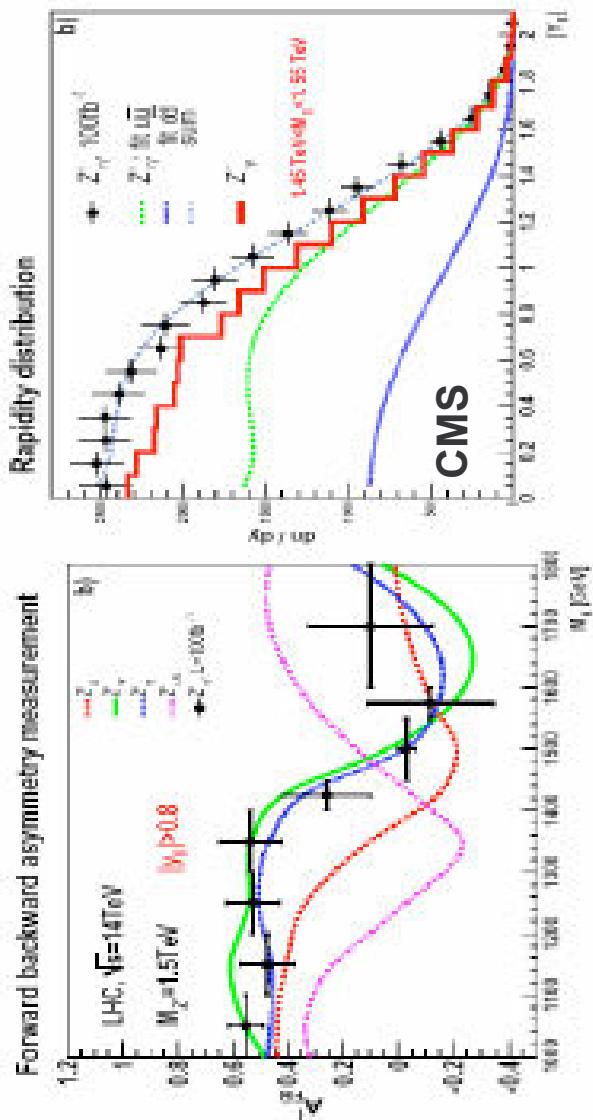
→ Also provides discrimination against other models like extra-D, little Higgs, ...

In one year at 10^{34} : reach 4-5 TeV & discriminating between models possible up to $m \sim 2.5$ TeV

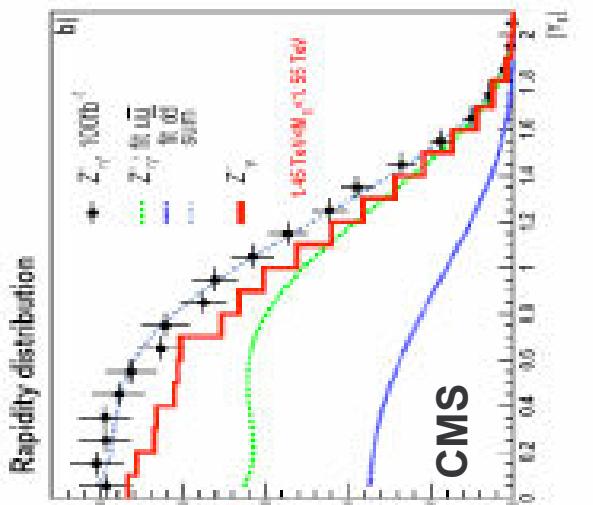
Model	$\sigma_{ } \times \Gamma_{ }$ (fb \times GeV)	A_{FB} Corrected at Z' peak	$M = 1.5 \text{ TeV}$
$Z' \rightarrow ee$ 100 fb $^{-1}$	36668 ± 138	$+0.108 \pm 0.027$	
SSM	828 ± 48	-0.361 ± 0.030	
χ_{LR}	1515 ± 75	$+0.186 \pm 0.032$	



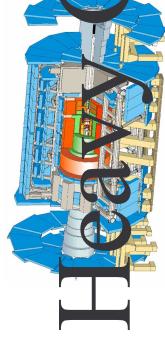
Dilepton invariant mass spectrum



Forward backward asymmetry measurement

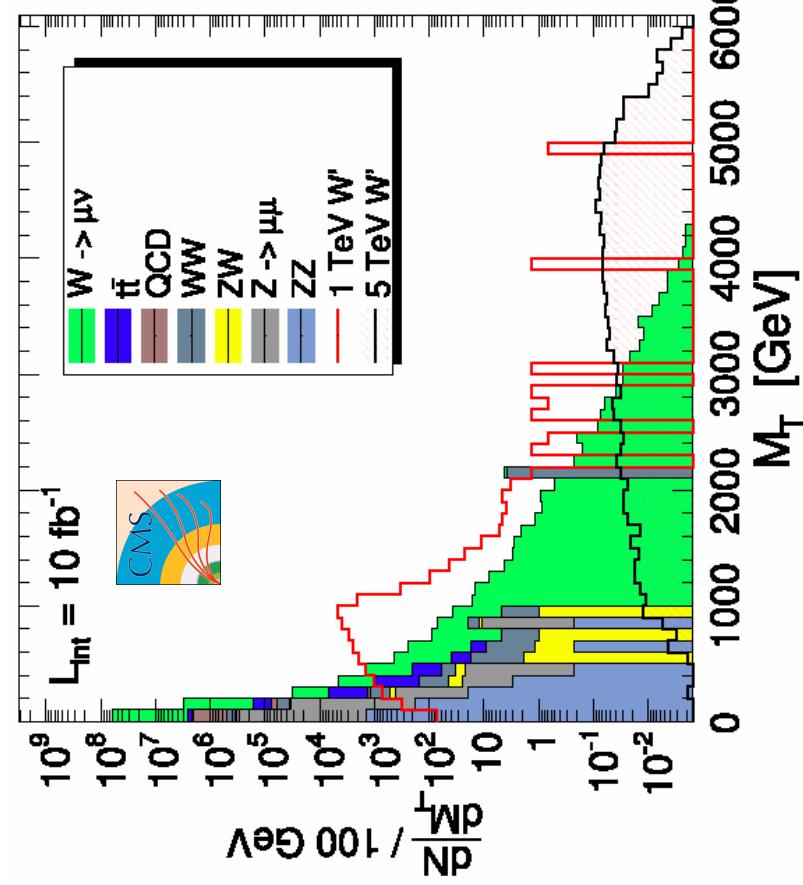


Rapidity distribution

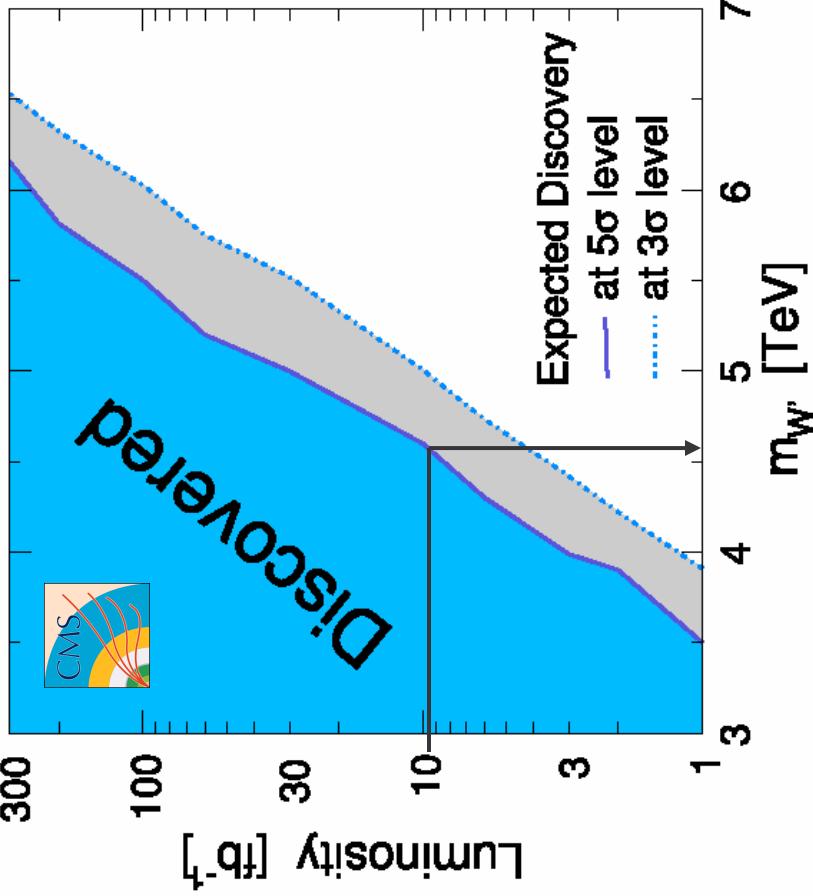


Heavy Gauge Bosons: W'

- General Model by Altarelli, Mele, Ruiz-Altaba with same W' couplings to fermions as for W
- Selection: one-muon event with track isolation reqt around μ
 - + missing transverse energy
- Background: mostly $W \rightarrow \mu \nu$

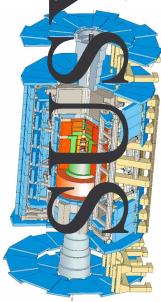


Flavour Workshop, CERN



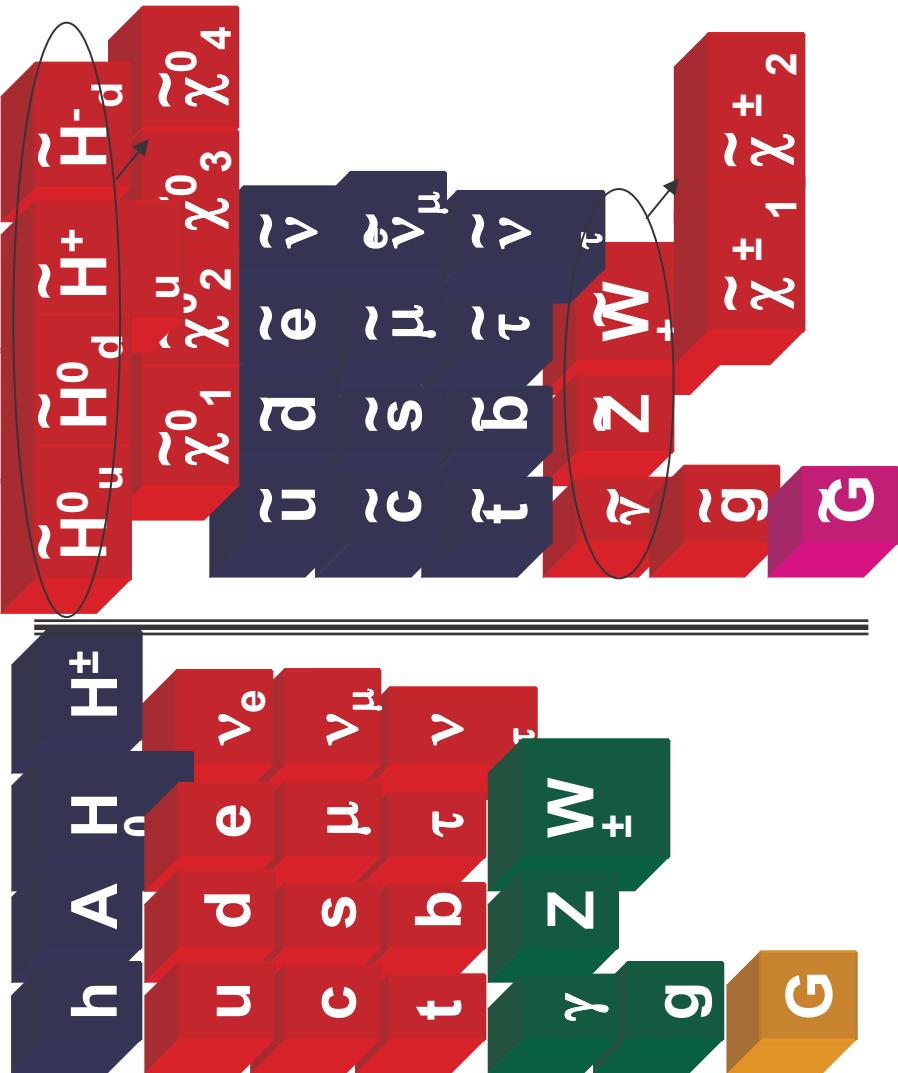
Tracey Berry (RHUL)

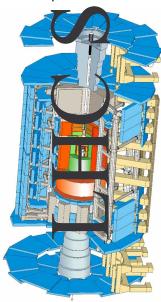
8



SUSY Spectrum

- SUSY gives rise to partners of SM states with opposite spin-statistics but otherwise same Quantum Numbers.
spin-1/2 matter particles (fermions) \leftrightarrow spin-1 force carriers (bosons)
- Expect SUSY partners to have same masses as SM states
 - Not observed
 - SUSY must be a broken symmetry
- Different mechanisms of SUSY breaking lead to different models MSSM, mSugra, MSSM, GMSB, AMSB
- R-Parity $R_p = (-1)^{3B+2S+L}$
 - Conservation of R_p causes LSP to be stable

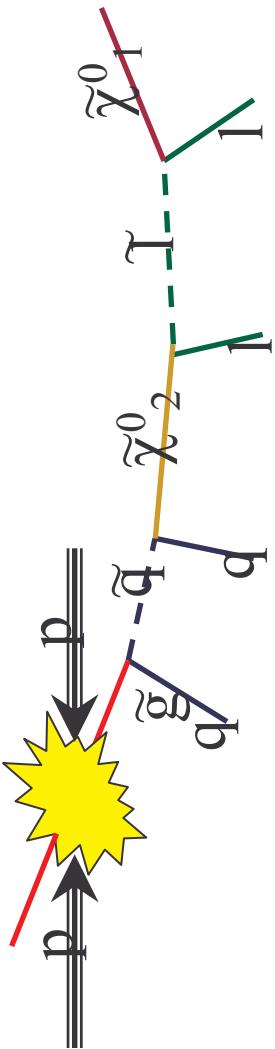




SUSY Searches



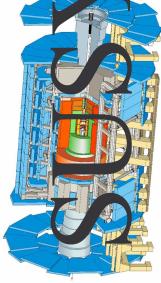
Typical Signature/decay chain:



- Strongly interacting sparticles (squarks, gluinos) dominate production.
- Heavier than sleptons, gauginos etc. g cascade decays to LSP.
- Potentially long decay chains and large mass differences
 - **Many high p_T objects observed (leptons, jets, b-jets).**
- If R-Parity conserved LSP (lightest neutralino in mSUGRA) stable and sparticles pair produced.
 - **Large E_T^{miss} signature (c.f. $Wg|\nu$).**

SUSY Searches

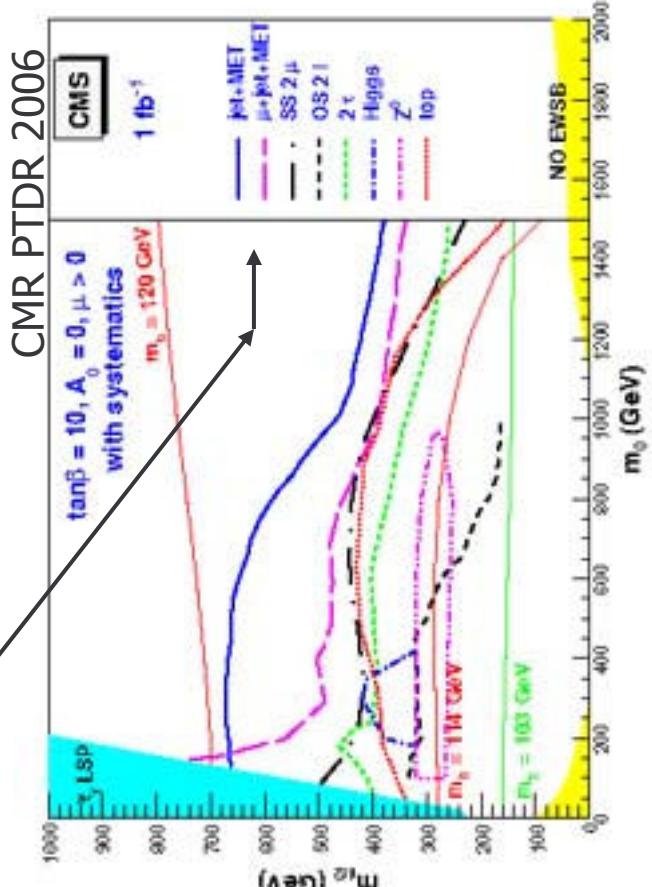
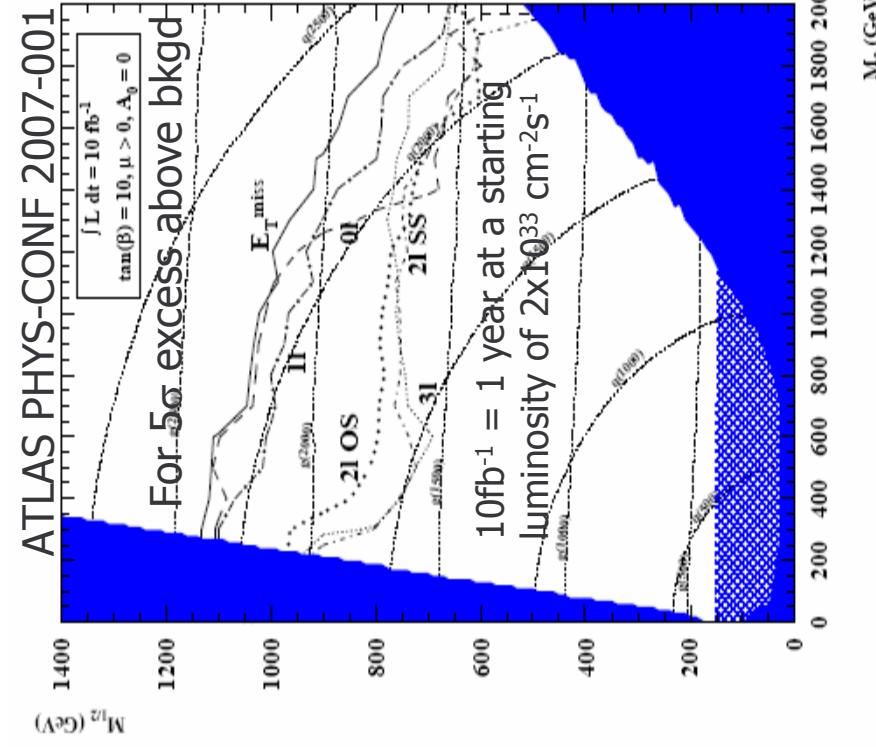
- **Inclusive** searches to detect SUSY with first data
- **Exclusive** studies – performed with more data to determine model parameters e.g. masses etc from end point measurements...



SUSY discovery with first data

Inclusive searches to detect SUSY with first data

Look for deviations from SM predictions – requires good knowledge of SM processes at that energy scale and on an understanding of the detectors' performances
Best reach obtained with the most inclusive channels:

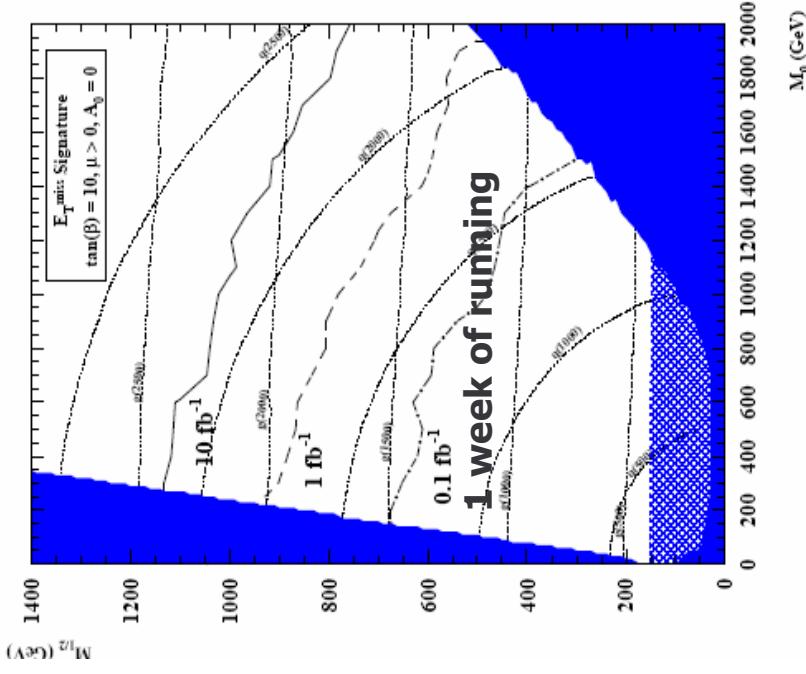


The range of gluino and squark masses up to about 1.5 TeV can be probed with $L=1\text{fb}^{-1}$ and up to about 2 TeV with 10fb^{-1}



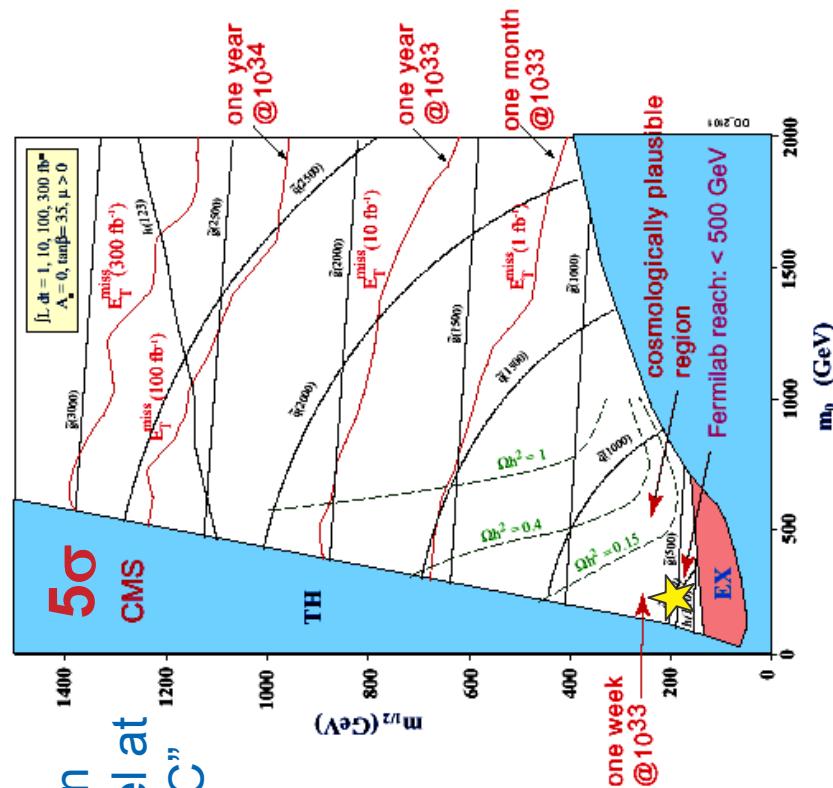
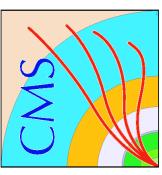
LHC Jets + E_T^{miss}

- Inclusive searches with Jets + n leptons + E_T^{miss} channel.
- Map statistical discovery reach in mSUGRA m_0 - $m_{1/2}$ parameter space.
- Sensitivity only weakly dependent on A_0 , $\tan(\beta)$ and sign(μ).

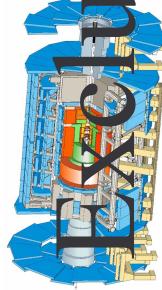


ATLAS PHYS-CONF 2007-001

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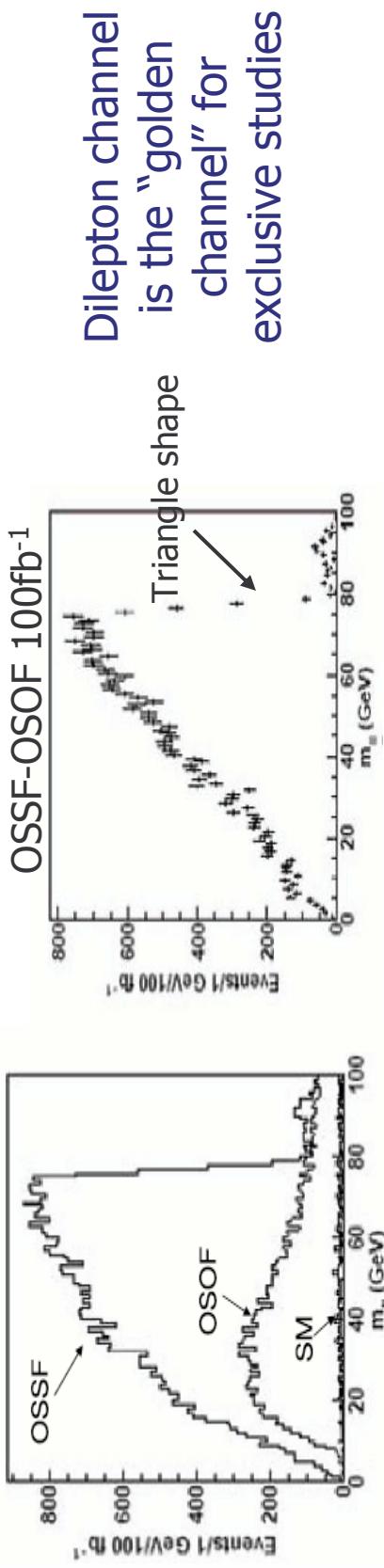
Tracey Berry (RHUL)



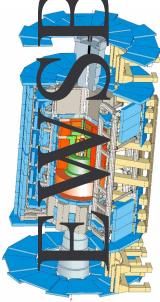
Inclusive SUSY Searches

- **Exclusive** studies – reconstruct specific decay channels in order to estimate physical parameters characterising the decay:
e.g. reconstructing kinematic endpoints (edges & thresholds) in invariant mass distributions to extract masses

Simultaneous observation of a signal in various topologies would allow measurements of cross-sections of sub-processes and their ratios and help to determine the underlying physics.



