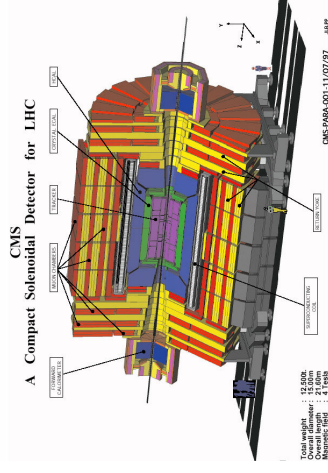


Extra Dimensions



Dr Tracey Berry
Royal Holloway
University of London

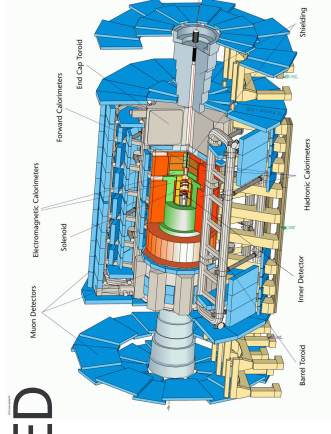
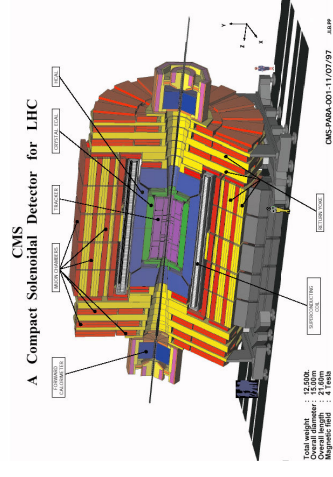
FEI

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Oct 23-27, 2006

LHC & Extra Dimensions



- Theoretical Motivations
- Extra Dimensional Models Considered
- Signatures Covered
- Search Facilities: ATLAS & CMS
- Uncertainties
- Present Constraints and Discovery Limits for ED (ADD, RS, TeV^{-1} , UED)
- Summary of LHC Start-up Expectations
- Conclusions



Extra Dimensions: Motivations



In the late 90's Large Extra Dimensions (LED) were proposed as a solution to the hierarchy problem $M_{EW} (1 \text{ TeV}) \ll M_{Planck} (10^{19} \text{ GeV})?$

ADD

Arkani-Hamed, Dimopoulos, Dvali,
Phys Lett B429 (98)

Many (δ) large compactified EDs
In which G can propagate

$$M_{Pl}^2 \sim R^\delta M_{Pl(4+\delta)}^{(2+\delta)}$$

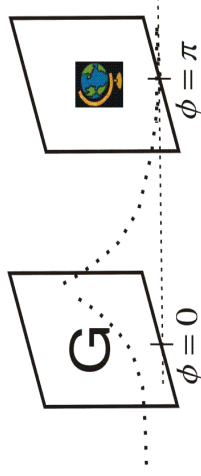
Effective $M_{Pl} \sim 1 \text{ TeV} \rightarrow$ if compact space (R^δ) is large

RS

Randall, Sundrum,
Phys Rev Lett 83 (99)

1 highly curved ED
Gravity localised in the ED

Planck TeV brane



$$\Lambda_\pi = M_{Pl} e^{-kR_c \pi}$$

$$\Lambda_\pi \sim \text{TeV}$$

if warp factor $kR_c \sim 11-12$

Since then, new Extra Dimensional models have been developed and been used to solved other problems:

Dark Matter, Dark Energy, SUSY Breaking, etc
Some of these models can be/have been experimentally tested at high energy colliders

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Extra Dimensional Models



ADD

Arkani-Hamed, Dimopoulos, Dvali,
Phys Lett B429 (98)

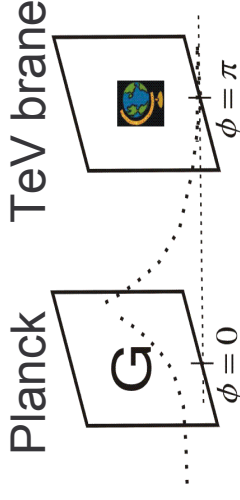
(Many) Large flat Extra-Dimensions (LED) could be as large as a few μm
In which G can propagate, SM particles restricted to 3D brane



RS

Randall, Sundrum,
Phys Rev Lett 83 (99)b

Small highly curved extra spatial dimension
(RS1 – two branes) Gravity localised in the ED



TeV⁻¹
sized EDs

Dienes, Dudas, Gherghetta,
Nucl Phys B537 (99)

Bosons could also propagate in the bulk
Fermions are localized at the same (opposite) orbifold point: destructive
(constructive) interference between SM gauge bosons and KK excitations

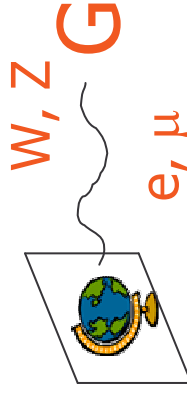
SM chiral fermions

SM Gauge Bosons

W, Z, γ, g

UED

All SM particles propagate in “Universal” ED
often embedded in large ED