Cleanest signature is the opposite sign same flavour invariant mass distributions of lepton (e/μ) pairs. **From its endpoints can extract precise constraints on difference in masses between s-particles in decay chain with a statistical precision of 0.1% or better with L=100fb⁻¹**

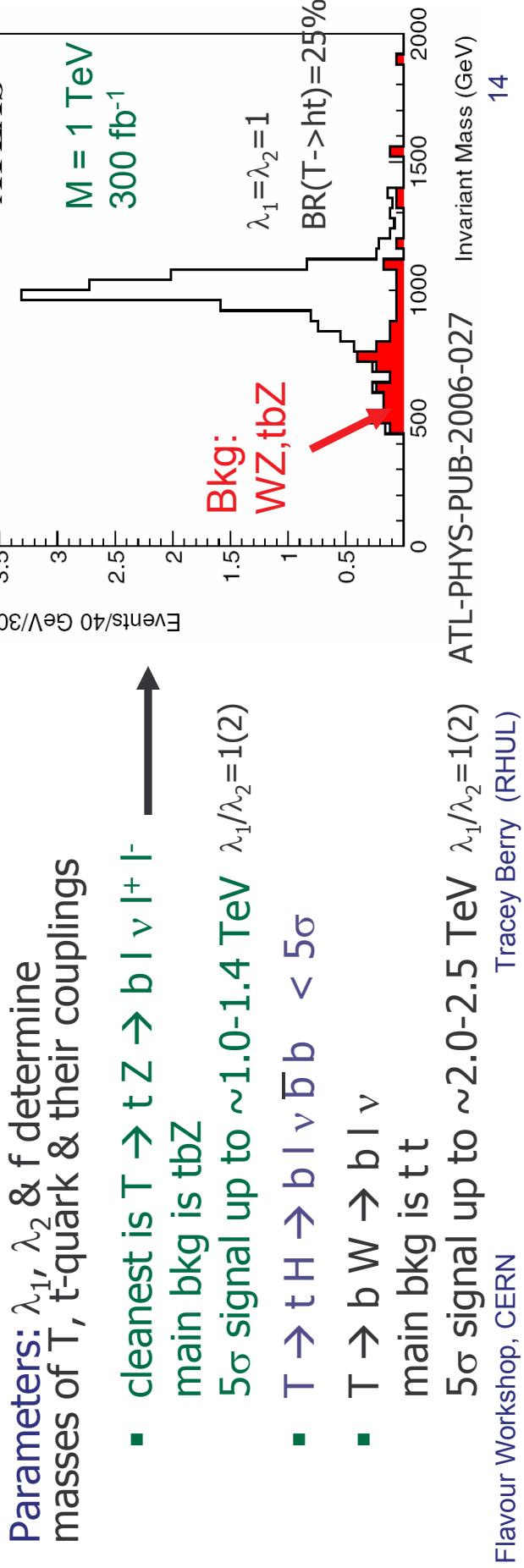


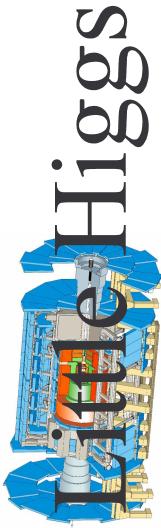
FYS-B: Little Higgs

- Models with Higgs as pseudo-Goldstone boson from a broken global symmetry ($SU(5)$ in "littlest Higgs model")
 - Extra $Q=2/3$ heavy quark (T) and heavy gauge bosons (A_H, W_H, Z_H)
 - Quadratic divergences cancel top and VB divergences to Higgs mass

Heavy Quark T

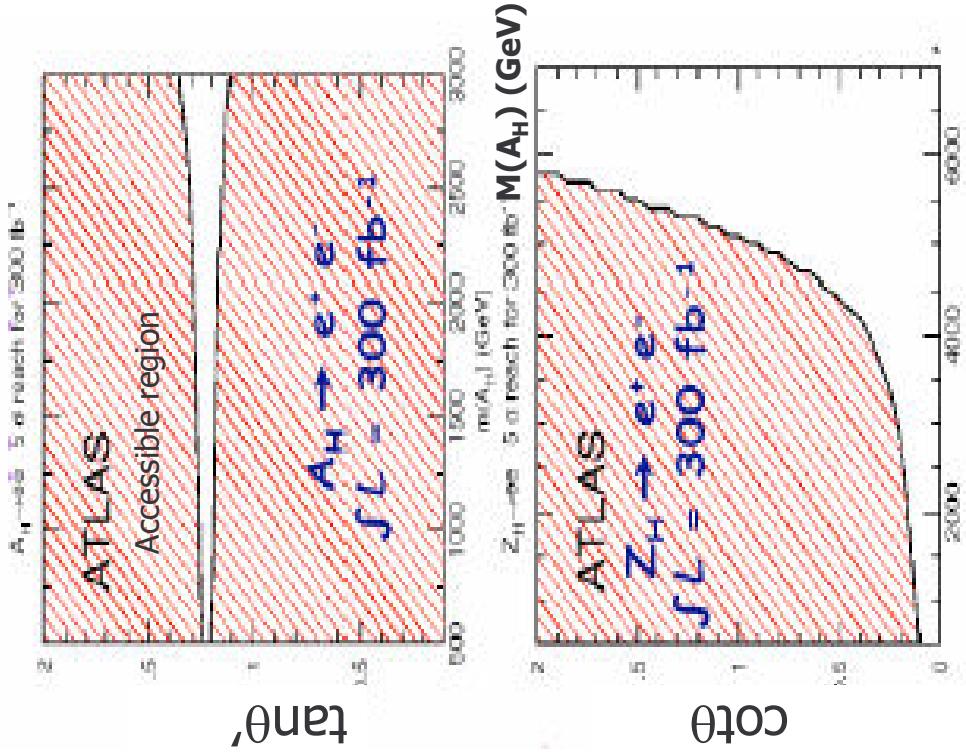
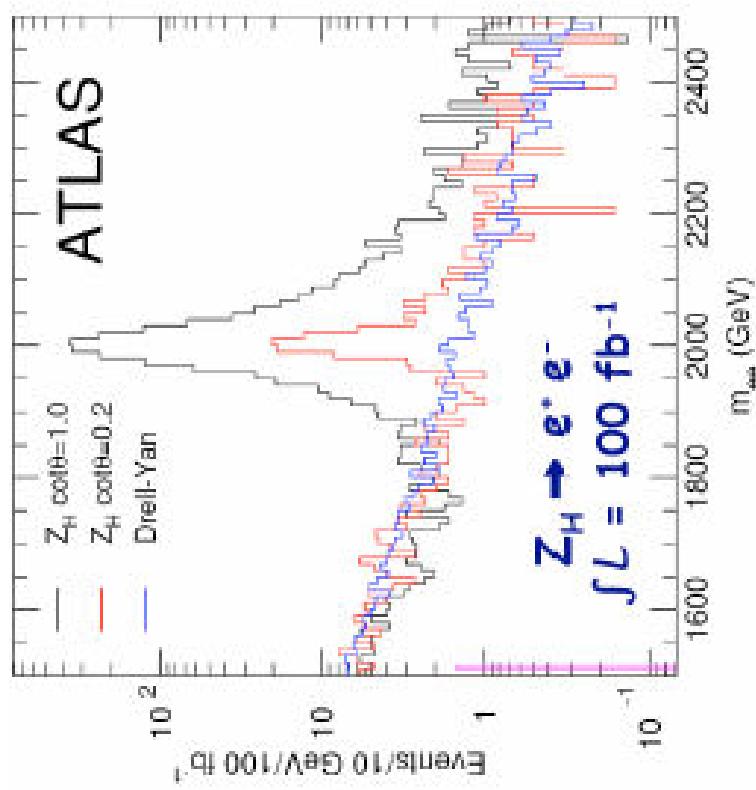
- Production:
 - via QCD ($gg \rightarrow T\bar{T}$, $q\bar{q} \rightarrow T\bar{T}$)
 - via W exchange ($qb \rightarrow q' T$) dominant for $M_T > 700$ GeV
- Decays: $T \rightarrow tZ, T \rightarrow tH, T \rightarrow bW$





- A_H and Z_H discovery in lepton modes

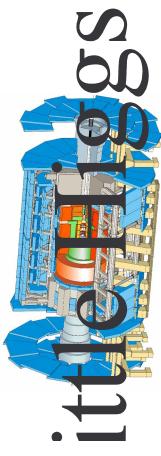
Signal: dilepton resonance



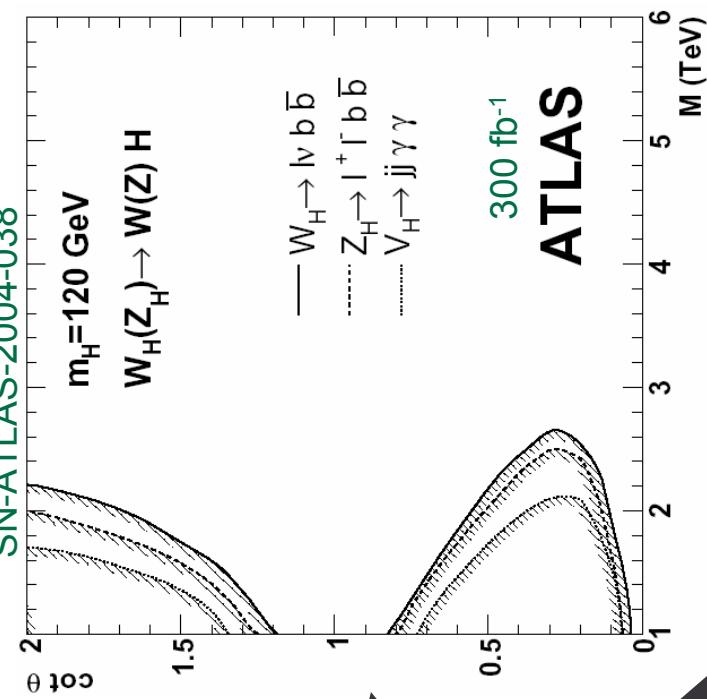
- Reach up to $\mathcal{M} \sim 6$ TeV (depending on mixing angle $\cot \theta$)
- New gauge bosons observable over almost the entire mass range

Flavour Workshop, CERN

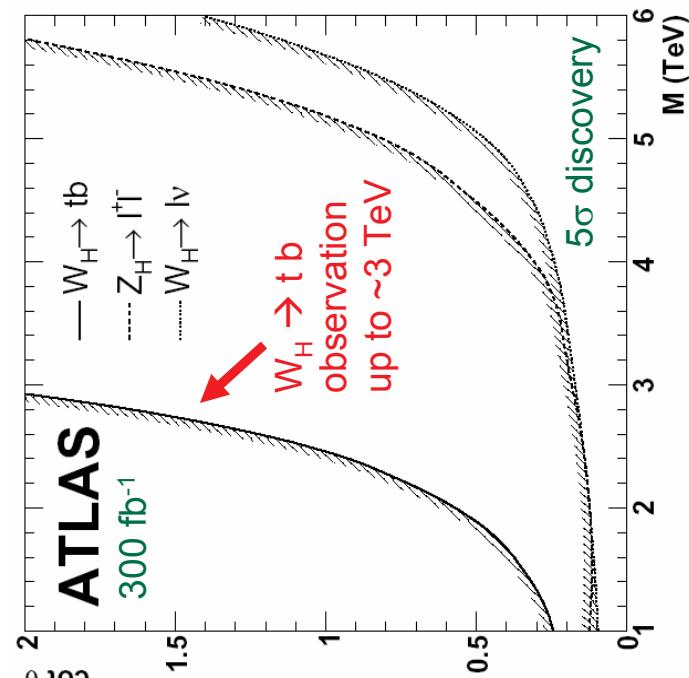
Tracey Berry (RHUL)



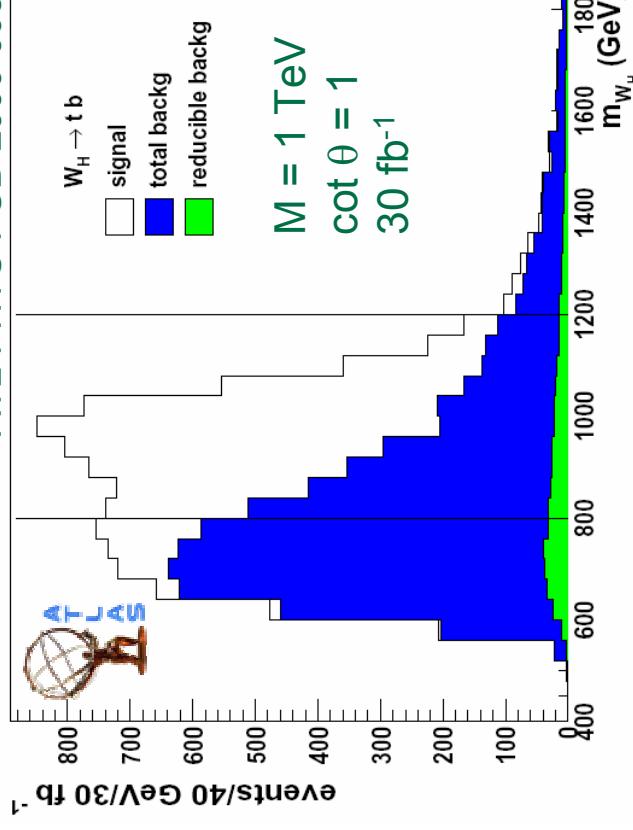
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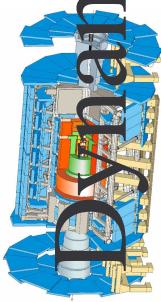


- Discrimination against other models predicting dilepton resonances via observation of decay modes like $W_H \rightarrow W H$, $Z_H \rightarrow Z H$, and $W_H \rightarrow t \bar{b}$ (important at $\cot \theta \approx 1$, where $\text{BR}(W_H/Z_H \rightarrow W/Z H)$ vanishes)



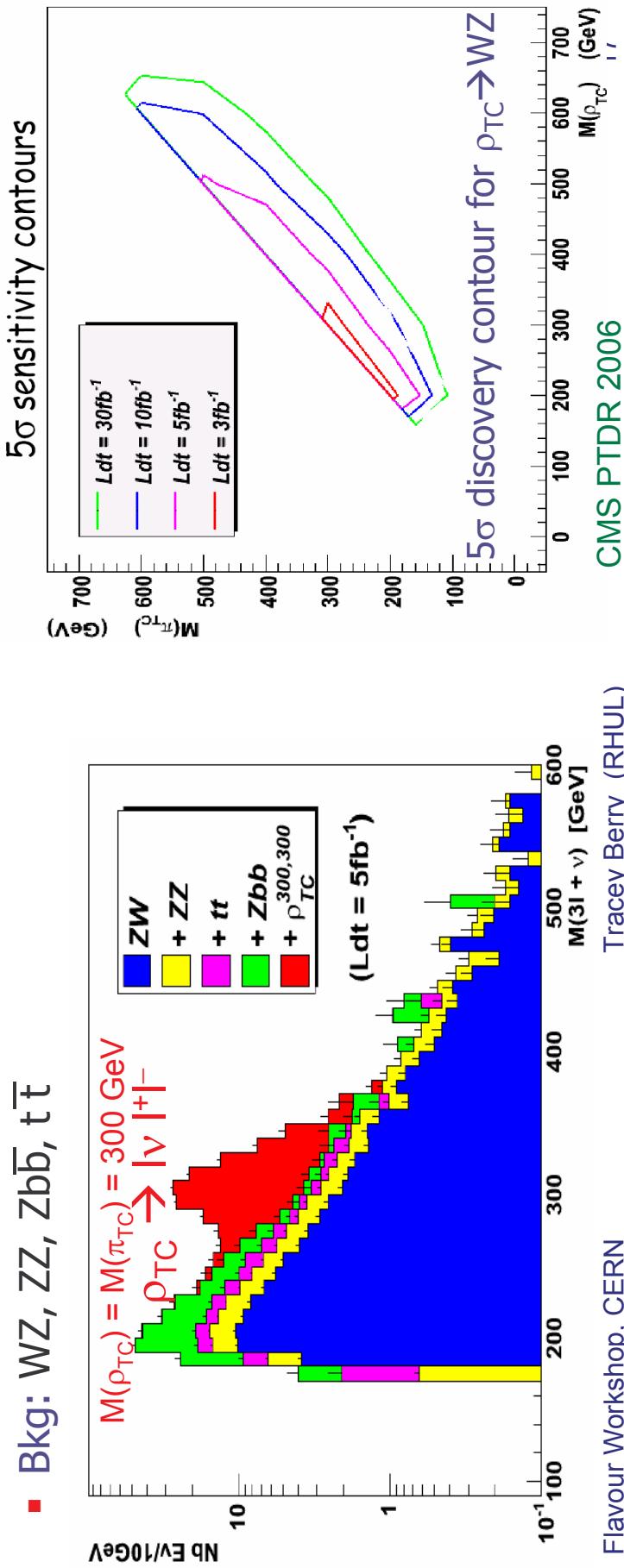
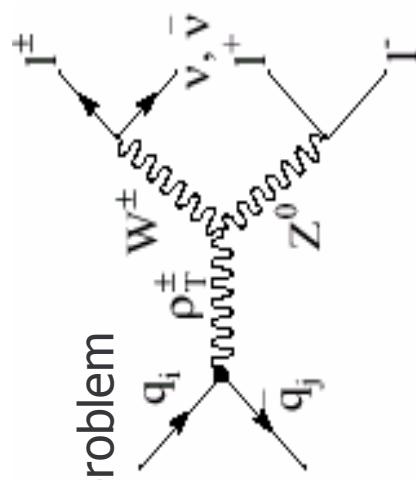
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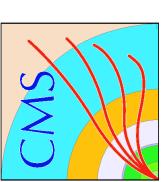




Dynamical EWSB: Technicolor

- Dynamical EWSB via new strong interaction
 - No need for Higgs boson → removes fine tuning problem
 - Predict new technifermions, technihadrons
 - Study $\rho_{TC} \rightarrow W Z$ process
 - Select isolated leptons, measure missing E_T & apply W and Z kinematical constraints
- Bkg: $WZ, ZZ, Zb\bar{b}, t\bar{t}$



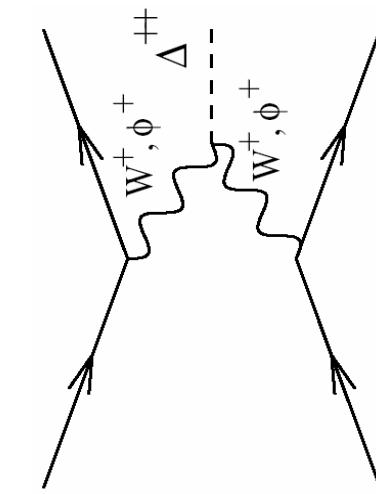
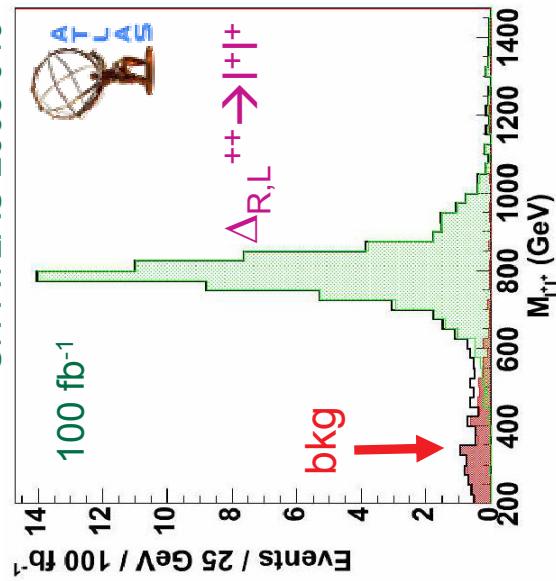


Doubly Charged Higgs in the L-R Symmetric Model

- Left-Right Symmetric Model based on $SU(2)_L \otimes SU(2)_R \otimes U(1)_{B-L}$
- Features triplet of Higgs fields ($\Delta_R^0, \Delta_R^+, \Delta_R^{++}$) + two doublets ϕ
- Predicts new gauge bosons (W_R and Z_R) & new fermions (heavy Majorana ν, ν_R)
- Addresses origin of pure left-handed charged weak interaction
+ origin of light neutrino masses (via see-saw mech. & heavy ν_R)
- Production:

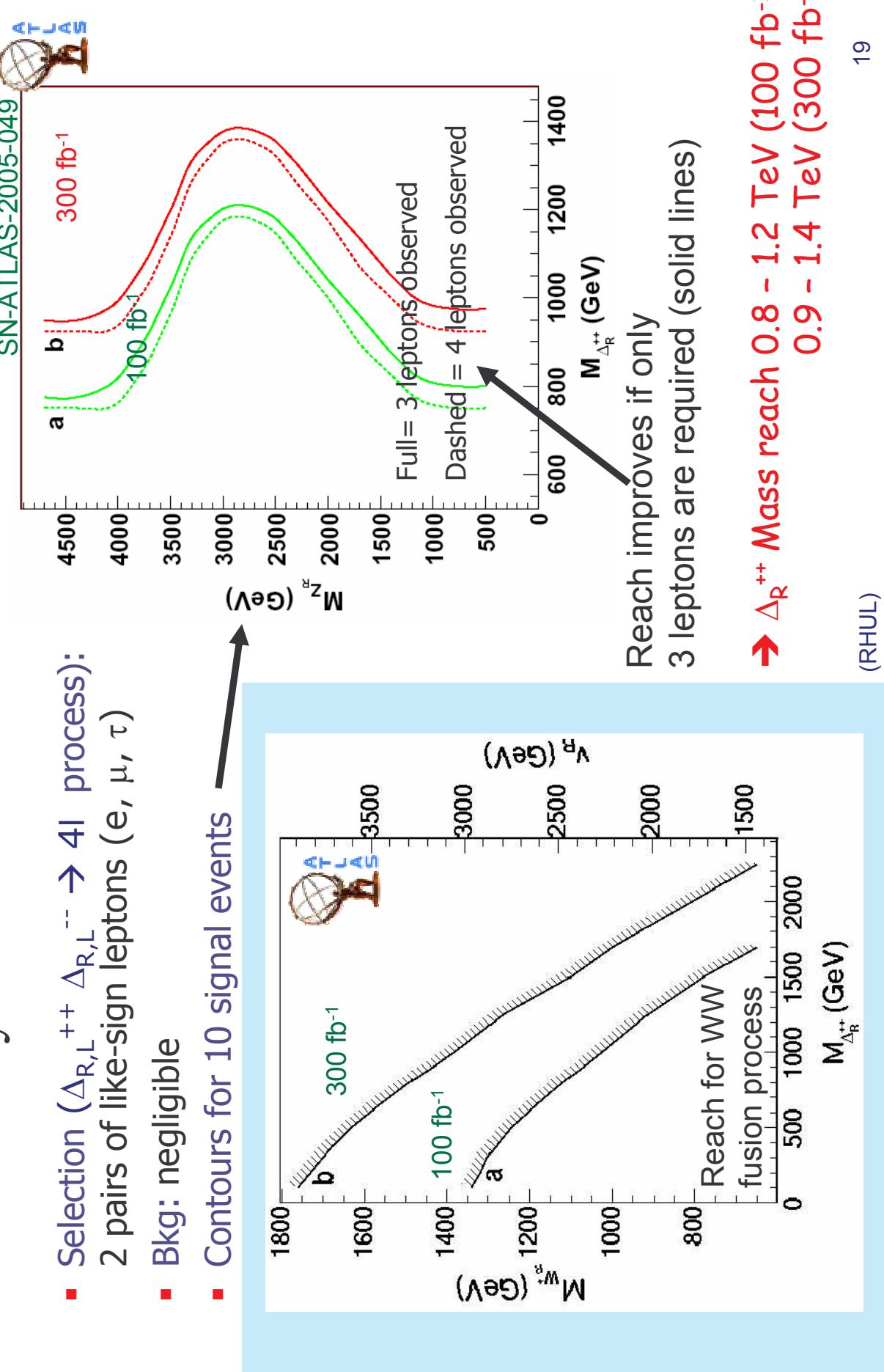
$$q\bar{q} \rightarrow q'q' W_{R,L}^+ W_{R,L}^+ \rightarrow q'q' \Delta_{R,L}^{++}$$

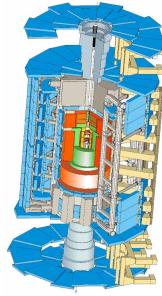
$$q\bar{q} \rightarrow \gamma^*/Z/Z_{R,L} \rightarrow \Delta_{R,L}^{++} \Delta_{R,L}^{--}$$
- Decay: $\Delta_{R,L}^{++} \rightarrow l^+ l^+$
- Selection (WW fusion):
 2 like-sign leptons (e, μ, τ)
 + "forward" jets
- Bkg: $W^+ W^+ q q, W t \bar{t}$



Doubly Charged Higgs in the E-R Symmetric Model

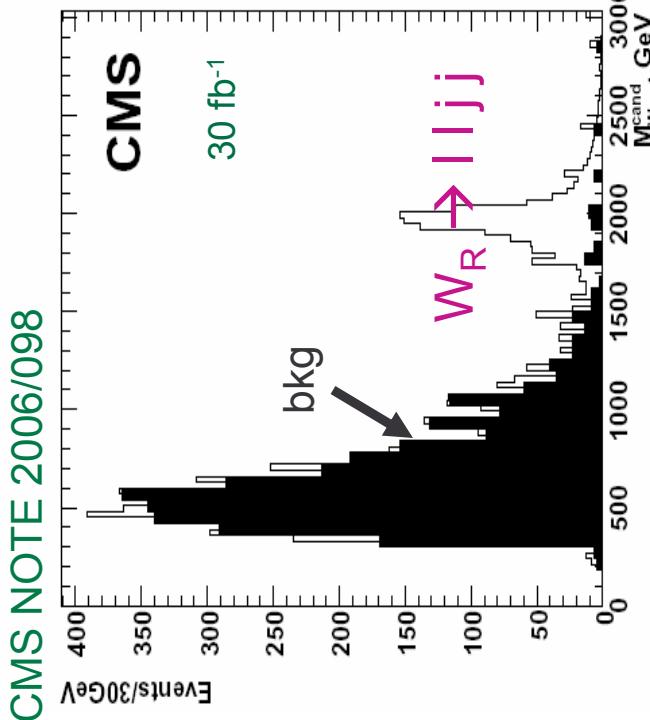
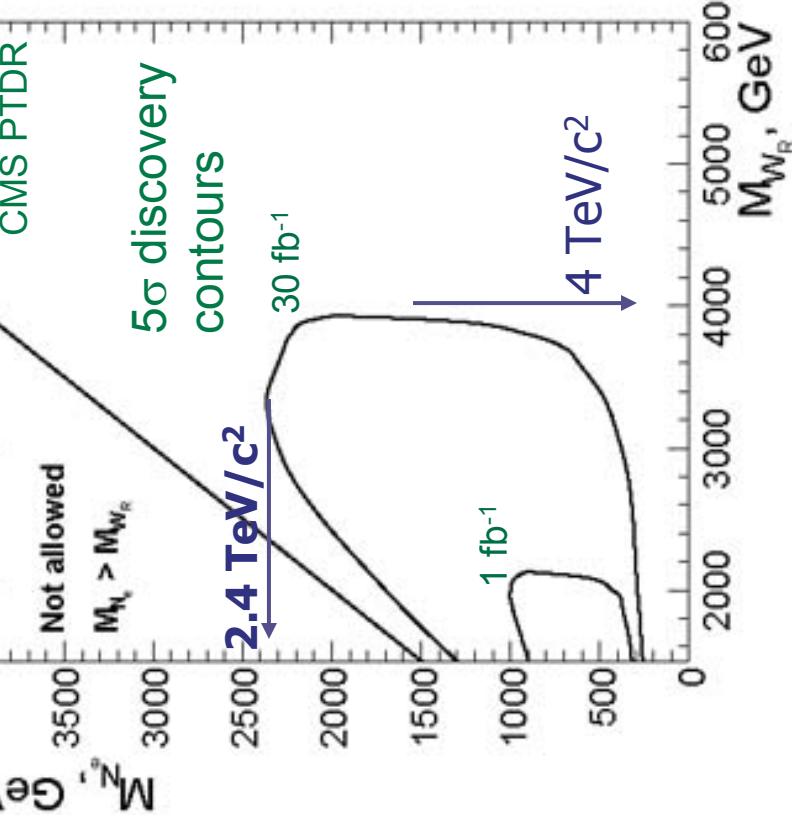
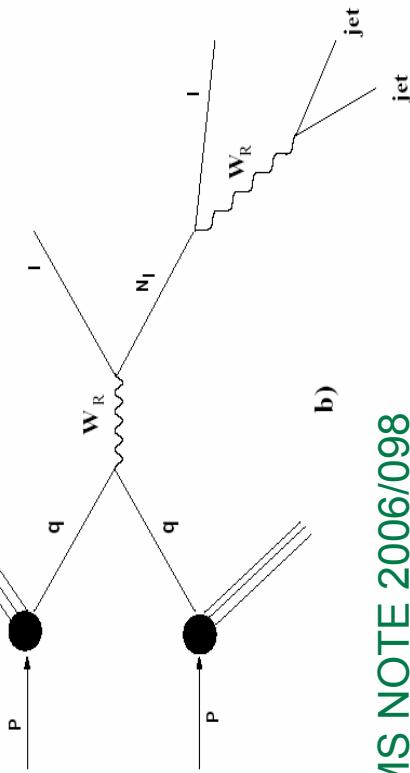
- Selection ($\Delta_{R,L}^{++} \Delta_{R,L}^{--} \rightarrow 4l$ process):
2 pairs of like-sign leptons (e, μ, τ)
 - Bkg: negligible
 - Contours for 10 signal events



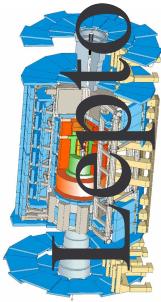


W_R and Majorana Neutrinos

- Left-right Symmetric Model
- Signature:
 - lepton + 2 jets for heavy neutrino N_i
 - dilepton + 2 jets for W_R



Even after just 1 month of running at low lum. can achieve $W_R = 2 \text{ TeV}$ $N_e = 500 \text{ GeV}/c^2$

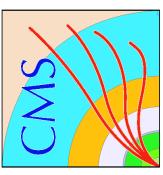


Leptoquarks

SM extension: lepton-quark symmetry

Generic prediction of GUTs, composite models, technicolor schemes, superstring-inspired E_6 models & SUSY with R-parity violation

⇒ LQ: color triplets with couplings to quarks and leptons



Study: Pair Production of Scalar LQ

$LQ\bar{LQ} \rightarrow l^+q\bar{q}$

Signature: 2 jets + 2 leptons: 1st and 2nd gen.

Sensitivity: $M_{LQ} \sim 1.3 \text{ TeV } \beta = 1$

$M_{LQ} \sim 1 \text{ TeV } \beta = 0.5 \text{ L} = 30 \text{ fb}^{-1}$

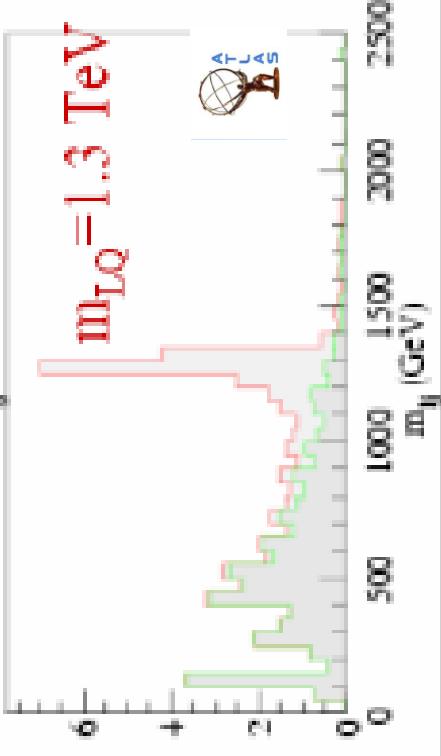
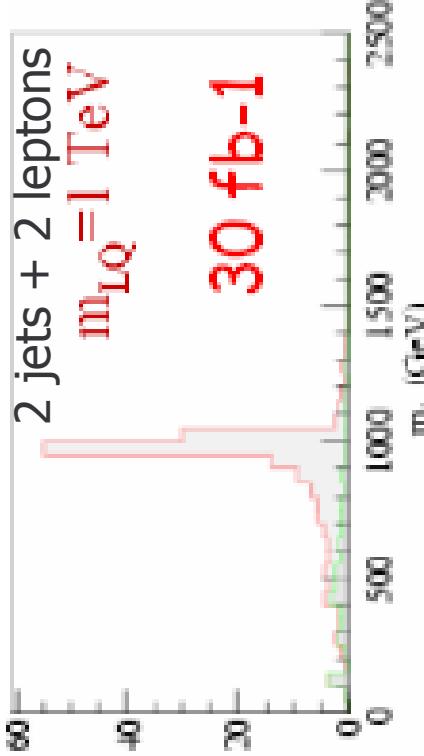
$\beta = BR(LQ \rightarrow lq)$

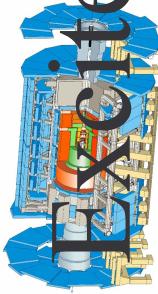
$LQ\bar{LQ} \rightarrow \nu_\tau b\bar{\nu}_\tau \bar{b}$

Signature 2 jets + Et : 3rd generation

Sensitivity: $M_{LQ} \sim 1 \text{ TeV } \beta = 0.5$

$M_{LQ} \sim 1.3 \text{ TeV } \beta = 0$





Excited Quarks

If quarks & leptons have substructure...
there should exist new interactions among them at the scale of
the constituents binding energies

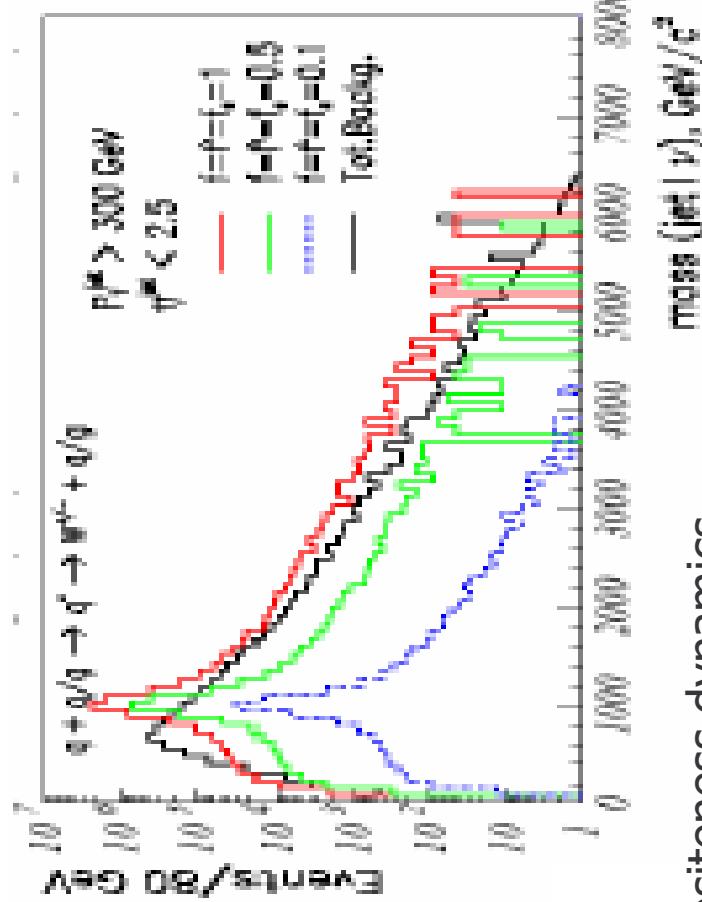
→ a possible consequence is the existence of excited quarks/leptons
with masses of order the compositeness scale Λ .

Excited Quarks

Production: quark-gluon fusion

Decay Signature: photons+jets
quarks + gauge bosons
quarks + gluons

- $L=300\text{fb}^{-1}, \Lambda=m^*, f=f'=1$
- Reach limit for $q^*\rightarrow q\gamma$: 6.5 TeV
 - Reach limit for $q^*\rightarrow qW$: 7 TeV
 - Reach limit for $q^*\rightarrow qZ$: 4.5 TeV
 - Reach limit for $q^*\rightarrow qq$: 6 TeV



$f =$ Coupling determined by the compositeness dynamics
ATL-PHYS-PUB-2006-027

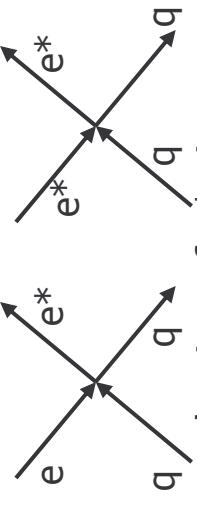
Flavour Workshop, CERN

Tracey Berry (RHUL)



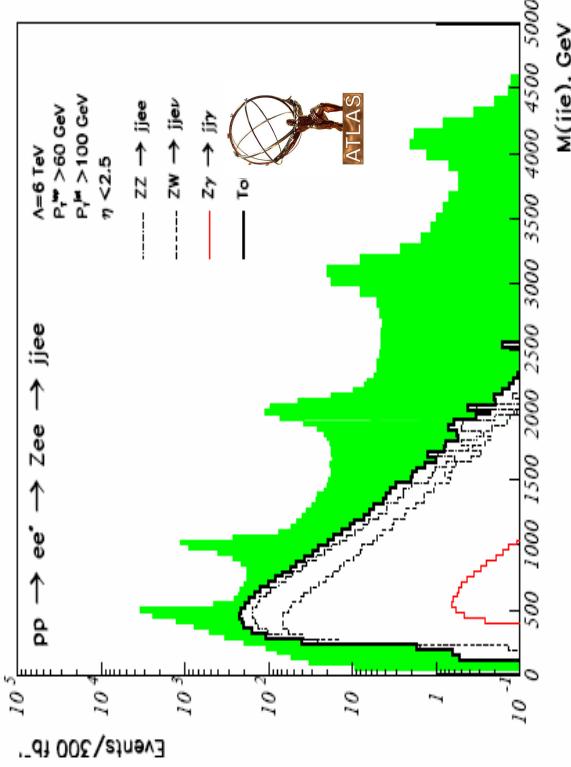
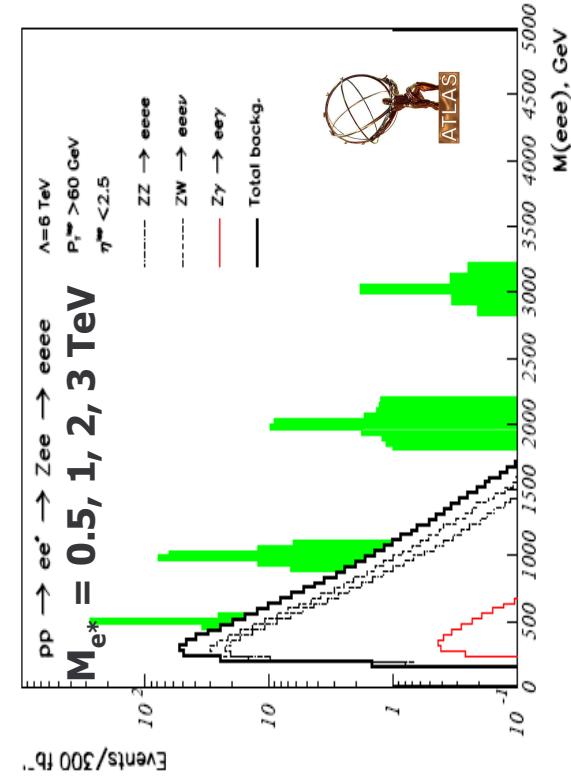
Excited Electrons

At $E < \Lambda$ compositeness scale Λ , quark-lepton interactions can be effectively approximated by contact interactions



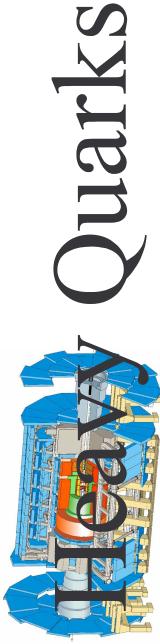
- Production and decay: $qq \rightarrow e^*e \rightarrow Zee$
 $\& Z \rightarrow ee \text{ or } jj$
- Primary backgrounds: $ZZ, Z + \text{photon}$

Single production of e^* via contact interactions
Decay via gauge interactions

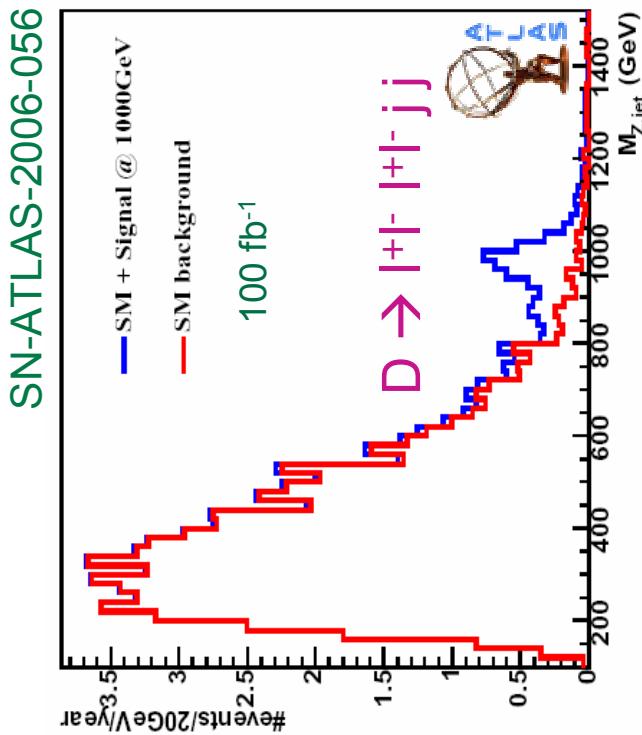
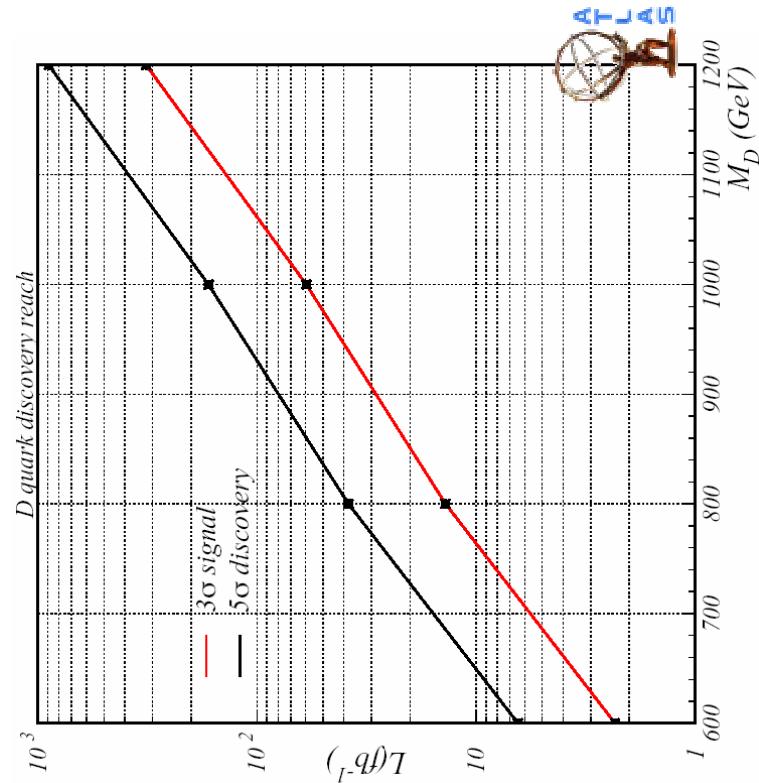
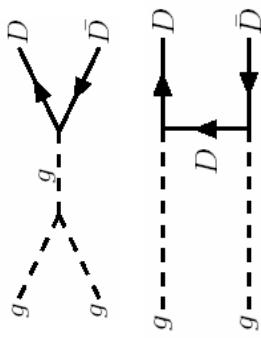
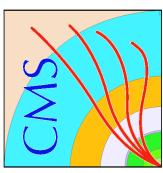


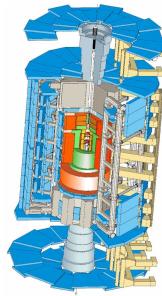
Reach ~3-4 TeV for $L=6 \text{ TeV} 300 \text{ fb}^{-1}$

ATLAS-PHYS-2002-014

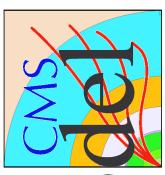


- Symmetry group E_6 favored by string-inspired GUTs (supergravity)
- Predicts new $Q=-1/3$ quark
- Production: $gg \rightarrow D\bar{D}$ (dominant for $M_D < 1.1$ TeV)
 $\bar{q}q \rightarrow D\bar{D}$ (dominant for $M_D > 1.1$ TeV)
- Decay: $D \rightarrow W^+ u$ or $D \rightarrow Z^+ d$ (for this study)
- Selection: 4 leptons (from Z)
+ 2 jets





Extra Dimensions ADD



Arkani-Hamed, Dimopoulos, Dvali, Phys Lett B429 (98), Nucl.Phys.B544(1999)

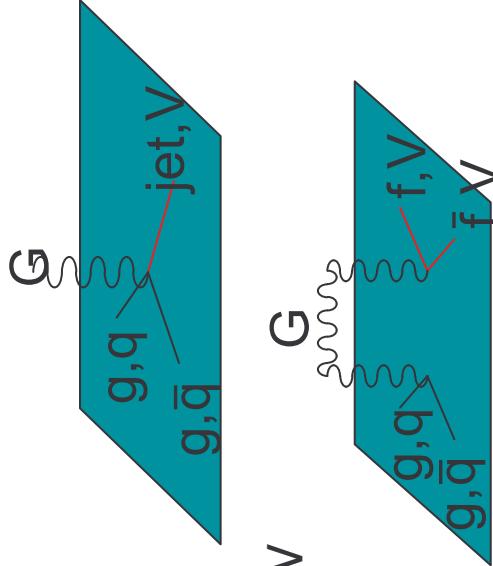
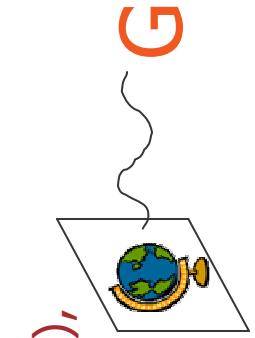
(Many) Large flat Extra-Dimensions (LED),

could be as large as a few μm

G can propagate in ED

SM particles restricted to 3D brane

The fundamental scale is not planckian: $M_D = M_{Pl(4+\delta)} \sim \text{Tev}$



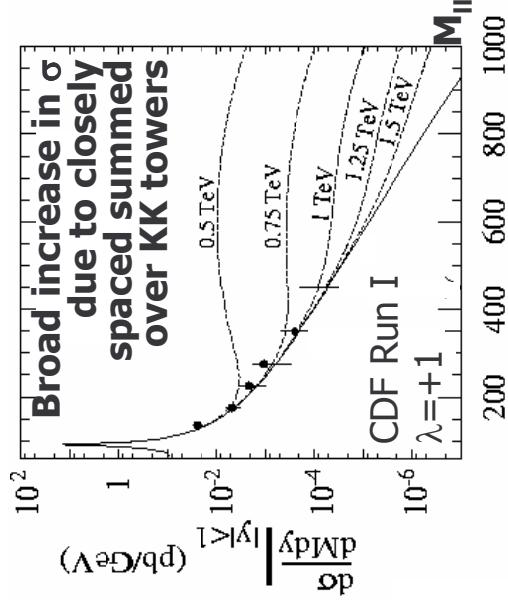
Model parameters are:

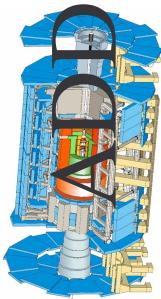
$$M_{Pl}^{-2} \sim R^\delta M_{Pl(4+\delta)}^{(2+\delta)}$$

- δ = number of ED
- $M_{Pl}(4+\delta)$ = Planck mass in the 4+ δ dimensions

Signature:

- ❖ Real graviton emission: **in association with a vector-boson jets + missing ET, V+missing ET**
- ❖ Or deviations in virtual graviton exchange
e.g. Excess above di-lepton continuum





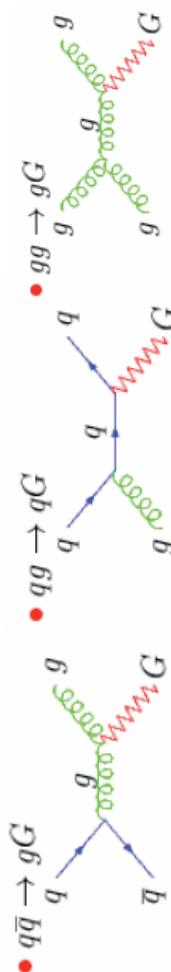
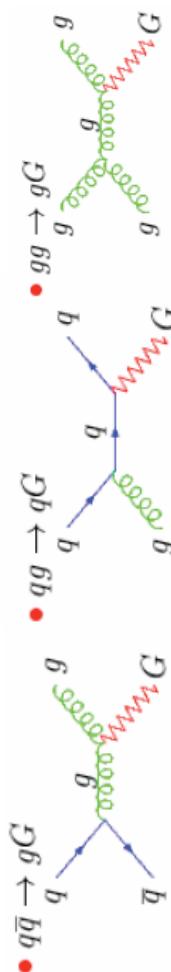
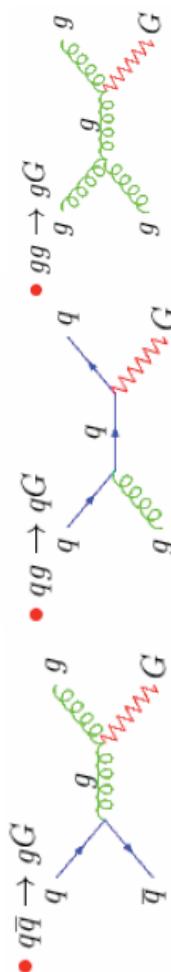
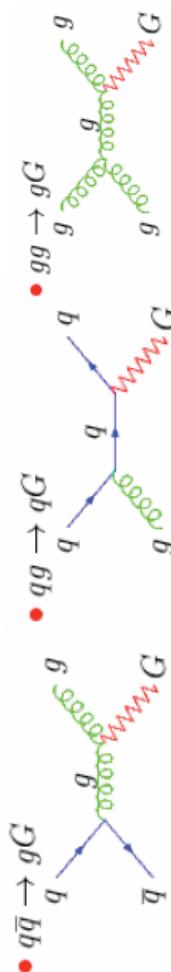
Discovery Limit: Real G Emission

$\text{pp} \rightarrow \gamma + \text{G}^{\text{KK}}$

Signature: high- p_T photon + high missing E_T

Main Bkgd: irreducible $Z\gamma \rightarrow \nu\nu\gamma$,
Also $W \rightarrow e(\mu, \tau)\nu, W\gamma \rightarrow e\nu, \gamma + \text{jets}$,

QCD, di- γ , $Z^0 + \text{jets}$

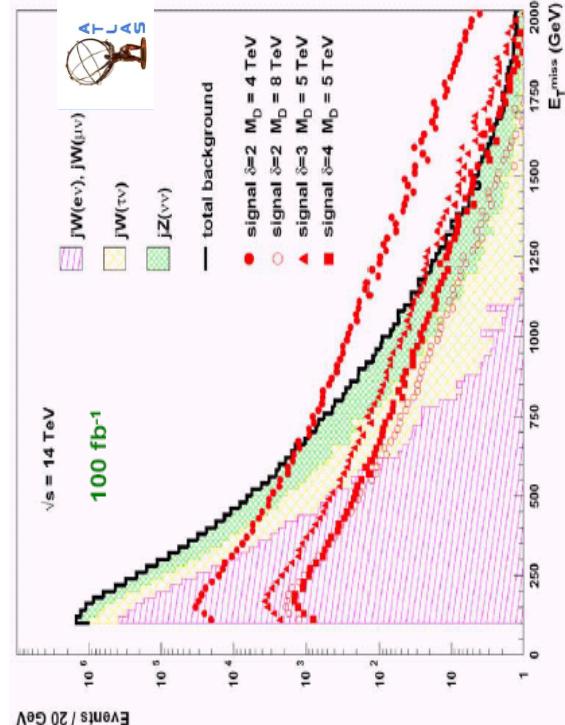


J. Weng et al. CMS NOTE 2006/129

$M_D =$	1 - 1.5 TeV for 1 fb^{-1}
	2 - 2.5 TeV for 10 fb^{-1}
	3 - 3.5 TeV for 60 fb^{-1}

Rates for $M_D \geq 3.5 \text{ TeV}$ are very low –
too low for 5 σ discovery

$M_D^{\text{MAX}} (\text{TeV})$	$\delta=2$
HL 100 fb^{-1}	4



$\text{pp} \rightarrow \text{jet} + \text{G}^{\text{KK}}$

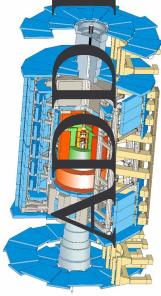
Signature: high E_T jet + large missing E_T

Bkgd: irreducible jet+Z/W \rightarrow jet+vv /jet+ $\nu\nu$

veto leptons: to reduce jet+W bkgd mainly

Discovery limits

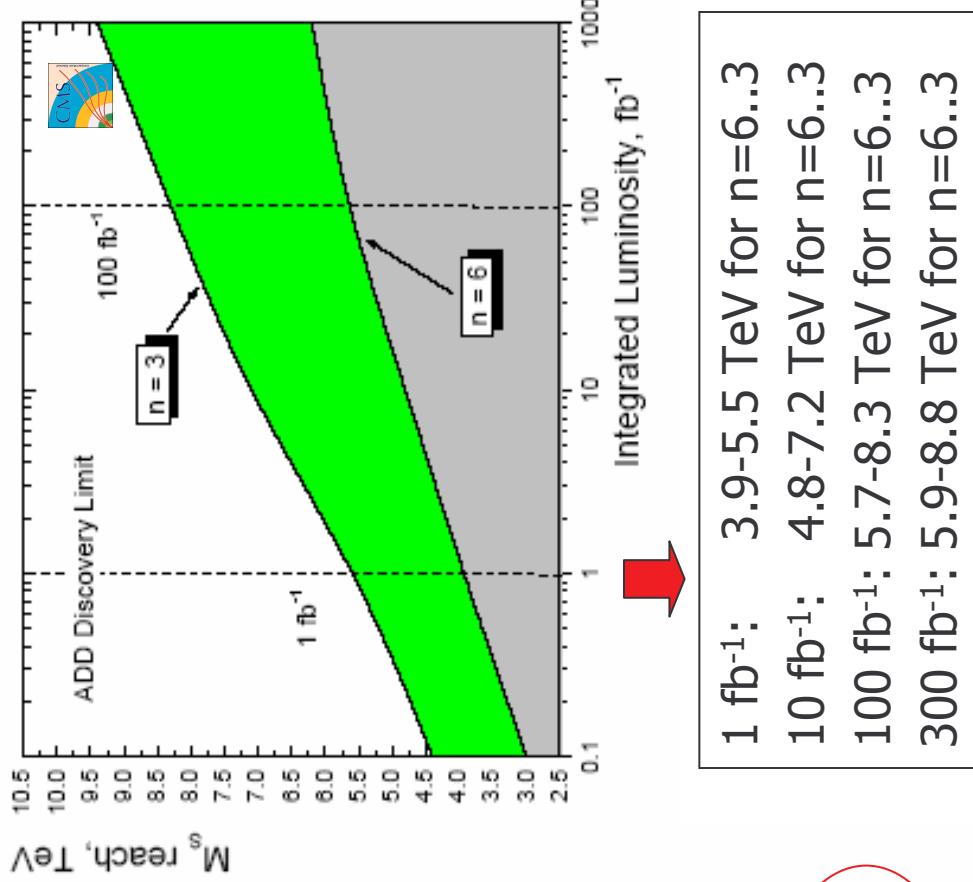
$M_{\text{Pl}}^{\text{MAX}} (\text{TeV})$	$\delta=2$	$\delta=3$	$\delta=4$
LL 30 fb^{-1}	7.7	6.2	5.2
HL 100 fb^{-1}	9.1	7.0	6.0



Add Discovery Limit: G Exchange

$pp \rightarrow GKK \rightarrow \mu\mu$

- Two opposite sign muons & $M_{\mu\mu} > 1$ TeV
- **Bkg:** Irreducible Drell-Yan, also ZZ, WW, WW, tt (suppressed after selection cuts)



channel	n	M_S^{\max} (TeV) S/B	2	3	4	5
$\gamma\gamma$	10 fb^{-1}	6.3 36/18	5.6 36/18	5.1 39/25	4.9 34/13	
$\gamma\gamma$	100 fb^{-1}	7.9 50/53	7.3 62/96	6.7 55/72	6.3 51/53	
t^+t^-	10 fb^{-1}	6.6 33/11	5.9 31/8	5.4 30/6	5.1 30/6	
t^+t^-	100 fb^{-1}	7.9 49/48	7.5 38/21	7.0 36/16	6.6 29/6	
Fast MC	10 fb^{-1}	M_S^{\max} (TeV)	7.0	6.3	5.7	5.4
$\gamma\gamma + l^+l^-$	100 fb^{-1}	M_S^{\max} (TeV)	8.1	7.9	7.4	7.0

V. Kabachenko et al.
ATL-PHYS-2001-012

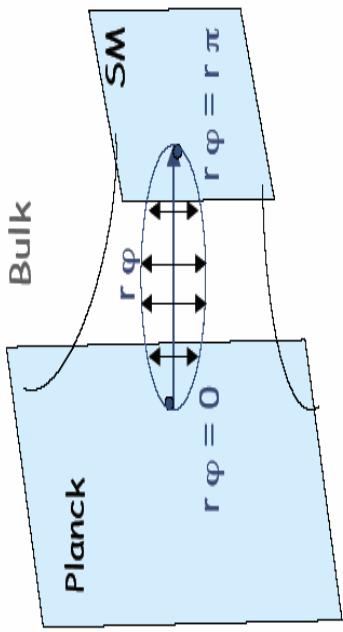
Belotelov et al.,
CMS NOTE 2006/076, CMS PTDR 2006

L)



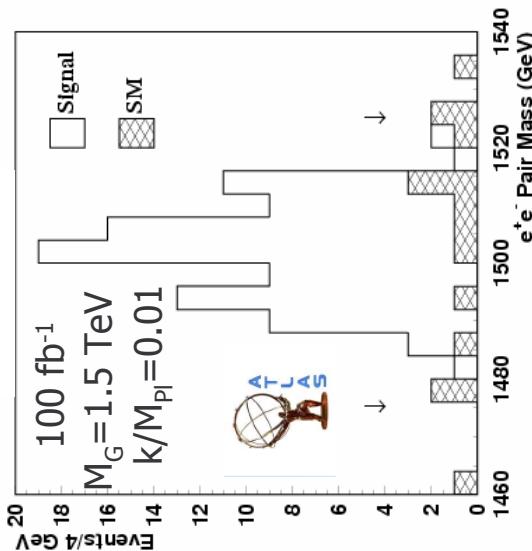
RS

1 extra
warped
dimension



Signature:
Narrow, high-mass resonance states in
dilepton/dijet/diboson channels

G₁→e⁺e⁻



Best channels to search in are

$G(1) \rightarrow e^+e^-$ and $G(1) \rightarrow \gamma\gamma$ due to the energy and angular resolutions of the LHC detectors

$G(1) \rightarrow e^+e^-$ best chance of discovery
 $G(1) \rightarrow \gamma\gamma$ due to relatively small bkgd, from
Drell-Yan

Allenach et al, hep-ph0211205 Allenach et al, hep-ph0006114

Tracey Berry (RHUL)

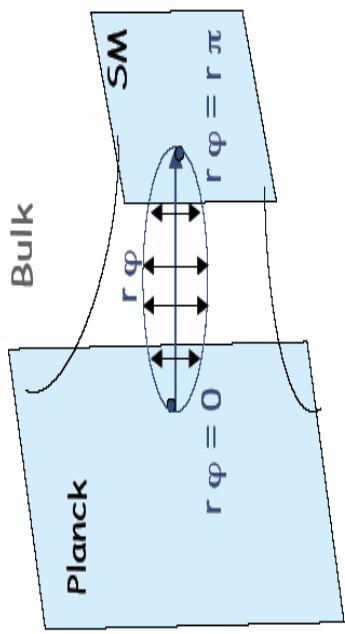
Flavour Workshop, CERN



RS

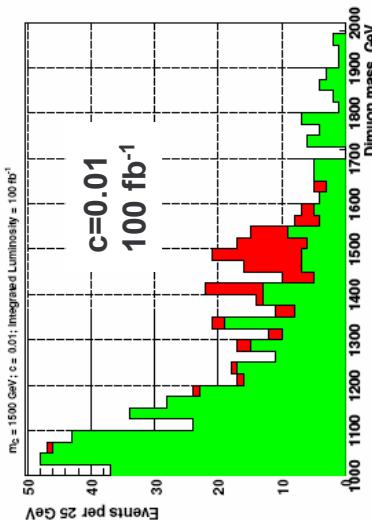
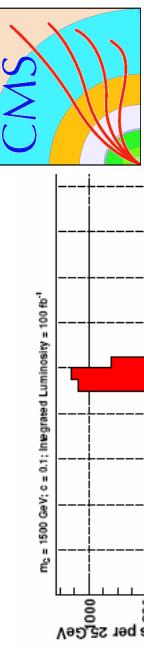
Model

1 extra
warped
dimension



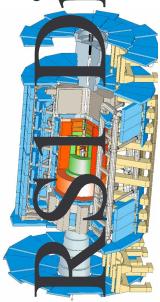
Signature:

Narrow, high-mass resonance states in
dilepton/dijet/diboson channels

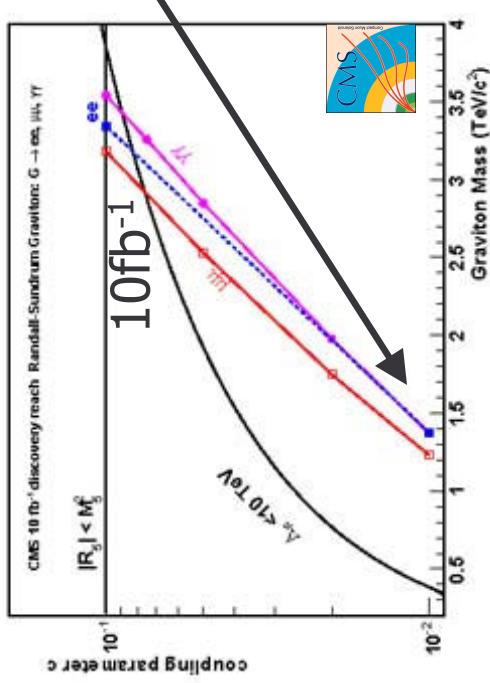


- Best channels to search in are $G(1) \rightarrow e^+e^-$ and $G(1) \rightarrow \gamma\gamma$ due to the energy and angular resolutions of the LHC detectors
- $G(1) \rightarrow e^+e^-$ best chance of discovery due to relatively small bkgd, from Drell-Yan

- Model parameters: $\Lambda_\pi = \overline{M}_p e^{-kR_c \pi}$
- Gravity Scale: $\Lambda_\pi = \overline{M}_p e^{-kR_c \pi}$ **Resonance position**
1st graviton excitation mass: $m_1 \rightarrow \text{position}$
 $\Lambda_\pi = m_1 \overline{M}_p / k x_1$, & $m_n = k x_n e^{k R_c \pi} (J_1(x_n) = 0)$
 - Coupling constant: $c = k/M_p$
 $\Gamma_1 = \rho m_1 x_1^2 (k/M_p)^2 \rightarrow \text{width}$
 $k = \text{curvature}, R = \text{compactification radius}$



R&D Discovery Limit

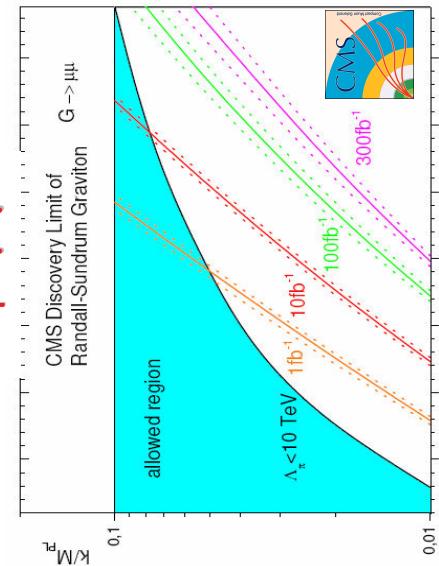


$$\text{BR}(G \rightarrow \gamma\gamma) = 2 * \text{BR}(G \rightarrow ee/\mu\mu)$$

Reach for e and γ is comparable for low coupling and G mass due to the QCD and prompt photon backgrounds in the $\gamma\gamma$ channel which are harder to efficiently suppress $\mu\mu$ channel trails ee channel due to resolution

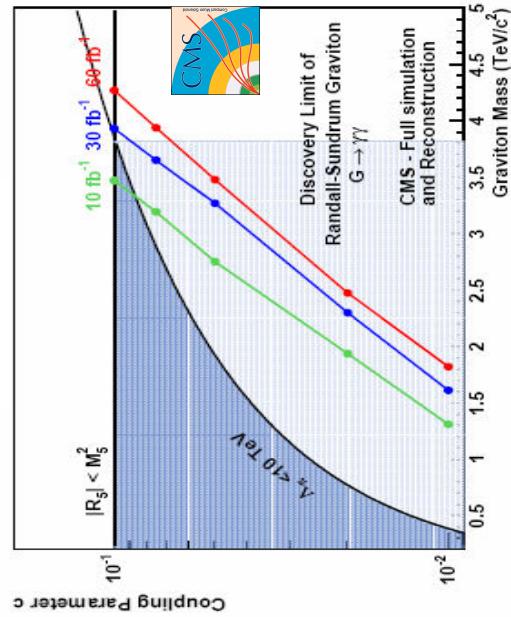
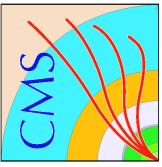
< 60 fb⁻¹ LHC completely covers the region of interest

$G_1 \rightarrow \mu^+ \mu^-$



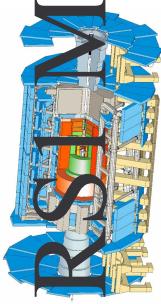
$c > 0.1$ disfavoured as bulk curvature becomes to large (larger than the 5-dim Planck scale)

Flavour Workshop, CERN



CMS PTDR 2006

Tracey Berry (RHUL)

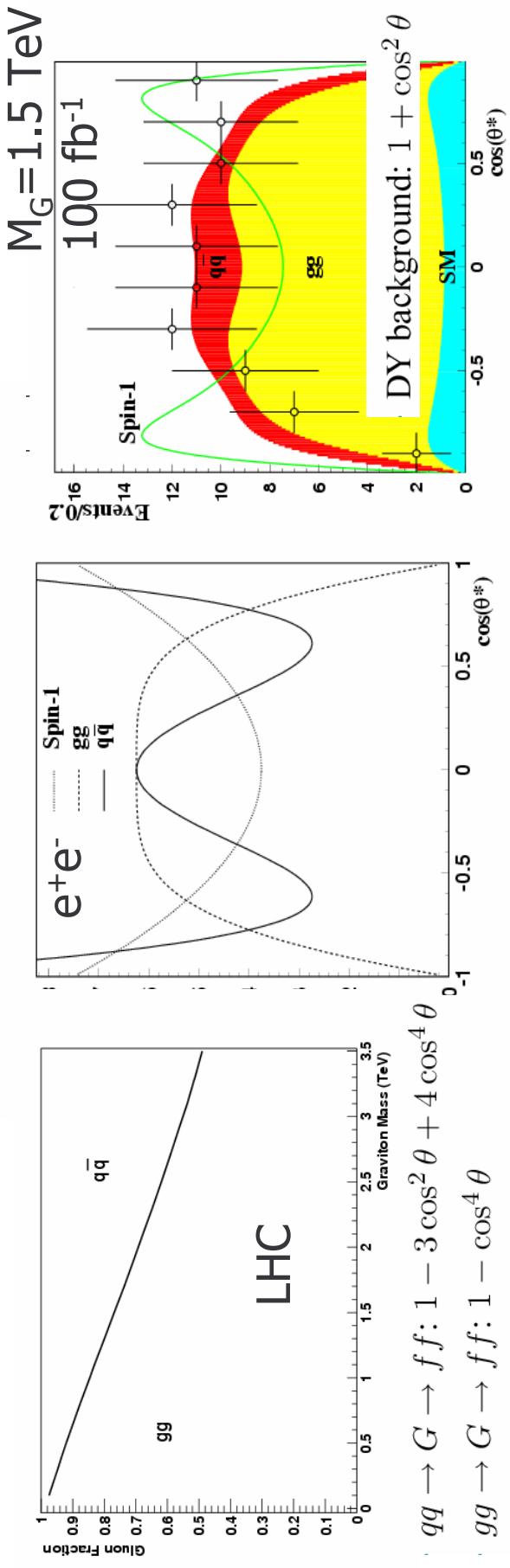


RS Model Determination

How could a RS G resonance be distinguished from a Z' resonance?

Potentially using Spin information:

G has spin 2: $pp \rightarrow G \rightarrow ee$ has 2 components: $gg \rightarrow G \rightarrow ee$ & $q\bar{q} \rightarrow G \rightarrow ee$:
each with different angular distributions:

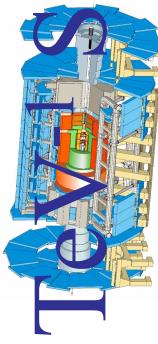


**Spin-2 could be determined (spin-1 ruled out) with 90% C.L.
up to $M_G = 1720 \text{ GeV}$ with 100 fb^{-1}**

Note: acceptance at large pseudo-rapidities is essential for spin discrimination
($1.5 < |\eta| < 2.5$)

Allanach et al, hep-ph 0006114

Tracey Berry (RHUL)



TeV SIZED Extra Dimensions

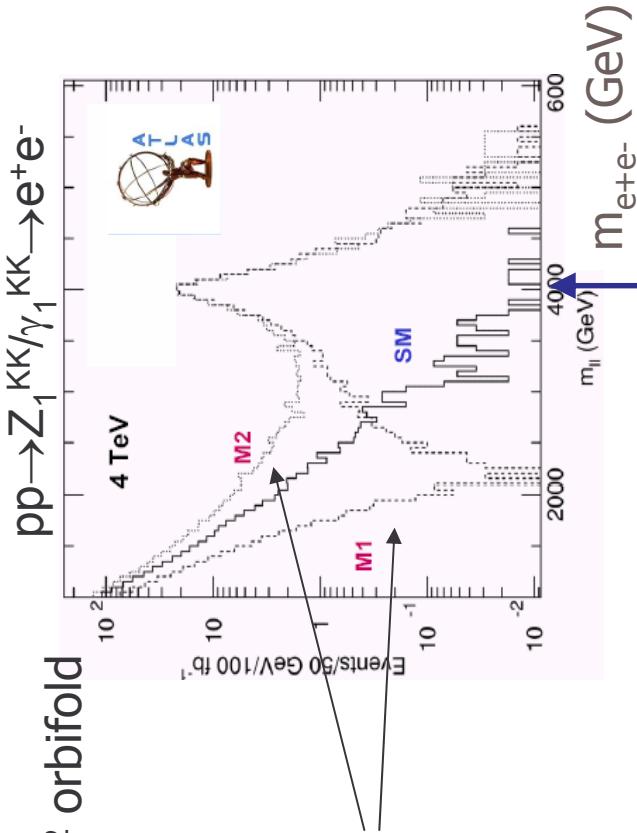
I. Antoniadis, PLB246 377 (1990)

- One extra dimension compactified on a S^1/Z^2 orbifold
- Radius of compactification small enough ->
Gauge bosons can travel in the bulk
- Fermions (quarks/leptons) localized
 - at a fixed point (M_1) or
 - opposite (M_2) points
- ⇒ destructive (M_1) or constructive (M_2) interference of the KK excitations with SM model gauge bosons

Signature:

KK excitations of the gauge bosons ($Z^{(k)}, W^{(k)}$) appear as resonances with masses :
 $M_k = \sqrt{(M_0^2 + k^2/R^2)}$ where ($k=1, 2, \dots$) & also interference effects!

- Look for I^+I^- decays of γ and Z^0 KK modes.
Also in decays (m_T) of $W^{+/}$ KK modes.
Or evidence of g^* via dijet σ or bb, tt s



$$M_n = M_0$$

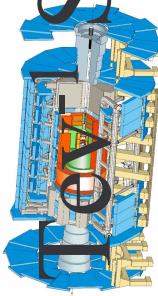
New Parameters

- $R = M_C^{-1}$: size of compact dimension
- M_C : compactification scale
- M_0 : mass of the SM gauge boson

G. Azuelos, G. Polesello
EPJ Direct 10.1140 (2004)
32

Tracey Berry (RHUL)

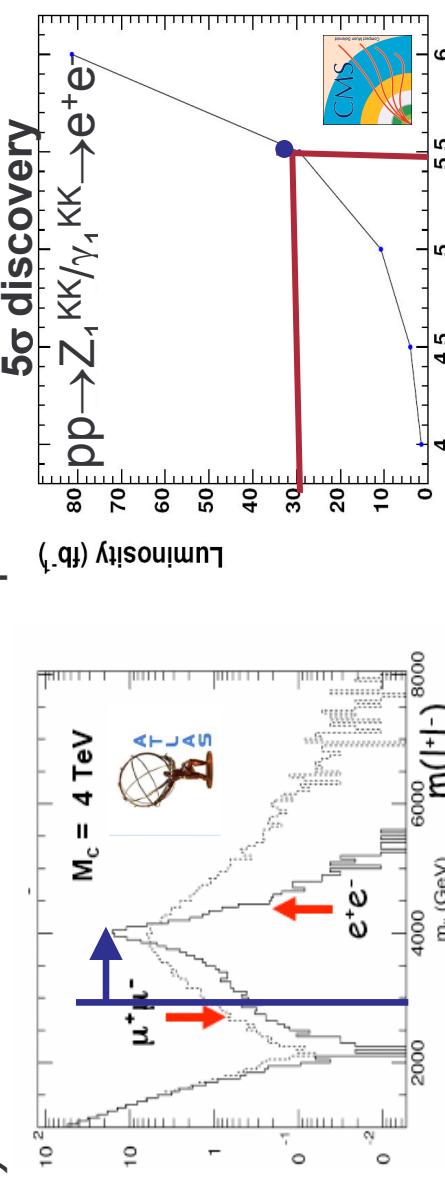
Flavour Workshop, CERN



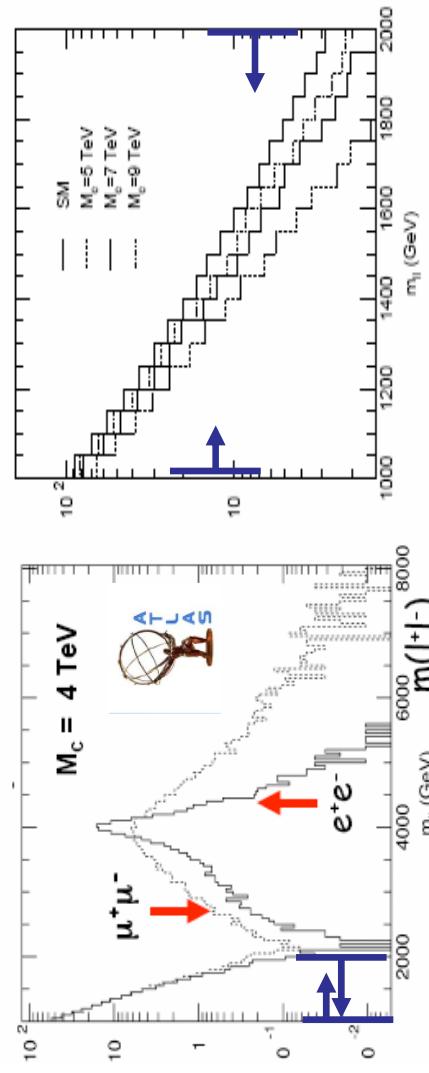
TeV-Sized Extra Dimensions

Look for resonances/deviations in the \parallel spectrum

1) Search for the resonance peak



2) Search for interference in a mass window



$ee + \mu\mu$: ATLAS 5 σ reach is ~ 8 TeV for $\mathcal{L}=100$ fb⁻¹ and ~ 10.5 TeV for 300 fb⁻¹

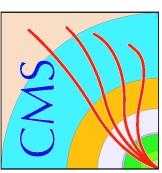
Flavour Workshop, CERN

Tracey Berry (RHUL)

$\mathcal{L}=30/80$ fb⁻¹ CMS will be able to detect a peak in the e^+e^- invariant mass if $M_c < 5.5/6$ TeV.

ATLAS: $\mathcal{L}=100$ fb⁻¹
 $M_c(R^{-1}) < 5.8$ TeV
 $(ee + \mu\mu)$

3) Fit to kinematics of signal
 With 300 fb⁻¹ can reach 13.5 TeV ($ee + \mu\mu$)





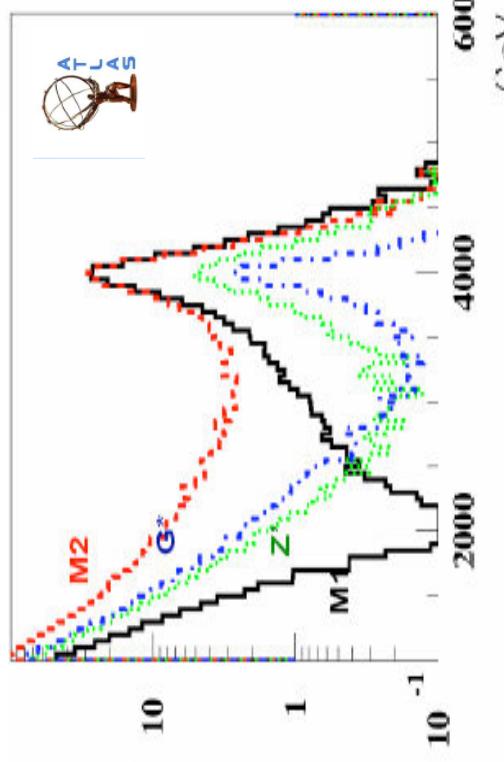
Distinguishing $Z^{(1)}/\gamma^{(1)}$ from Z'/RS

Distinguish

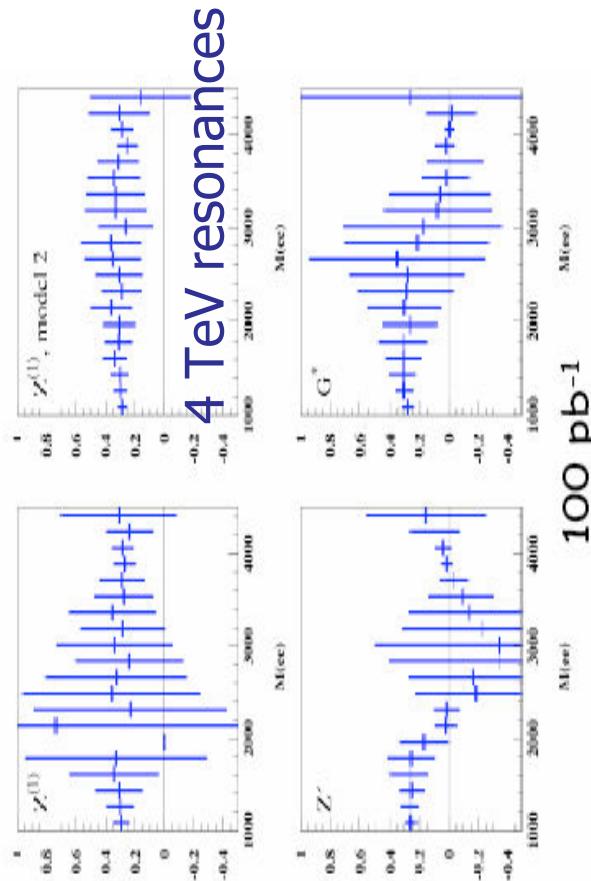
- spin-1 $Z^{(1)}$ from spin-2 G : angular distribution of decay products
- spin-1 $Z^{(1)}$ from spin-1 Z' with SM-like couplings: forward-backward asymmetry due to contributions of the higher lying states, interference terms and additional $\sqrt{2}$ factor in its coupling to SM fermions.

The $Z^{(1)}$ can be discriminated for masses up to about 5 TeV with $L=300\text{fb}^{-1}$.

4 TeV $Z^{(1)})/\gamma^{(1)}$ or Z' or RS Graviton?



Forward-backward asymmetries:

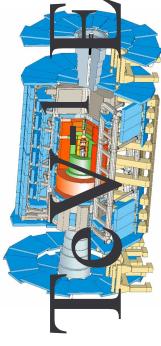


G. Azuelos, G. Polesello
EPJ Direct 10.1140 (2004)

34

Tracey Berry (RHUL)

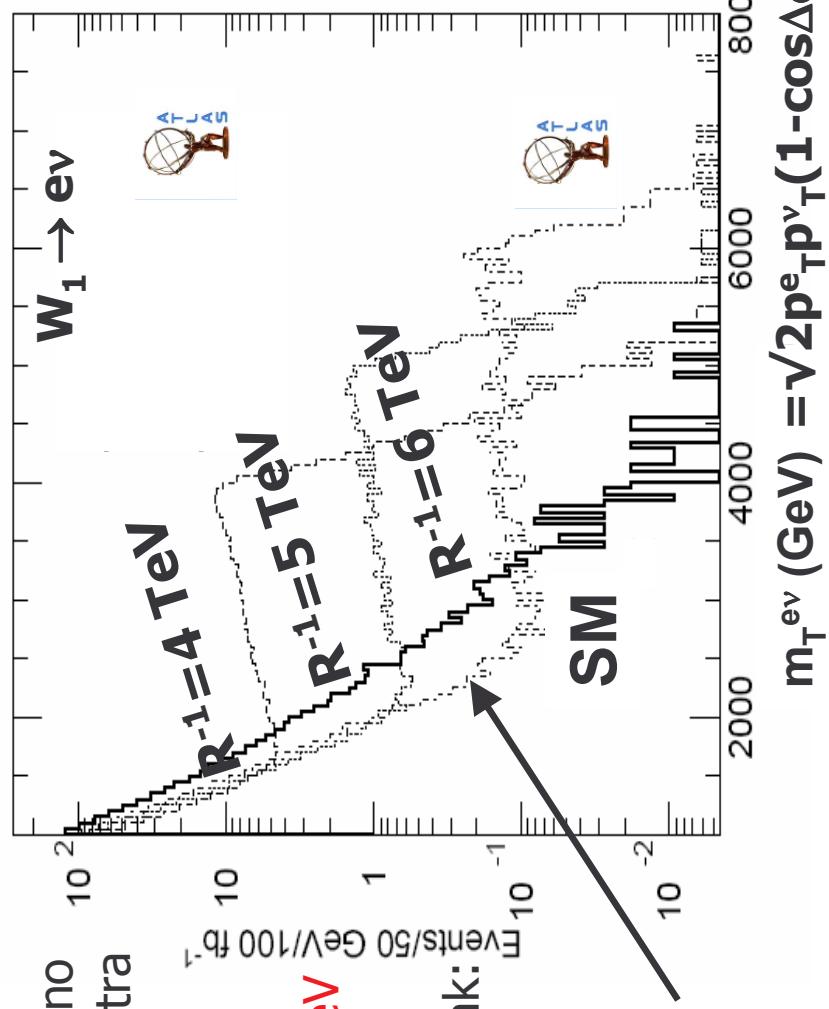
Flavour Workshop, CERN



T_eV_NEED Discovery Limits

W_{KK} decays

Search for deviations in lepton-neutrino transverse invariant mass ($m_T^{l\nu}$) spectra



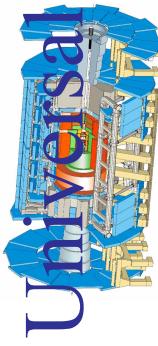
- 1) Search for peak:
 $L=100 \text{ fb}^{-1}$ detect a **peak if compactification scale ($M_C = R^{-1}$) < 6 TeV**
Sum over 2 lepton flavours

- 2) Studying distribution **below** the peak:
in $m_T^{e\nu}$ spectra
 $L=100 \text{ fb}^{-1}$ a limit of **$M_C > 11.7 \text{ TeV}$**

-ve interference sizable even for M_C above the ones accessible to a direct detection of the mass peak.

If a peak is detected, a measurement of the couplings of the boson to the leptons and quarks can be performed for M_C up to $\sim 5 \text{ TeV}$.

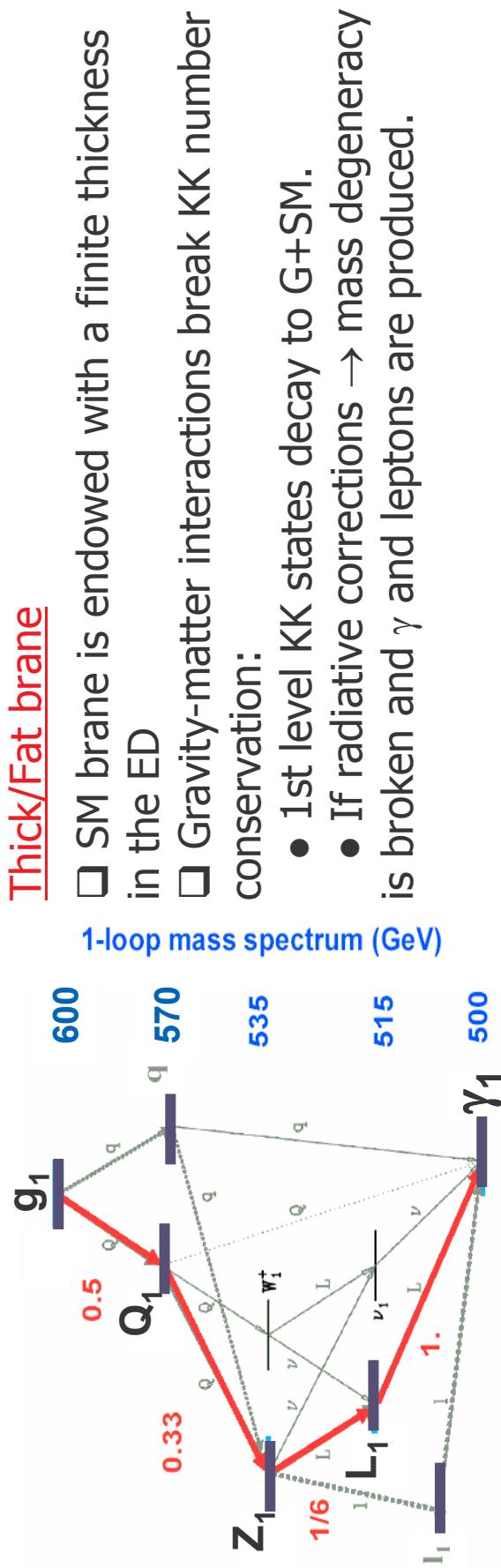
G. Polesello, M. Patra
EPJ Direct, ATLAS 2003-023
G. Polesello, M. Patra
EPJ Direct C 32 Sup.2 (2004) pp.55-67

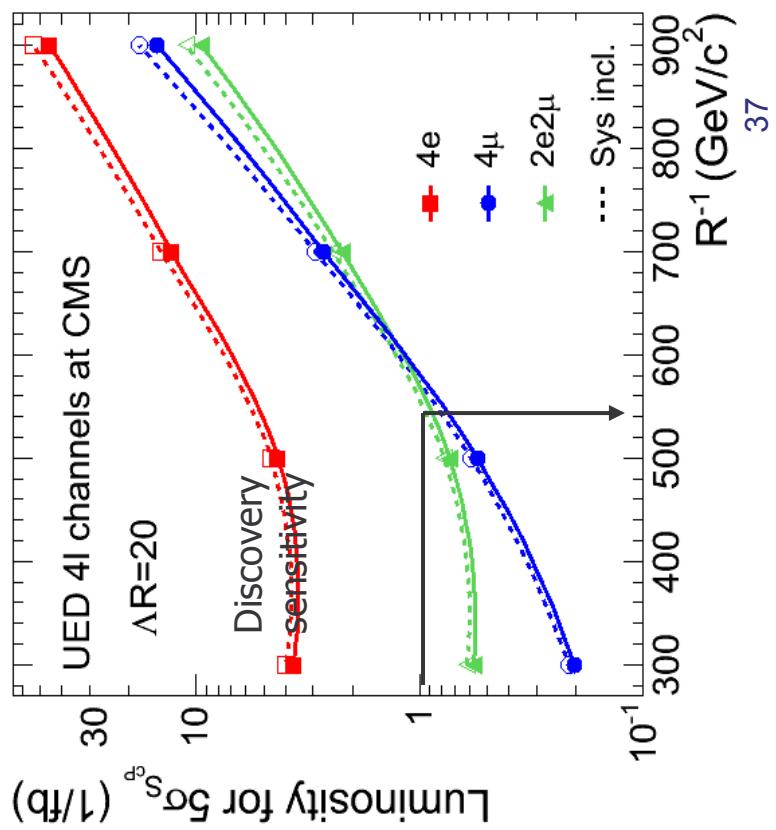
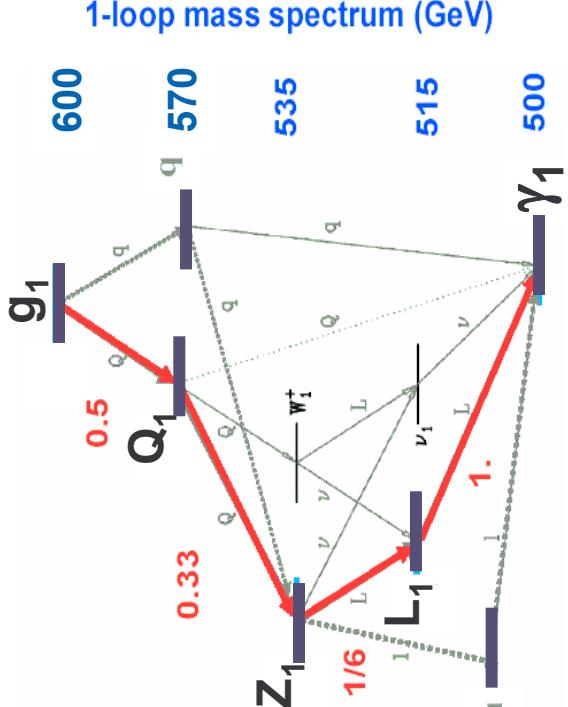
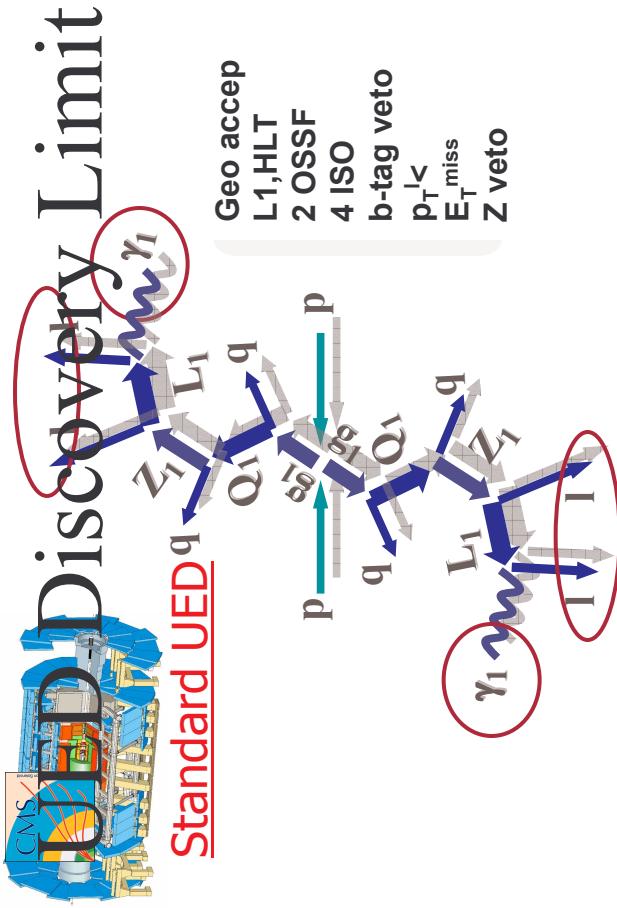


Universal Extra Dimensions Standard/Minimal UED

- All particles can travel into the bulk, so each SM particle has an infinite tower of KK partners
- Spin of the KK particles is the same as their SM partners
- In minimal UED: 1 ED compactified in an orbifold (S_1/Z_2) of size R
 - KK parity conservation → the lightest massive KK particle (LKP) is stable (dark matter candidate).
 - Level one KK states must be pair produced
 - Mass degeneration except if radiative corrections included

The model parameters: compactification radius R , cut-off scale Λ , m_h



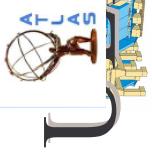


$pp \rightarrow g_1 g_1 \rightarrow 4l + 4q + 2LKP \rightarrow 4l + 4\text{jets} + p_T$
 $pp \rightarrow g_1 Q_1 \rightarrow 4l + 3q + 2LKP \rightarrow 4l + 3\text{jets} + p_T$
 $pp \rightarrow Q_1 Q_1 \rightarrow 4l + 2q + 2LKP \rightarrow 4l + 2\text{jets} + p_T$

Signature: 4 low- p_T isolated leptons (2 pairs of opposite sign, same flavour leptons) + n jets + missing E_T (from 2 undetected γ_1)

Irreducible Bckg: $t\bar{t}\text{bar}$ + n jets($n = 0, 1, 2$),
4 b-quarks, ZZ, Zbbar

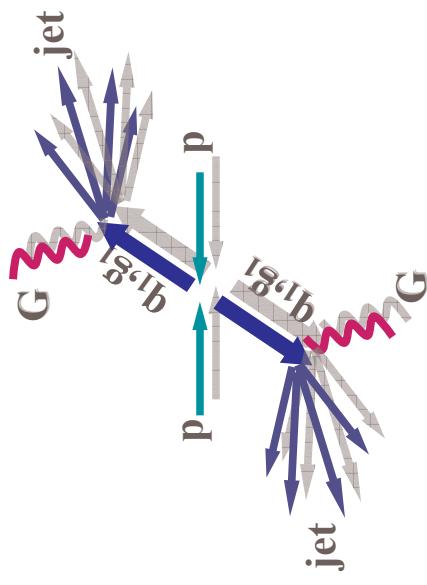
Studied for low lum run $\sim 2 \times 10^{33} \text{ cm}^{-2} \text{s}^{-1}$



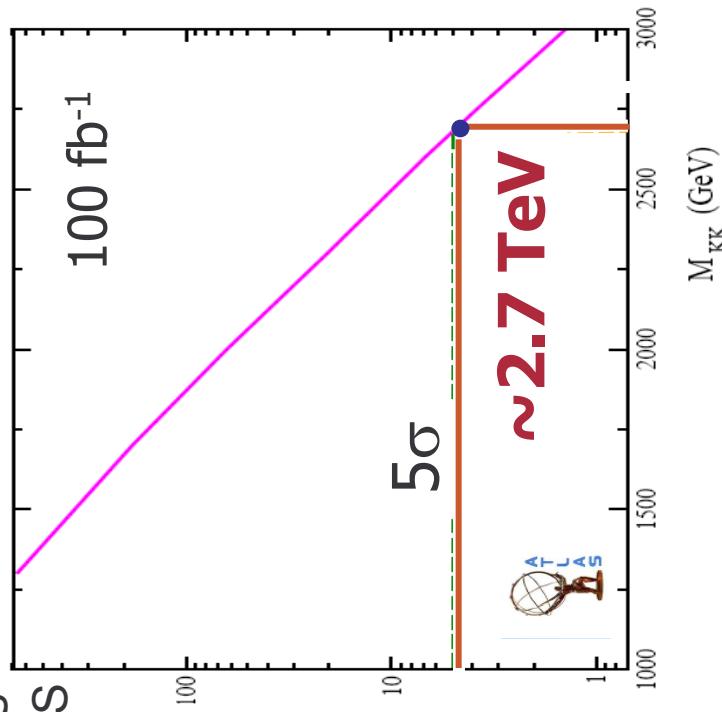
Discovery Limit

Thick brane in UED with TeV⁻¹ size

$$pp \rightarrow g_1 g_1 / q_1 q_1 \rightarrow 2 jets + \cancel{E}_T$$



Significance vs Mass of 1st KK excitation

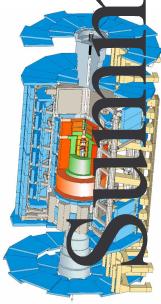


$\sim 2.7 \text{ TeV}$

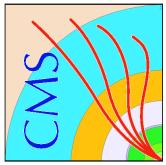
Signature:
2 back-to back jets + missing $E_T (> 775 \text{ GeV})$

Irreducible Bckg: $Z(\rightarrow \nu\nu) jj, W(\rightarrow l\nu) jj$

5σ discovery possible at ATLAS with 100 fb^{-1} if first KK excitation mass $< 2.7 \text{ TeV}$



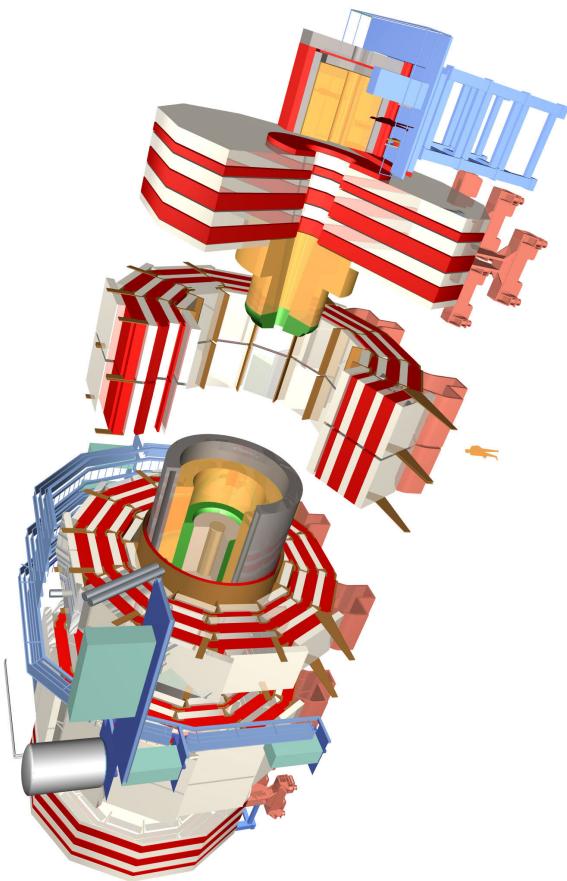
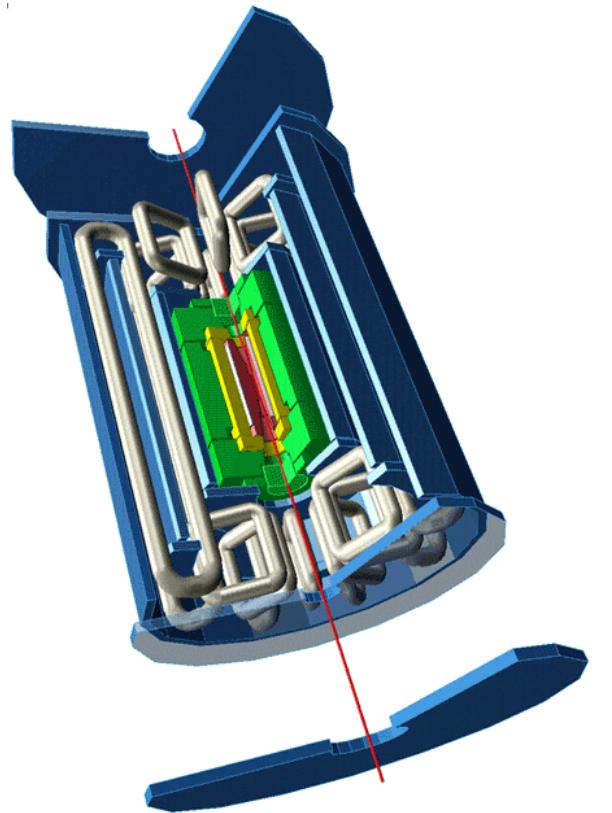
Summary & Outlook



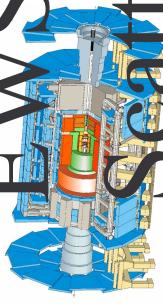
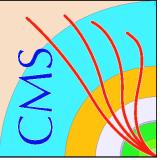
- ATLAS & CMS have significant discovery potential related to fundamental symmetries, Electroweak symmetry breaking, quark-lepton family structure and extra dimensions.
 - Heavy gauge bosons up to \sim 5-6 TeV
 - Little Higgs T quark up to \sim 2 TeV
 - Vector boson resonances; Technihadron p_{Tc} mass up to \sim 600 GeV
 - Doubly-charged Higgs up to \sim 2 TeV
 - Heavy neutrino up to \sim 2.5 TeV,
 - RS Model ED up to \sim 4 TeV, $Z^{(\text{KK})}/\gamma^{(\text{KK})}$ up to \sim 13 TeV
 - *Many more topics not covered*
- ATLAS & CMS increasing focus on first year of data taking
 - Understand/optimize detector performance (calibration, alignment, ...)
 - Understand/measure Standard Model processes (bkg sources)

Once these are achieved ATLAS could potentially have new physics results within months!
- **Eager to start exploration of TeV scale!**

Backup Slides



EWSSB: Resonant Vector Boson Scattering

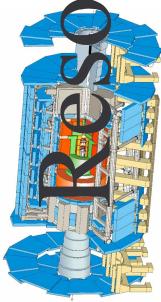


- SM cross section for $W_{\text{long}} W_{\text{long}}$ scattering diverges at high energy if there is no Higgs \rightarrow new physics via diboson resonances?
- Chiral Lagrangian Model

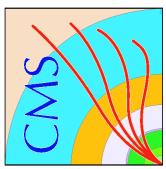
- low-energy effective description of electroweak interactions
 \rightarrow yields interaction terms describing VB scattering with arb. coeffs.
- respects chiral symmetry via $SU(2)_L \otimes SU(2)_R$
- choose parameters such that new resonance $M = 1.15$ TeV

- Study $W Z$ scattering (cleaner than $WW +$ to reconstruct mass):

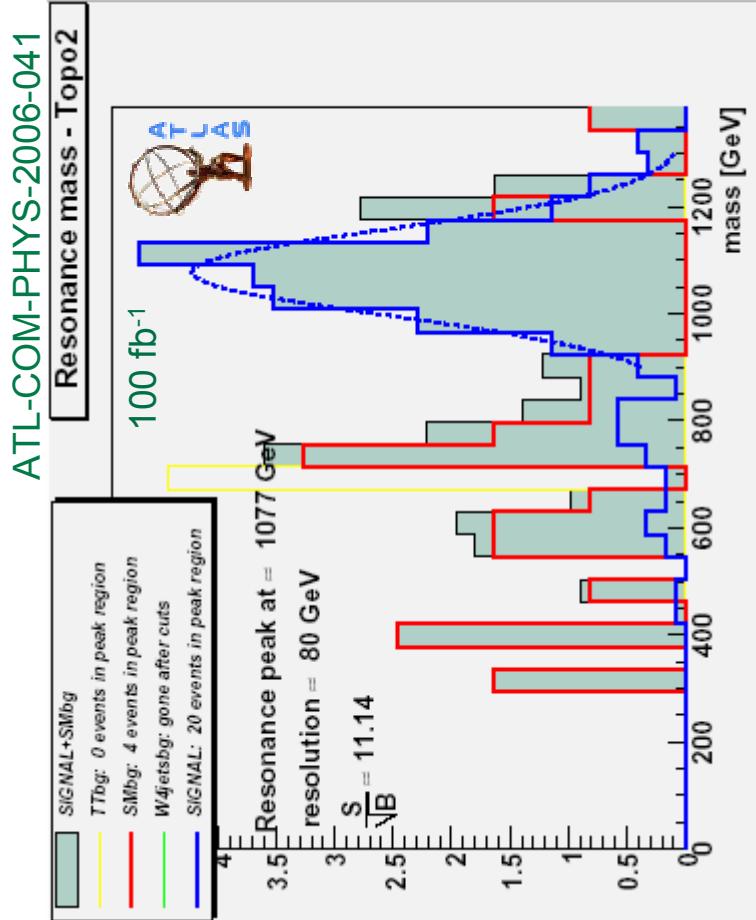
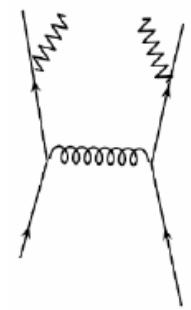
- $qq \rightarrow qqWZ \rightarrow qq l\nu \parallel (\sigma \times BR = 1.3 \text{ fb})$
- $qq \rightarrow qqWZ \rightarrow qq jj \parallel (\sigma \times BR = 4.1 \text{ fb})$
- $qq \rightarrow qqWZ \rightarrow qq l\nu jj (\sigma \times BR = 14 \text{ fb})$



Resonant Vector Boson Scattering



- Selection:
 - 2 forward jets + central jets and/or leptons
 - + missing E_T (for $W \rightarrow l\nu$)
- Require no additional central jet & b-jet veto (for jet mode)
- Bkg: gluon and g/Z exchange with W and Z radiation
 - also $t\bar{t}$ & $W+4$ jets (need more stats)
- Promising sensitivity for
 - jet modes at 100 fb^{-1}
 - (need 300 fb^{-1} for $WZ \rightarrow l\nu ll$)
- study is ongoing



$WZ \rightarrow jj \parallel$

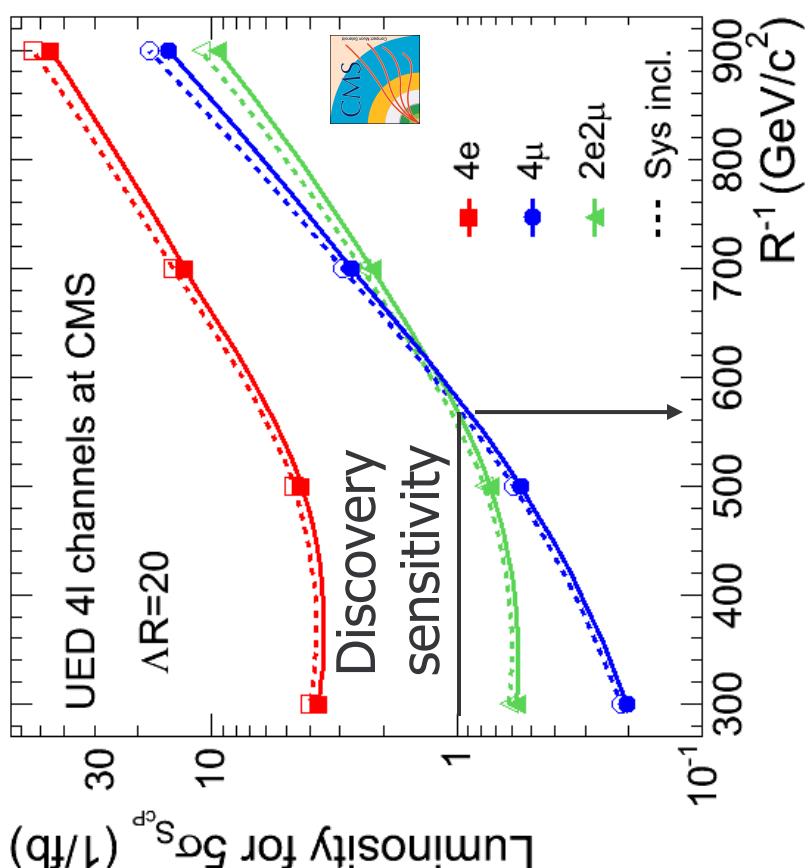


- $pp \rightarrow g_1g_1 \rightarrow 4l + 4q + 2LKP \rightarrow 4l + 4\text{jets} + P_T$
- $pp \rightarrow g_1Q_1 \rightarrow 4l + 3q + 2LKP \rightarrow 4l + 3\text{jets} + P_T$
- $pp \rightarrow Q_1Q_1 \rightarrow 4l + 2q + 2LKP \rightarrow 4l + 2\text{jets} + P_T$

□ 4 leptons in the final state + missing P_T

- Irreducible Bckg: ttbar + n jets ($n = 0, 1, 2$), 4 b-quarks, ZZ, Zbbar
- To improve bkgd rejection over signal: apply b-tagging and Z-tagging vetoes
- CompHEP for signal and CompHEP, PYTHIA, Alpgen for bckg. with CTEQ5L
- Full simulation/reco + L1 + HLT(trigger) cuts
- Theoretical and experimental uncert.

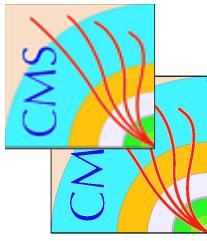
Studied for low lum run $\sim 2 \times 10^{33} \text{cm}^{-2}\text{s}^{-1}$



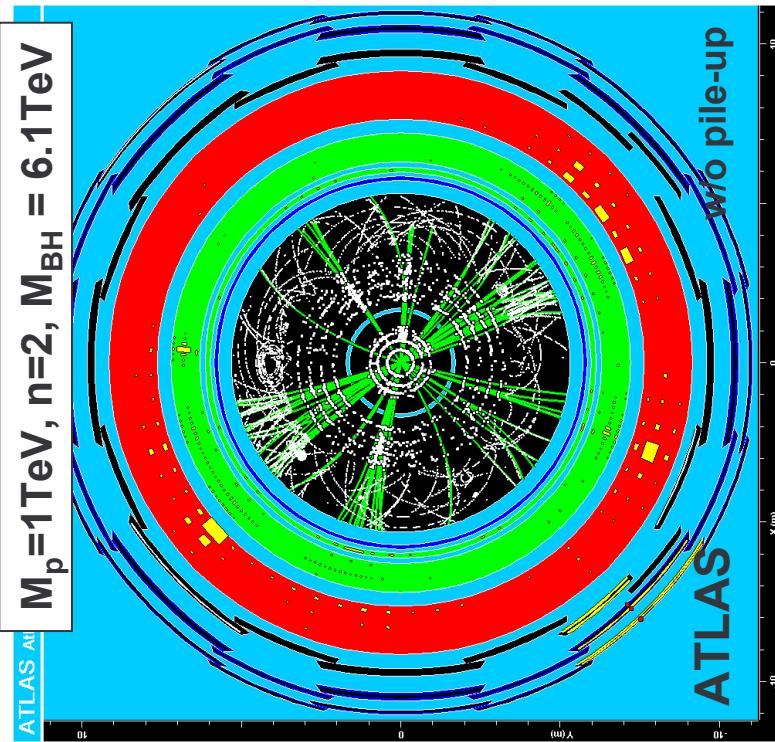


Higgs Black Hole Signatures

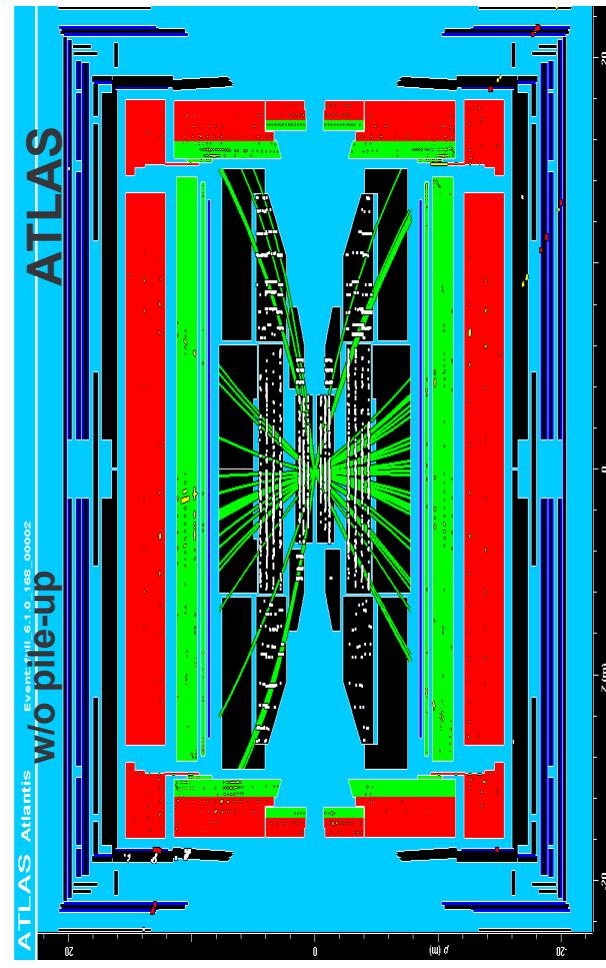
Dimopoulos and Landsberg PRL87 (2001) 161602



- In large ED (ADD) scenario, when impact parameter smaller than Schwartzschild radius Black Hole produced with potentially large x-sec (~ 100 pb).
- Decays democratically through Black Body radiation of SM states – Boltzmann energy distribution.



- Discovery potential (preliminary)
 - $M_p < \sim 4$ TeV $\rightarrow < \sim 1$ day
 - $M_p < \sim 6$ TeV $\rightarrow < \sim 1$ year
- Studies continue ...





T₂KND-g* Discovery Limits

This is more challenging than Z/W which have leptonic decay modes

Detect KK gluon excitations (g^*) by reconstructing their hadronic decays (no leptonic decays).

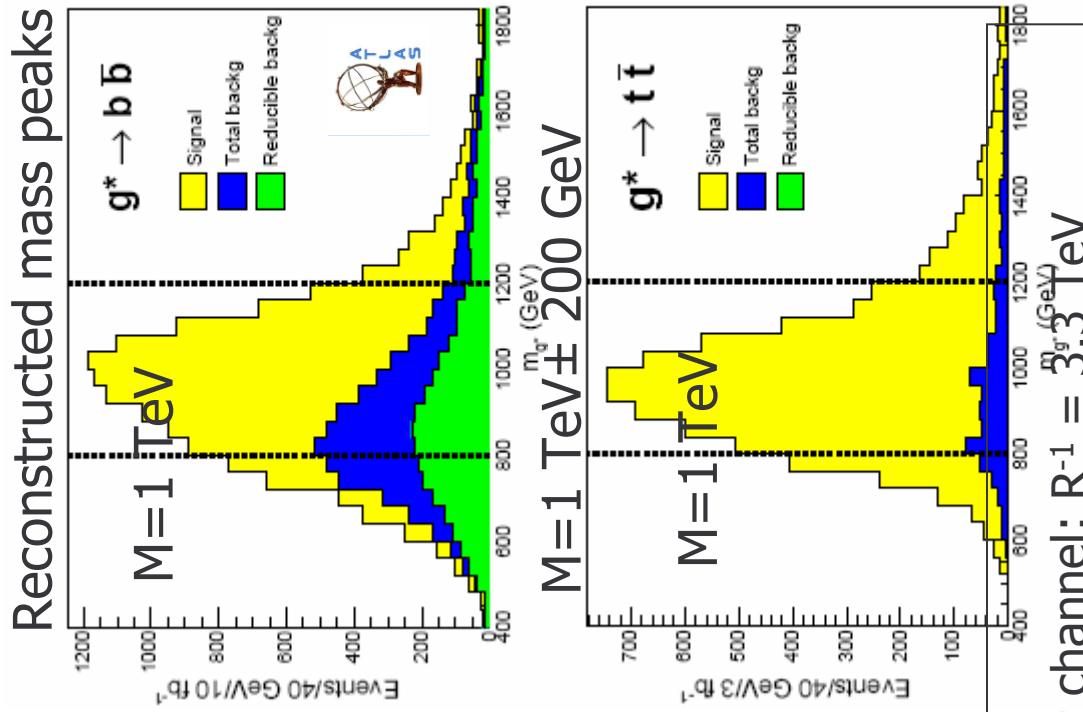
Detect g^* by (1) deviation in dijet σ

(2) analysing its decays into heavy quarks
Gluon excitation decays

$$q\bar{q} \rightarrow g^* \rightarrow b\bar{b}, q\bar{q} \rightarrow g^* \rightarrow t\bar{t}$$

- bbar or ttbar jets
- For ttbar one t is forced to decay leptonically
- Bckg: SM continuum bbar, ttbar, 2 jets, W + jets
- PYTHIA
- Fast simulation/reco

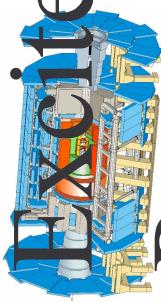
With 300 fb^{-1} Significance of 5 achieved for:
bbar channel: $R^{-1} = 2.7 \text{ TeV}$ ttbar channel: $R^{-1} = 3.3 \text{ TeV}$



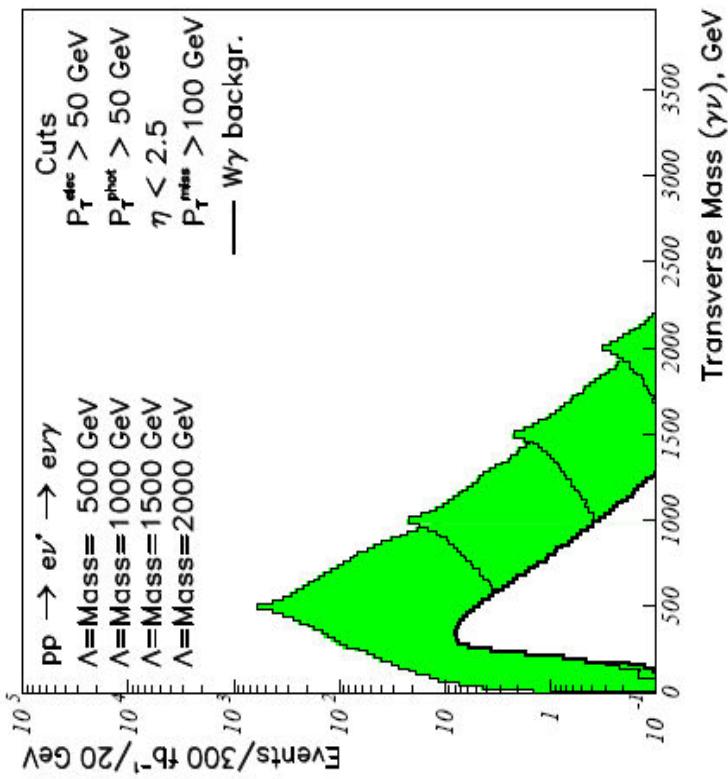
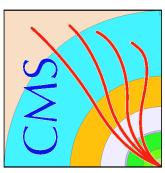
L. March, E. Ros, B. Salvachua,
ATL-PHYS-PUB-2006-002

Tracey Berry (RHUL)

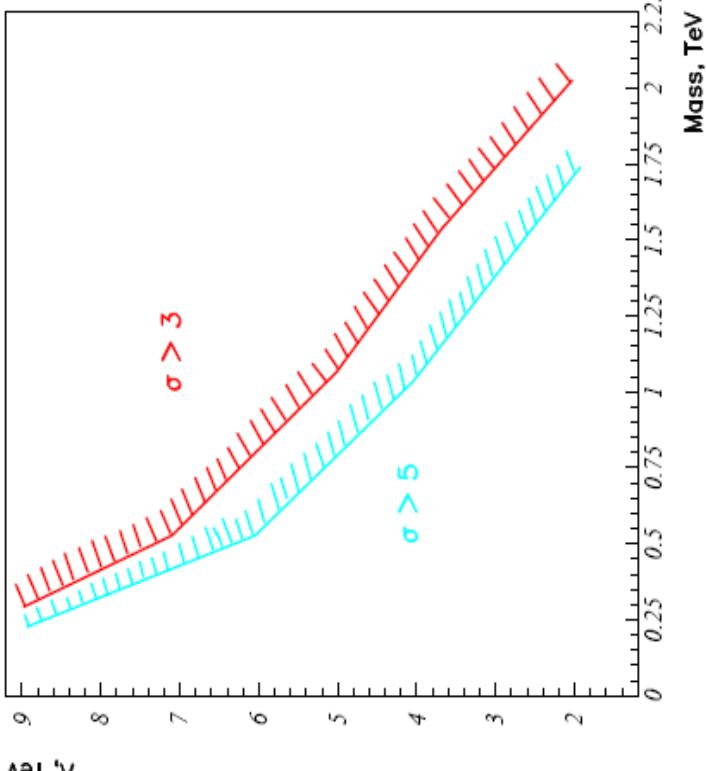
Flavour Workshop, CERN



ExNeD Neutrino Discovery Potential



Λ, TeV

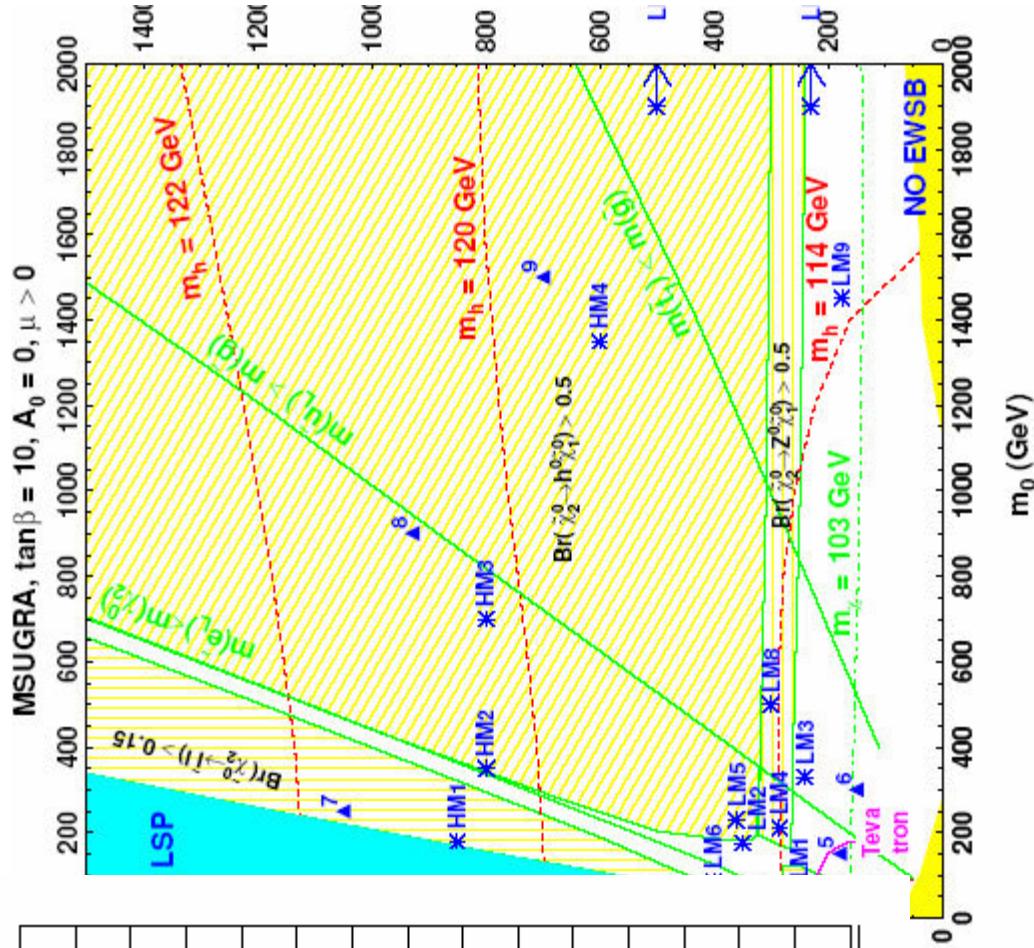


With 300 fb^{-1}

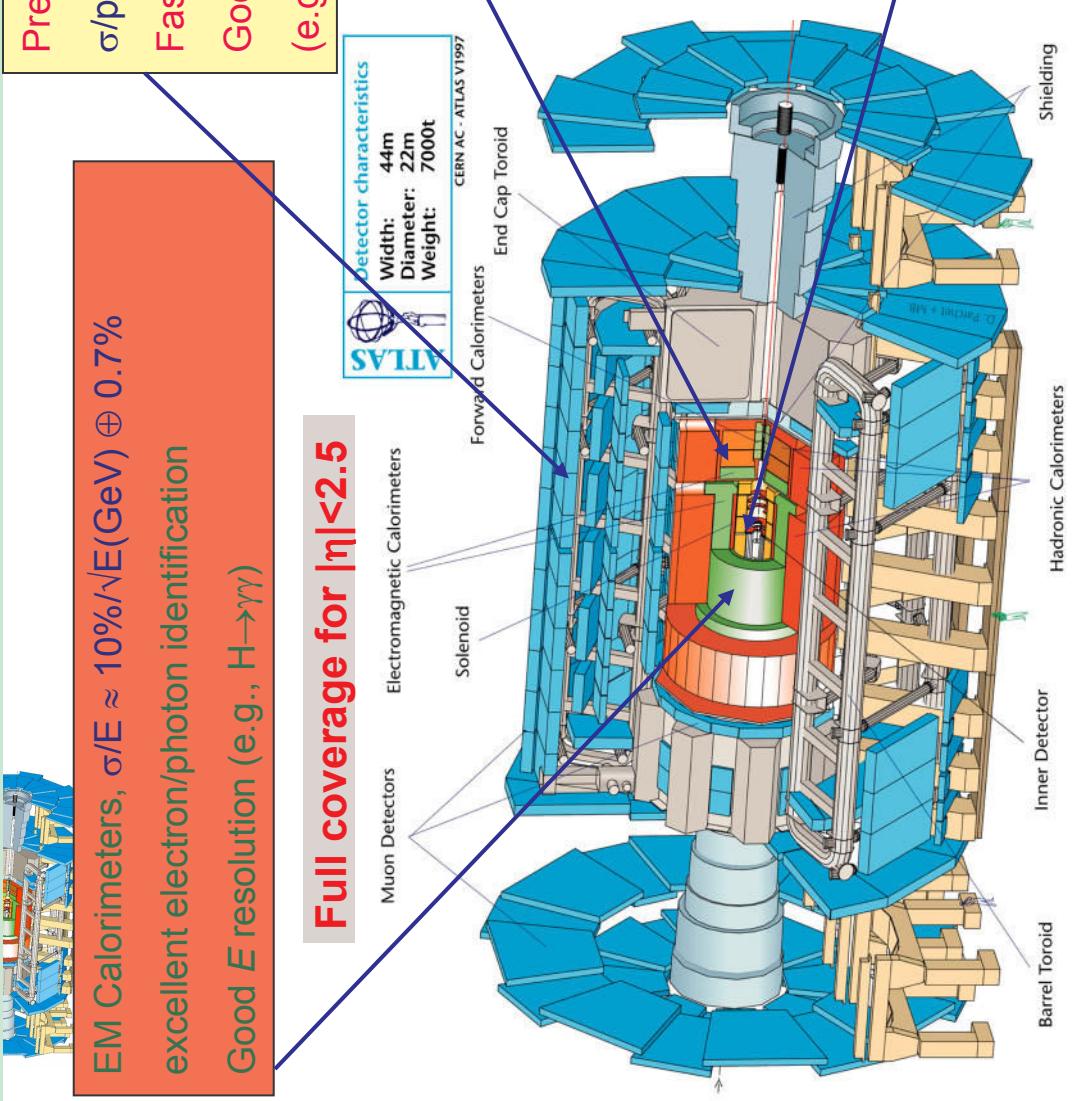


GeV/c^2

Point	m_0	$m_{1/2}$	$\tan \beta$	$\text{sgn}(\mu)$	A_0
LM1	60	250	10	+	0
LM2	185	350	35	+	0
LM3	330	240	20	+	0
LM4	210	285	10	+	0
LM5	230	360	10	+	0
LM6	85	400	10	+	0
LM7	3000	230	10	+	0
LM8	500	300	10	+	-300
LM9	1450	175	50	+	0
LM10	3000	500	10	+	0
HM1	180	850	10	+	0
HM2	350	800	35	+	0
HM3	700	800	10	+	0
HM4	1350	600	10	+	0



A Toroidal LHC ApparatuS (ATLAS) DETECTOR



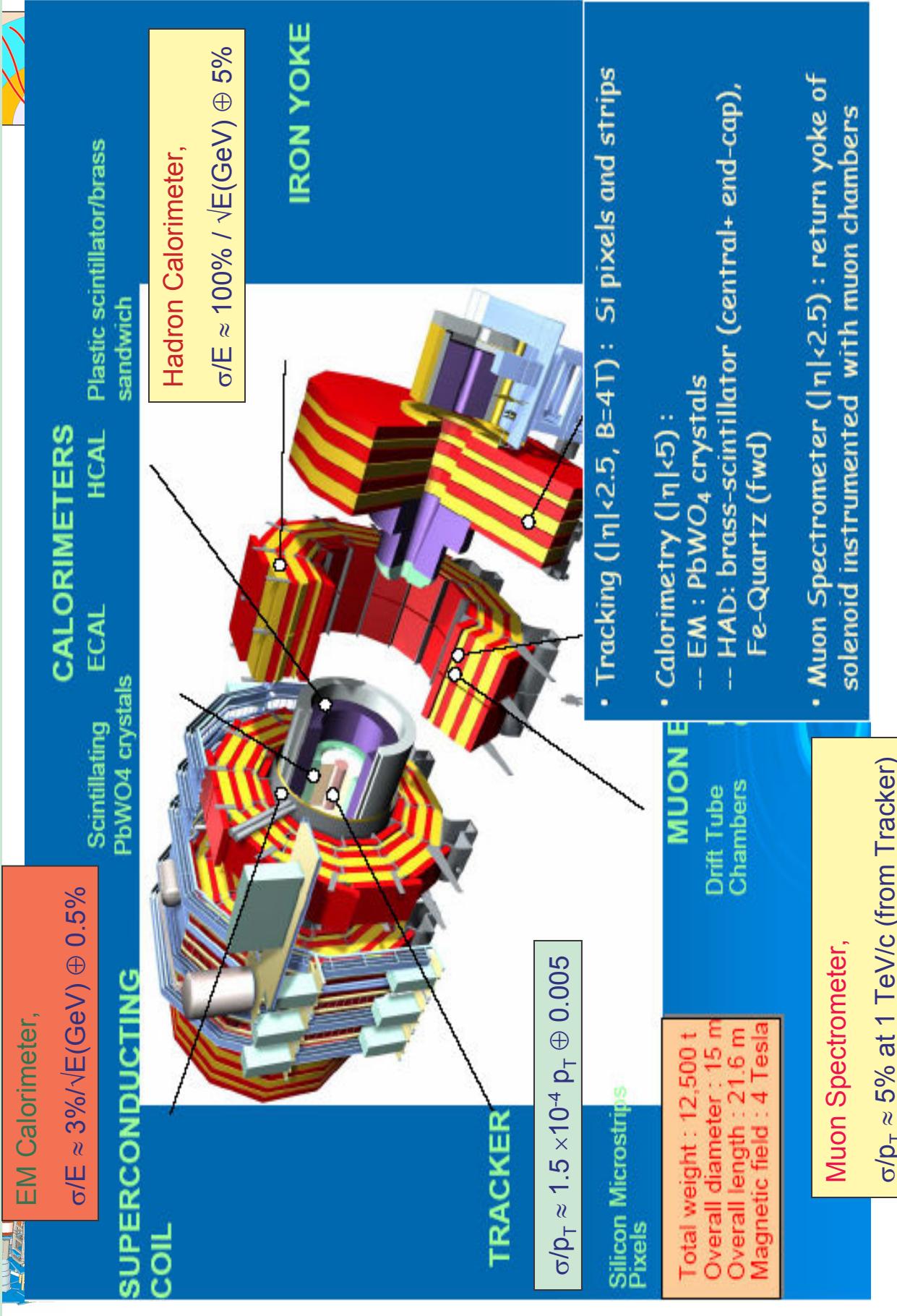
Precision Muon Spectrometer,
 $\sigma/p_T \approx 10\%$ at 1 TeV/c
Fast response for trigger
Good p resolution
(e.g., $A/Z' \rightarrow \mu\mu$, $H \rightarrow 4\mu$)

Hadron Calorimeters,
 $\sigma/E \approx 50\% / \sqrt{E(\text{GeV})} \oplus 3\%$
Good jet and E_T miss performance
(e.g., $H \rightarrow \tau\tau$)

Inner Detector:
Si Pixel and strips (SCT) &
Transition radiation tracker (TRT)
 $\sigma/p_T \approx 5 \times 10^{-4} p_T \oplus 0.001$
Good impact parameter res.
 $\sigma(d_0) = 15\mu\text{m}$ @ 20 GeV (e.g. $H \rightarrow b\bar{b}$)

Magnets: solenoid (Inner Detector) 2T, air-core toroids (Muon Spectrometer) ~0.5T

Compact Muon Solenoid (CMS) DETECTOR



ATLAS Inclusive Trigger Selection Signatures

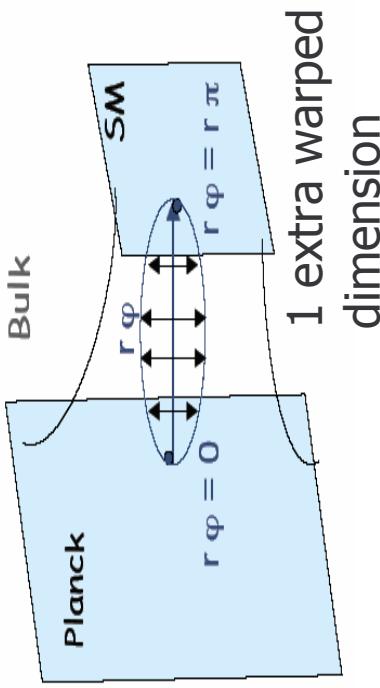
- To select an extremely broad spectrum of “expected” and “unexpected” Physics signals (hopefully!).
- The selection of Physics signals requires the identification of **objects** that can be **distinguished** from the high particle density environment.

Object	Examples of physics coverage	Nomenclature
Electrons	Higgs (SM, MSSM), new gauge bosons, extra dimensions, SUSY, W/Z, top	e25i, 2e15i
Photons	Higgs (SM, MSSM), extra dimensions, SUSY	γ 60i, 2 γ 20i
Muons	Higgs (SM, MSSM), new gauge bosons, extra dimensions, SUSY, W/Z, top	μ 20i, 2 μ 10
Jets	SUSY, compositeness, resonances	j360, 3j150, 4j100
Jet+missing E_T	SUSY, leptoquarks, “large” extra dimensions	j60 + xE60
Tau+missing E_T	Extended Higgs models (e.g. MSSM), SUSY	τ 30 + xE40
also inclusive missingET, SumET, SumET_jet		& many prescaled and mixed triggers

The list must be non-biasing, flexible, include some redundancy, extendable, to account for the “unexpected”.



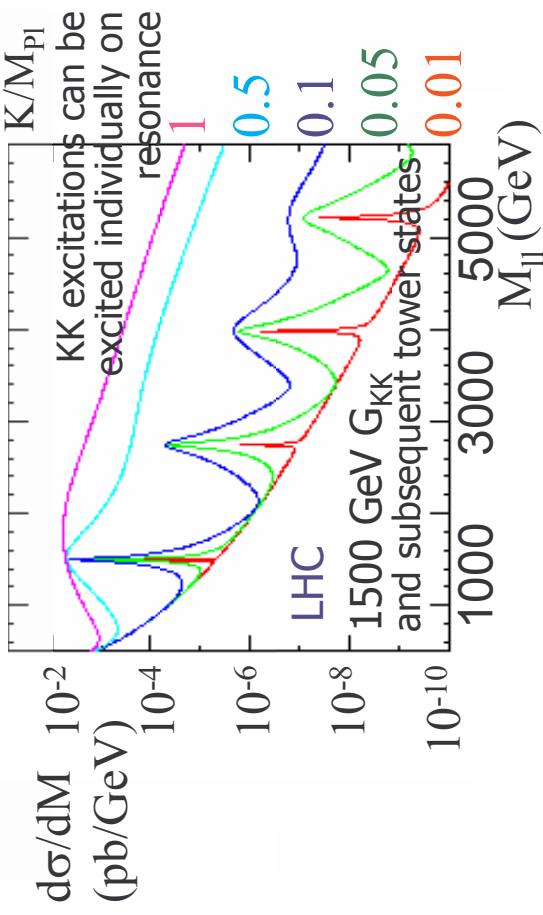
Experimental Signature for RS Model



Signature:

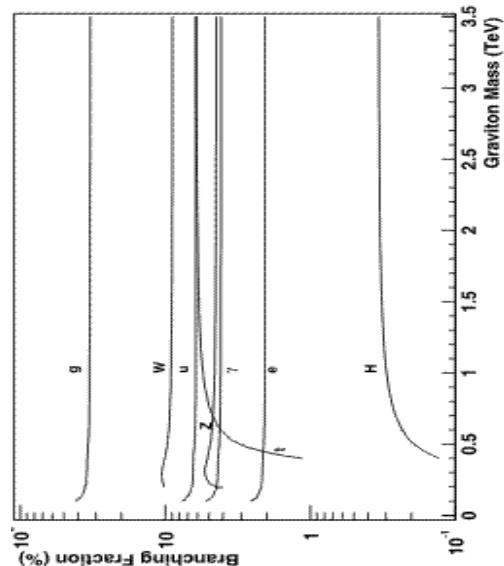
Narrow, high-mass resonance states
in dilepton/dijet/diboson channels

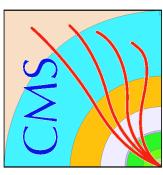
$$q\bar{q}, gg \rightarrow G_{KK} \rightarrow e^+ e^-, \mu^+ \mu^-, \tau^+ \tau^-$$



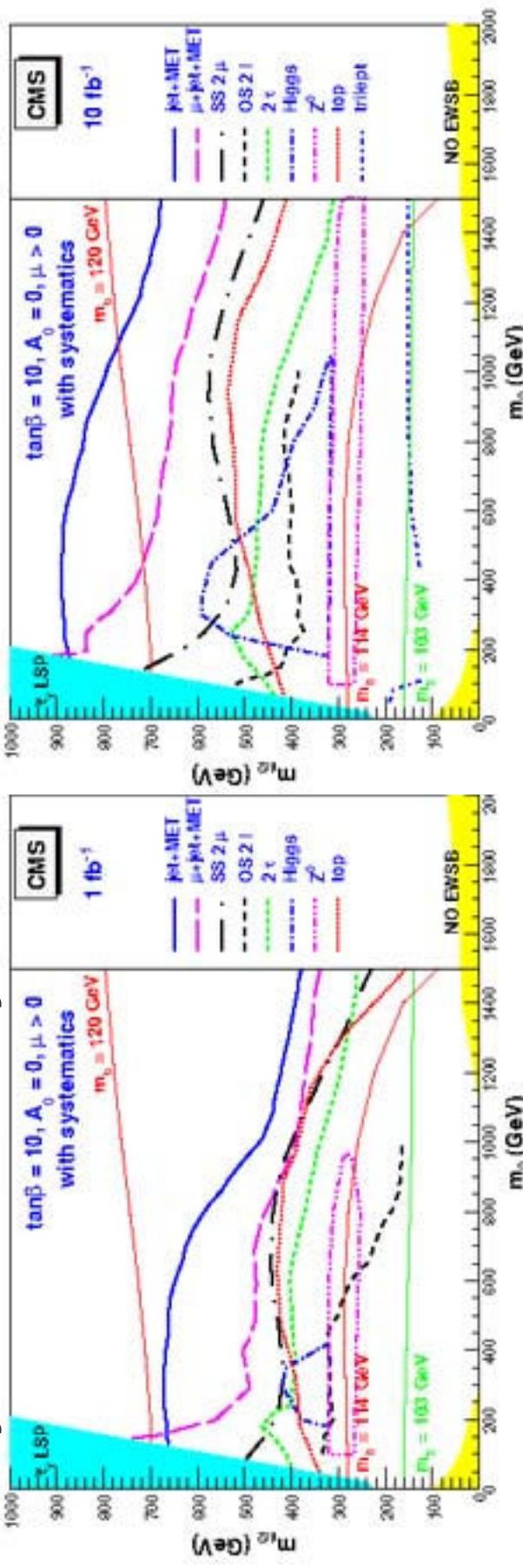
Model parameters:

- Gravity Scale: $\Lambda_\pi = \overline{M}_{pl} e^{-kR_{ct}}$ **Resonance Position**
 $m_1 = m_1 \overline{M}_{pl}/kx_1$, & $m_n = kx_n e^{kR_{ct}} (J_1(x_n) = 0)$
 - Coupling constant: $c = k/M_{Pl}$
 $\Gamma_1 = \rho m_1 x_1^2 (k/M_{pl})^2 \rightarrow \text{width}$
- $k = \text{curvature}, R = \text{compactification radius}$





Best reach obtained with the most inclusive channels:
jets+MET & muons+jet+MET



The range of gluino and squark masses up to about 1.5 TeV can be probed with $L=1\text{fb}^{-1}$ and up to about 2 TeV with 10fb^{-1}

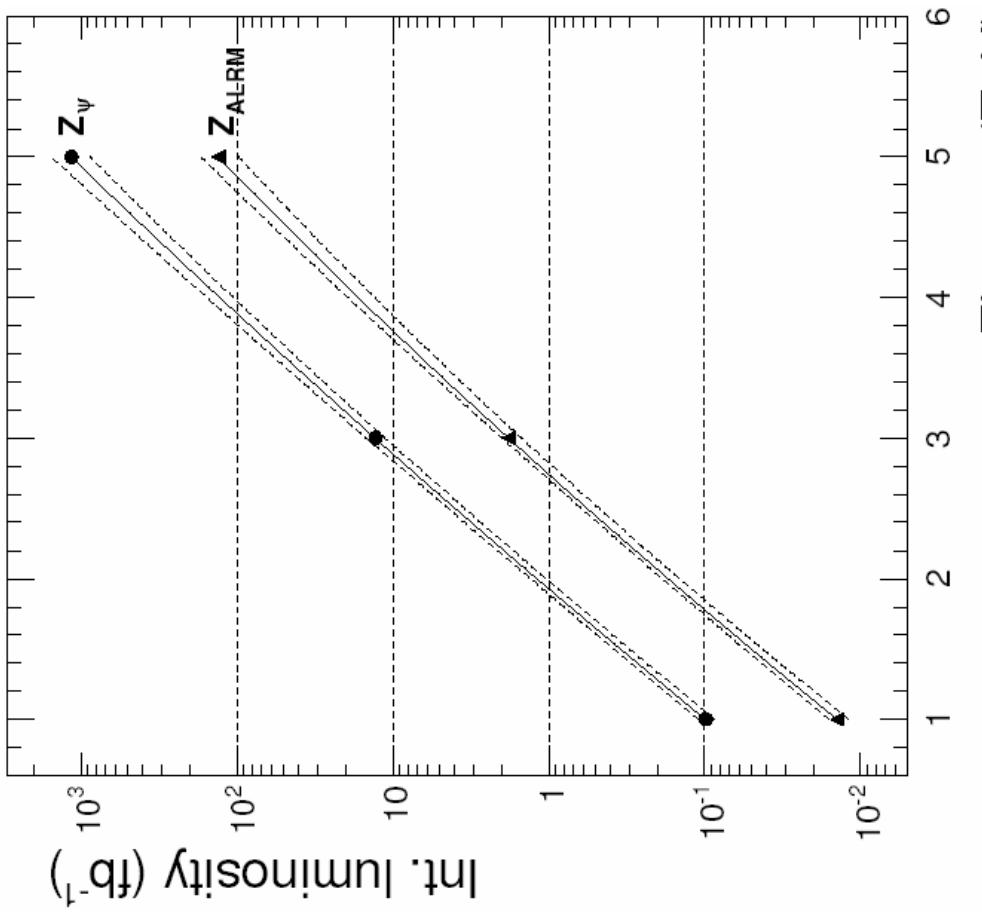
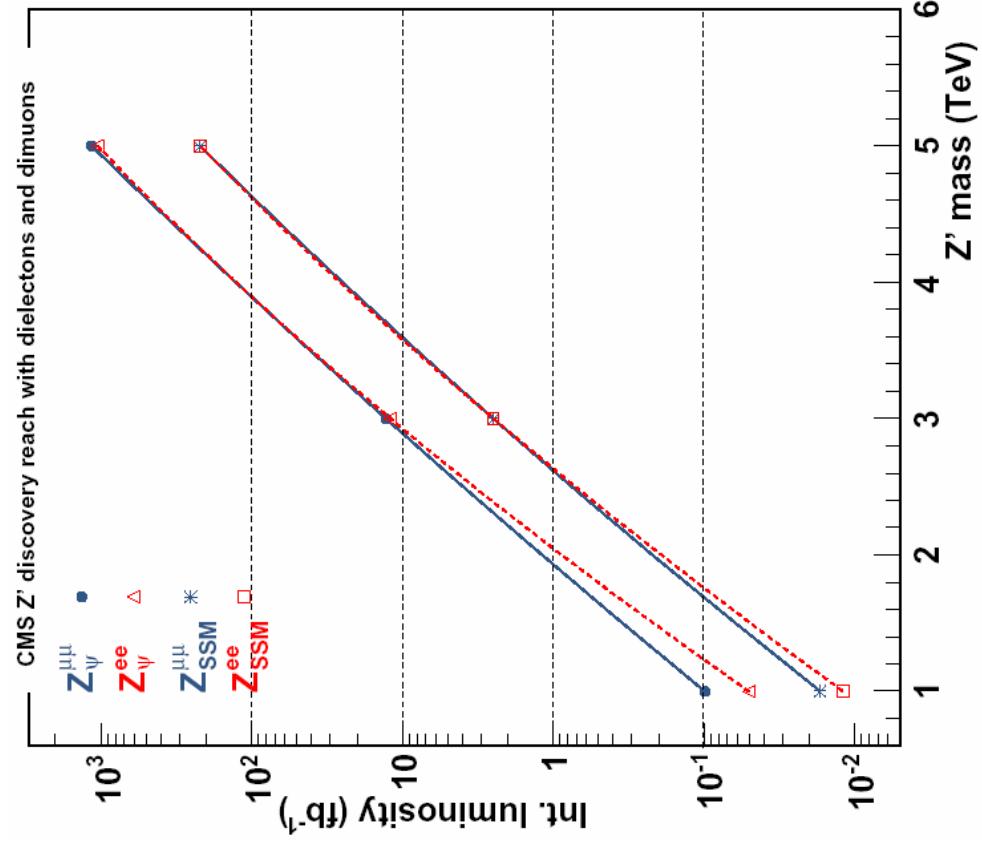
Simultaneous observation of a signal in various topologies would help determine the underlying physics. E.g. triangular dilepton mass distribution.

If discovered plan would be to do more exclusive analyses to reconstruct the sparticle masses and measure cross-sections of sub-processes and their ratios

Heavy Gauge Bosons: Z'

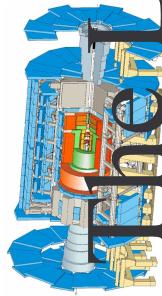


- CMS Z' studies (TDR): integrated luminosity needed for 5σ signal

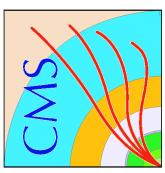


Z' 5σ reach $\text{vs. } \mu\mu$ channels

Z' 5σ reach: impact of theory uncertainties

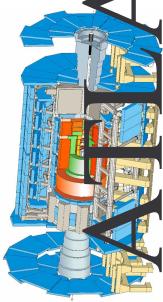


ATLAS Machine Schedule



<i>year</i>	<i>energy</i>	<i>luminosity</i>	<i>physics beam time</i>
2007	450+450 GeV	5x10³⁰	protons - 26 days at 30% overall efficiency 0.7*10⁶ seconds
2008	7+7 TeV	0.5x10³³	protons - starting beginning July 4*10⁶ seconds ions - end of run - 5 days at 50% overall efficiency → 0.2*10⁶ seconds
2009	7+7 TeV	1x10³³	protons: 50% better than 2008 → 6*10⁶ seconds ions: 20 days of beam at 50% efficiency → 10⁶ seconds
2010	7+7 TeV	1x10³⁴	TDR targets: protons: → 10⁷ seconds ions: → 2*10⁶ seconds

See many other talks at this meeting for ATLAS experimental details



ATLAS detector

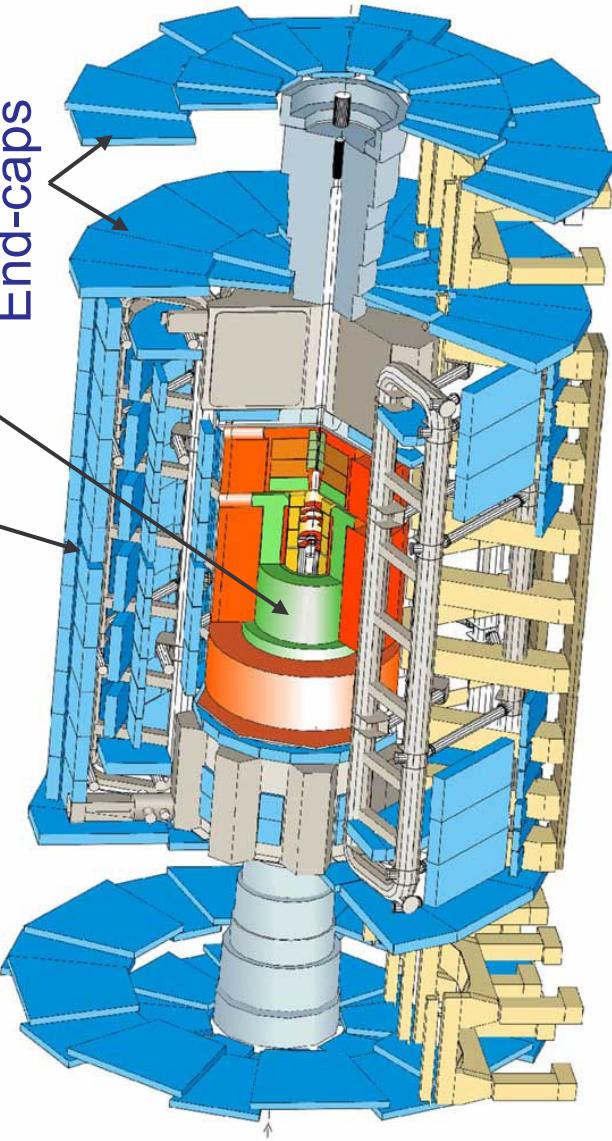
- High energy electrons are detected by LAr calorimeter.
- Muons are detected by the Muon System.
- Expected electron energy resolution is:
 - $\sim 0.6\%$ for $E=500\text{GeV}$,
 - $\sim 0.5\%$ for $E=1000\text{GeV}$.
- Muon transverse momentum (p_T) resolution is:
 - $\sim 6\%$ for $p_T=500\text{GeV}$,
 - $\sim 11\%$ for $p_T=1000\text{GeV}$.

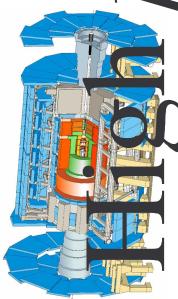
Electron energy resolution

$$\frac{\sigma(E)}{E} = \frac{9.5\%}{\sqrt{E}} \oplus 0.45\%$$

Muon System

LAr Calorimeter





High p_T leptons from Z' decay

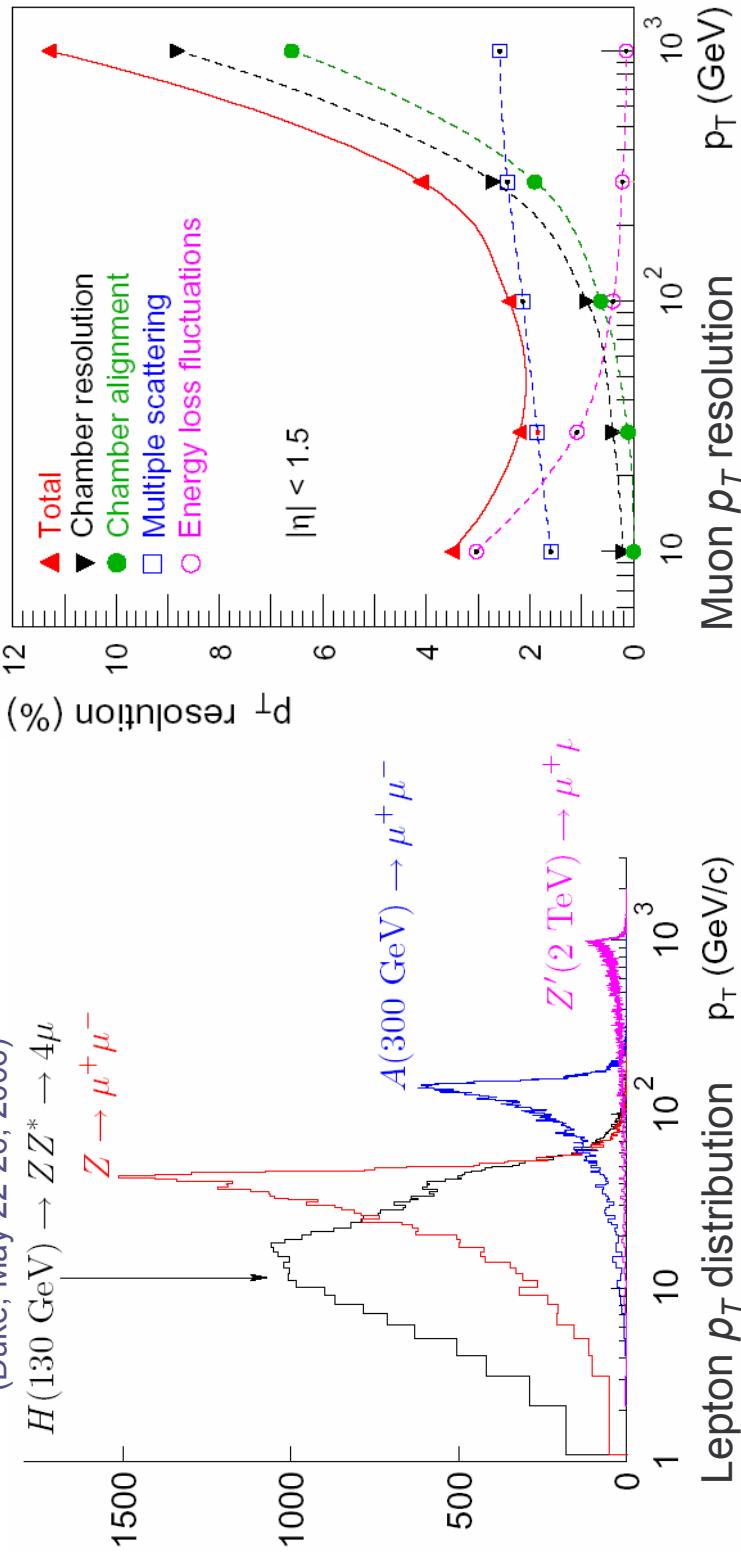
- The leptons p_T distribution from Z' decay has a Jacobian peak.
- At high p_T , the muon momentum resolution degrades.
- For the muon p_T resolution, calibration and alignment are critical.

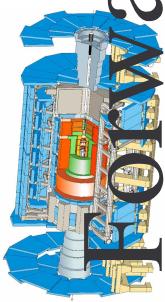
Oliver Kortner (MPI), HCP2006
(Duke, May 22-26, 2006)

$$H(130 \text{ GeV}) \rightarrow ZZ^* \rightarrow 4\mu$$

$$Z \rightarrow \mu^+ \mu^-$$

Muon spectrometer TDR
(CERN/LHCC 97-22)





Foward-backward asymmetry

- As a probe of the underlying model, one can measure the forward-backward asymmetry.
- The differential cross section of Z' depends on $\cos\theta^*$.
- And if Z' has spin 1, the differential cross section is given by:

- $$\frac{d\sigma}{A_{FB}(M_{l\nu}) d\cos\theta^*} \propto \frac{3}{8}(1 + \cos^2\theta^*) + A_{FB} \cos\theta^* \quad \text{if } \theta^* \text{ is angle between } l^- \text{ and quark in the CMS of the colliding partons.}$$
- This quantity $A_{FB}(M_{l\nu}) = \frac{N_+ - N_-}{N_+ + N_-}$ indent.
 N_+ : number of events with the lepton in the forward
 N_- : number of events with the lepton in the backward
- One can discriminate between the underlying models by measuring $A_{FB}(M_{l\nu})$.

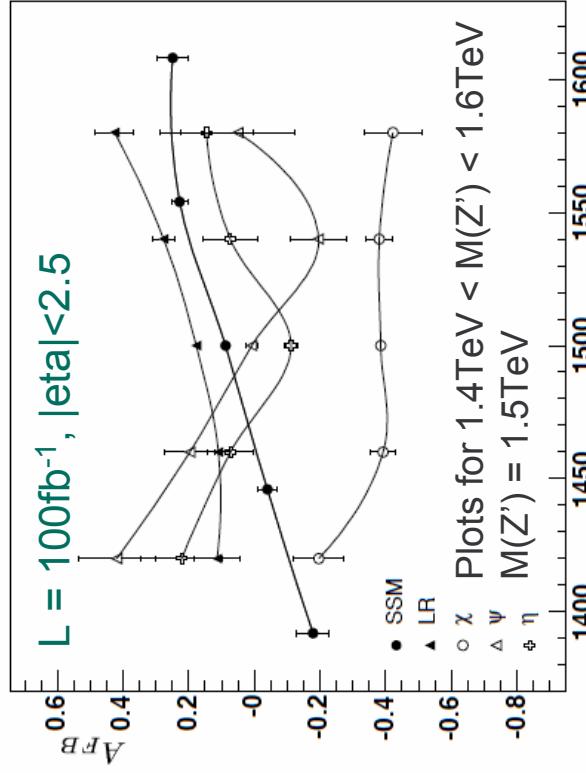
Ref: ATL-PHYS-PUB-2005-010



$\Delta TB(M_{l\nu})$ measurement(1)

▪ $Z' \rightarrow e^+ e^-$:

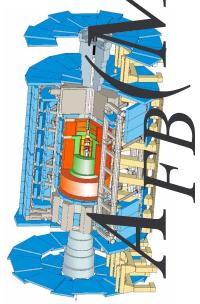
- high discriminating power of the asymmetry.



Asymmetry at generation level

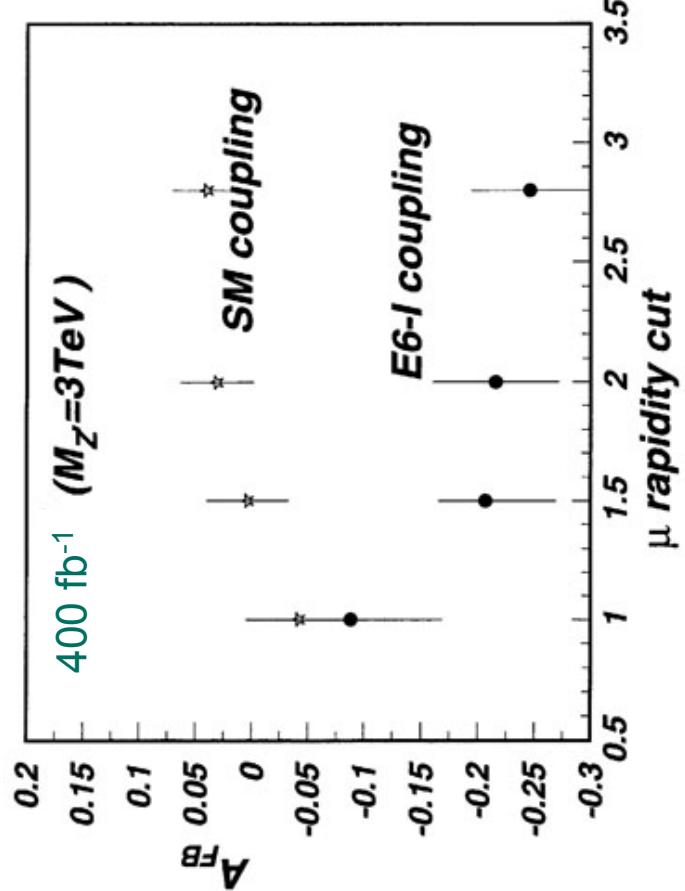
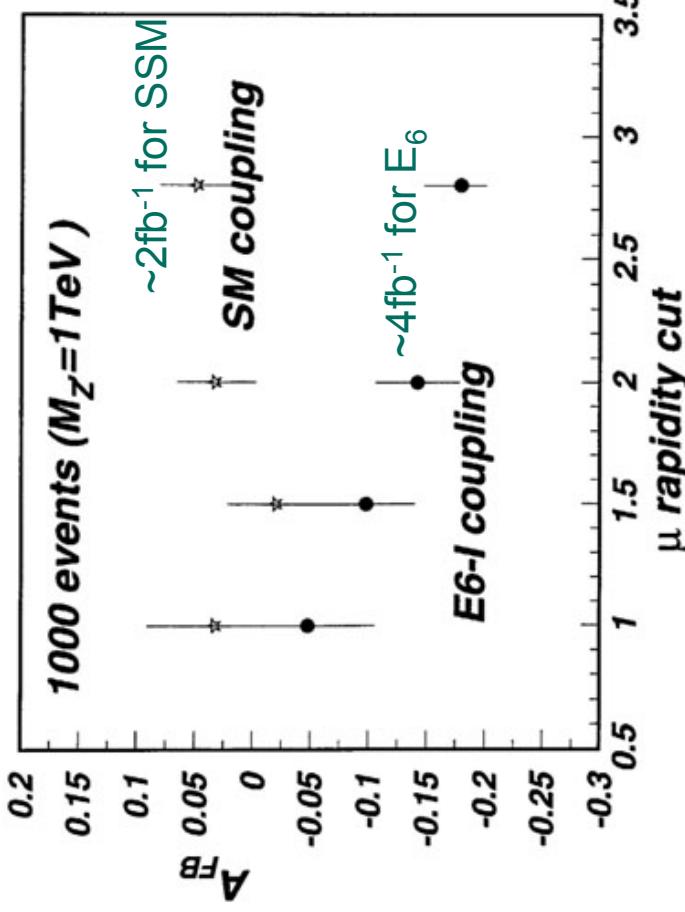
Model	$\int \mathcal{L}(fb^{-1})$	Generation	Observed	Corrected
1.5 TeV				
SSM	100	$+0.088 \pm 0.013$	$+0.060 \pm 0.022$	$+0.108 \pm 0.027$
χ	100	-0.386 ± 0.013	-0.144 ± 0.025	-0.361 ± 0.030
η	100	-0.112 ± 0.019	-0.067 ± 0.032	-0.204 ± 0.039
η	300	-0.090 ± 0.011	-0.050 ± 0.018	-0.120 ± 0.022
ψ	100	$+0.008 \pm 0.020$	-0.056 ± 0.033	-0.079 ± 0.042
ψ	300	$+0.010 \pm 0.011$	-0.019 ± 0.019	-0.011 ± 0.024
LR	100	$+0.177 \pm 0.016$	$+0.100 \pm 0.026$	$+0.186 \pm 0.032$
4 TeV				
SSM	500	$+0.138 \pm 0.099$	$+0.006 \pm 0.183$	$+0.265 \pm 0.260$
KK	500	$+0.491 \pm 0.028$	$+0.189 \pm 0.057$	$+0.457 \pm 0.073$

Ref: ATL-PHYS-PUB-2005-010

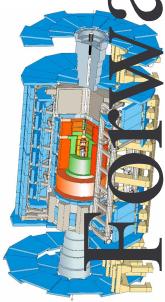


$M_{FB}(M_{l\bar{l}})$ measurement(2)

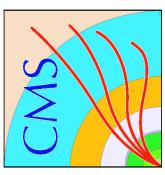
- $Z' \rightarrow \mu^+ \mu^-$
 - With 200fb^{-1} , the ATLAS can distinguishes the underlying theories with accuracy better than 3% using the asymmetry for $M(Z')$ less than 2TeV.
 - At higher masses, we need much more luminosity.



Ref: ATLAS Internal Note Muon-NO-161 23 May 1997



Forward and backward



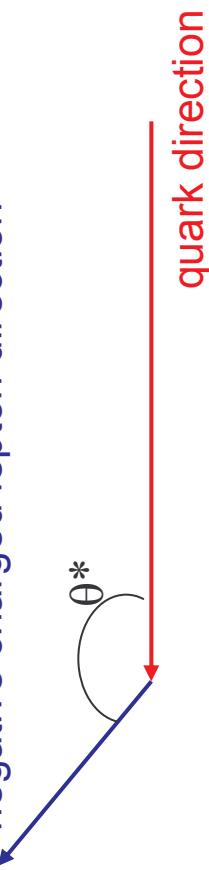
Forward

$$\cos \theta^* > 0$$



Backward

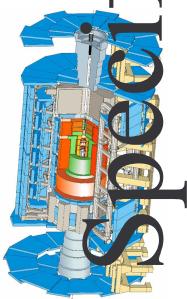
$$\cos \theta^* < 0$$



When $\cos \theta^*$ is positive, we call forward, and when $\cos \theta^*$ is negative we call backward.

The quark direction is not directly accessible in the data.

Therefore the Z' momentum defines the quark direction, because of the quark generally being at a higher momentum than the antiquark



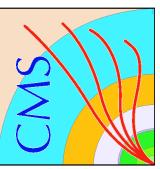
Specific models

- A popular model is:
 - Effective $SU(2) \times U(1)_Y \times U(1)'_\gamma$
 - There are two additional neutral gauge bosons.
 - The new gauge boson uniquely determined by:

$$Z' = \cos \theta Z'_\psi - \sin \theta Z'_\chi \quad \theta \text{ is a new mixing angle.}$$

- There are 3 special cases:
 - Z'_ψ model: $\theta=0, E_6 \rightarrow SO(10) \times U(1)_\psi$
 - Z'_χ model: $\theta=-\pi/2, E_6 \rightarrow SO(10) \times U(1)_\psi \rightarrow SU(5) \times U(1)_\chi \times U(1)_\psi$
 - Z'_η model: $\theta=\arctan(-\sqrt{5/3})+\pi/2, E_6 \rightarrow SU(3)_C \times SU(2)_L \times U(1)_Y \times U(1)_\eta = SM \times U(1)_\eta$
(E_6 breaks directly down to a rank 5 model.)
- Other popular models:
 - The Left-Right model from the breaking of the $SO(10)$ group,
 - The Kaluza-Klein model (Extra Dimension).
 - etc...

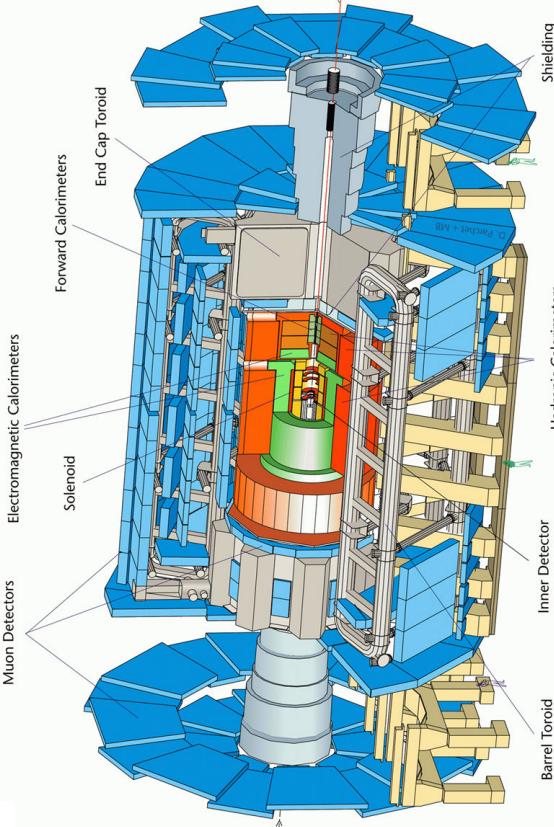
Ref: ATL-PHYS-PUB-2005-010



WAIS and CMS Experiments

Large general-purpose particle physics detectors

A Large Toroidal LHC Apparatus



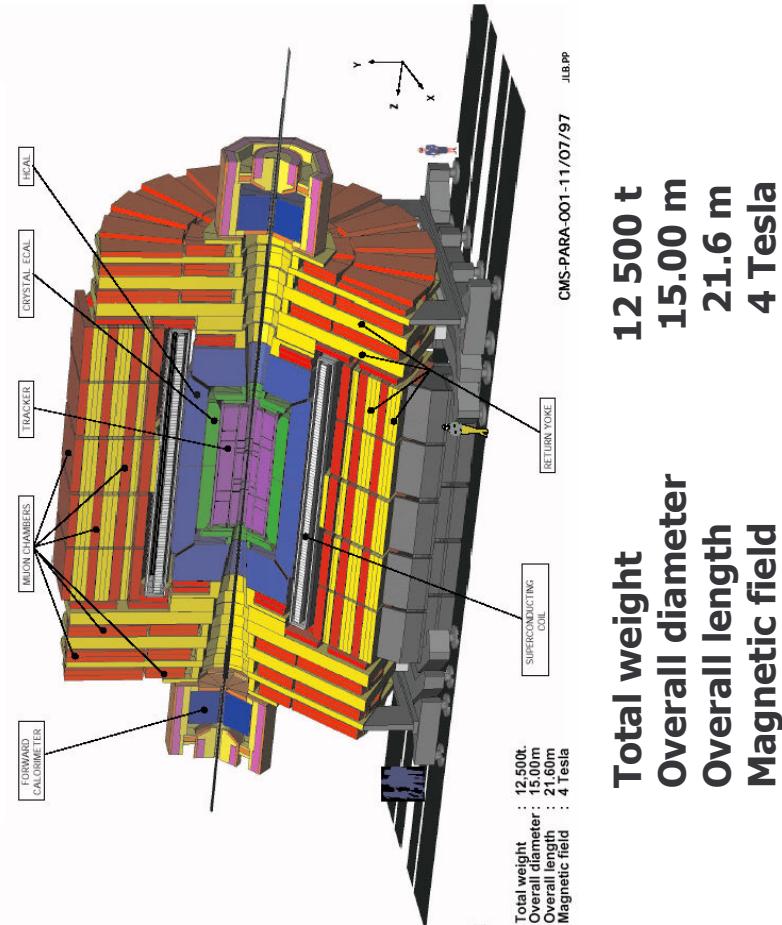
Total weight
Overall diameter
Barrel toroid length
End-cap end-wall chamber span
Magnetic field

7000 t
25 m
26 m
46 m
2 Tesla

Total weight:
12.500t.
Overall diameter:
15.00m
Overall length:
21.60m
Magnetic field:
4 Tesla

CMS PARA-001-11/07/97 J.B.PP

Compact Muon Solenoid



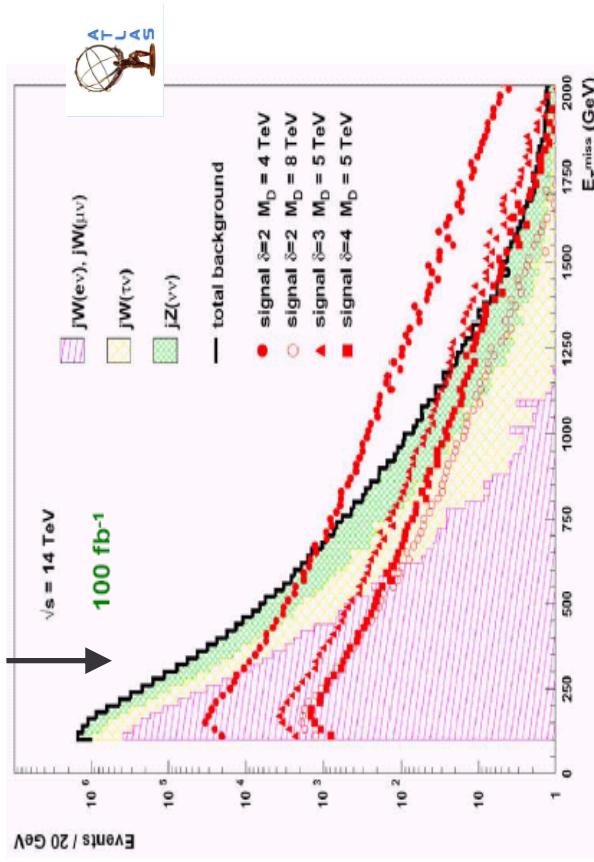
Detector subsystems are designed to measure:
energy and momentum of $\gamma, e, \mu, \text{jets}$, missing E_T up to a few TeV



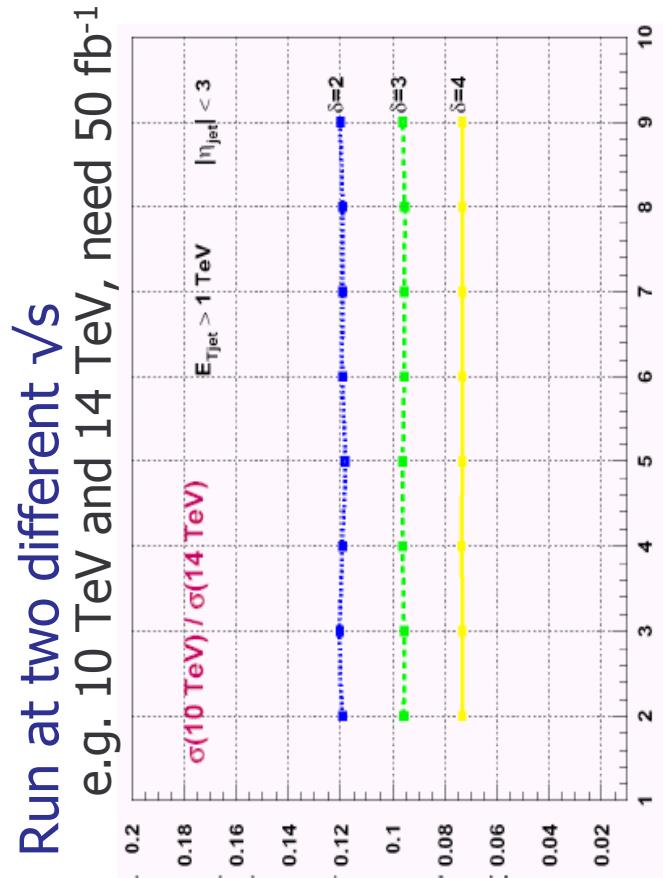
ADD Parameters: jet+G Emission

To characterise the model need to measure M_D and δ

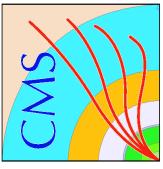
Measuring $\sigma(pp \rightarrow \text{jet}+\text{GKK})$ gives ambiguous results



Use variation of σ on \sqrt{s}
 σ at different \sqrt{s} almost
independent of M_D , varies with δ

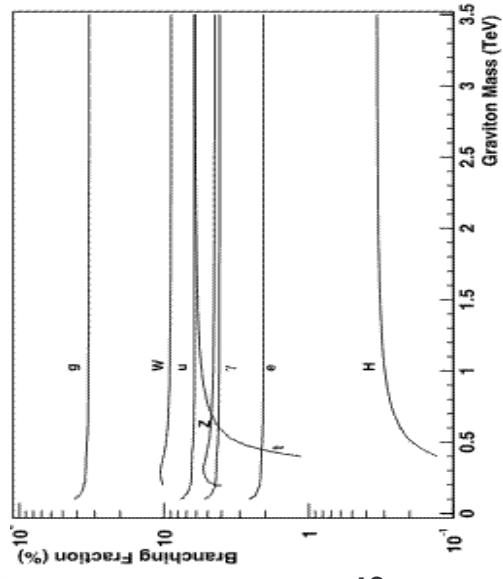


Rates at 14 TeV of $\delta=2$ $M_D=6$ TeV very
similar to $\delta=3$ $M_D=5$ TeV whereas
Rates at 10 TeV of ($\delta=2$ $M_D=6$ TeV) and
($\delta=3$ $M_D=5$ TeV) differ by \sim factor of 2



A resonance could be seen in many other channels: $\mu\mu, \gamma\gamma, jj, b\bar{b}b\bar{b}$, $t\bar{t}b\bar{b}$, WW, ZZ , hence allowing to check universality of its couplings:

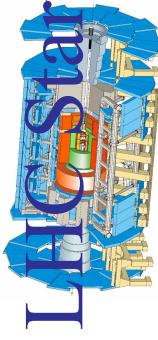
Channel	Point m_G, Λ_π (TeV)					
	1,10	1,20	1,30	2,10	2,20	2,30
e^+e^-	1.6	3.3	5.3	5.4	11.0	17.1
$\mu^+\mu^-$	1.9	4.5	8.2	6.2	15.2	28.2
$\gamma\gamma$	1.2	2.9	5.2	3.9	8.8	15.2
WW	11.6	44.9	-	38.2	-	-
ZZ	13.7	50.1	-	52.7	-	-
jj	19.0	77.0	-	31.0	-	-



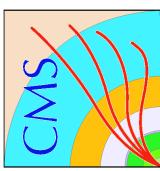
Relative precision achievable (in %) for measurements of σB in each channel for fixed points in the $M_{G'}\Lambda_\pi$ plane. Points with errors above 100% are not shown.

Also the size (R) of the ED could also be estimated from mass and cross-section measurements.

Allenach et al, hep-ph/0211205
Allenach et al, JHEP 9 19 (2000), JHEP 0212 39 (2002)

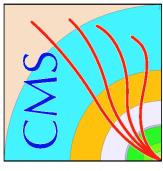


LHC Start-up Expectations



Model	Mass reach	Integrated luminosity (fb ⁻¹)	Systematic uncertainties
ADD Direct G_{KK}	$M_D \sim 1.5\text{-}1.0 \text{ TeV}, n = 3\text{-}6$	1	Theor.
ADD Virtual G_{KK}	$M_D \sim 4.3 - 3 \text{ TeV}, n = 3\text{-}6$ $M_D \sim 5 - 4 \text{ TeV}, n = 3\text{-}6$	0.1 1	Theor.+Exp.
RS1	di-electrons di-photons di-muons di-jets	$M_{G_1} \sim 1.35\text{-} 3.3 \text{ TeV}, c=0.01\text{-}0.1$ $M_{G_1} \sim 1.31\text{-} 3.47 \text{ TeV}, c=0.01\text{-}0.1$ $M_{G_1} \sim 0.8\text{-} 2.3 \text{ TeV}, c=0.01\text{-}0.1$ $M_{G_1} \sim 0.7\text{-} 0.8 \text{ TeV}, c=0.1$	10 10 1 0.1
Tev⁻¹ ($Z_{KK}^{(1)}$)	$M_{z_1} < 5 \text{ TeV}$	1	Theor.

Spin-1/Spin-2 Discrimination



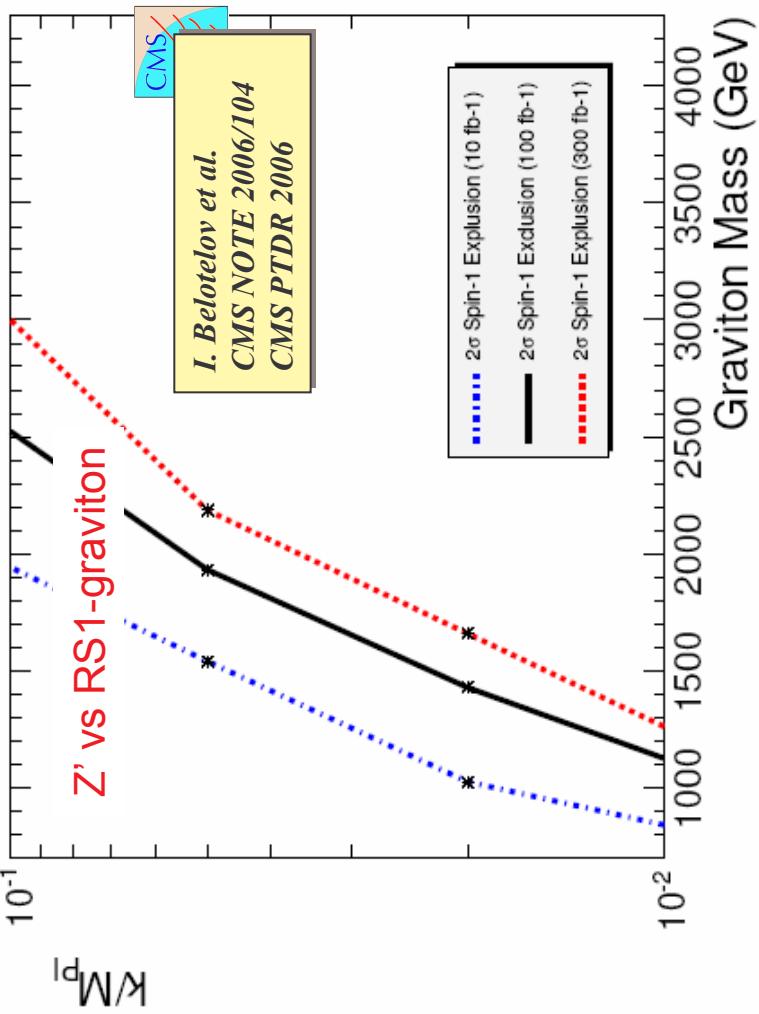
Spin-1 States: Z' from extended gauge models, Z_{KK}

Spin-2 States: RS1-graviton

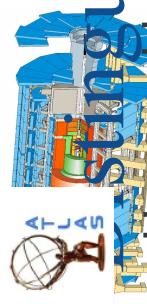
Method: unbinned likelihood ratio statistics incorporating the angles in of the decay products the Collins-Soper farme ([R.Cousins et al. JHEP11 \(2005\) 046](#)).
The statistical technique has been applied to fully simu/reco events.

Angular distributions

- $qq \rightarrow G \rightarrow ff: 1 - 3 \cos^2 \theta + 4 \cos^4 \theta$
- $gg \rightarrow G \rightarrow ff: 1 - \cos^4 \theta$
- $qq \rightarrow G \rightarrow VV: 1 - \cos^4 \theta$
- $gg \rightarrow G \rightarrow VV: 1 + 6 \cos^2 \theta + \cos^4 \theta$
- DY background: $1 + \cos^2 \theta$

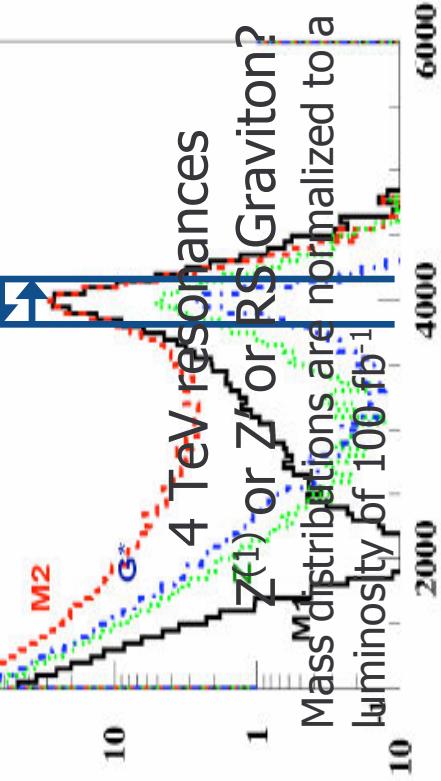


Older results on spin discrimination from ATLAS can be found [B.C. Allanach et al, JHEP 09 \(2000\) 019; ATL-PHYS-2000-029](#)



Distinguishing $Z^{(1)}$ from Z' , RS G

How could a 4 TeV $Z^{(1)}/\gamma^{(1)}$ resonance be distinguished from a 4 TeV Z' or Randall-Sundrum Model Graviton?



Select events around the peak of the resonance $3750 \text{ GeV} < M_{ee} < 4250 \text{ GeV}$

Plot cosine of the angle of the lepton, w.r.t the beam direction, the frame of the decaying resonance.

$Z^{(1)}$, model 2

$\cos(\Theta')$

$q q \rightarrow G$

$\cos(\Theta')$

$g g \rightarrow G^*$

Angular distributions are normalized to 116 events, the number predicted with a luminosity of 100 fb^{-1} for the $Z^{(1)}/\gamma^{(1)}$ case

(+ve direction was defined by the sign of reconstructed momentum in the dilepton system.)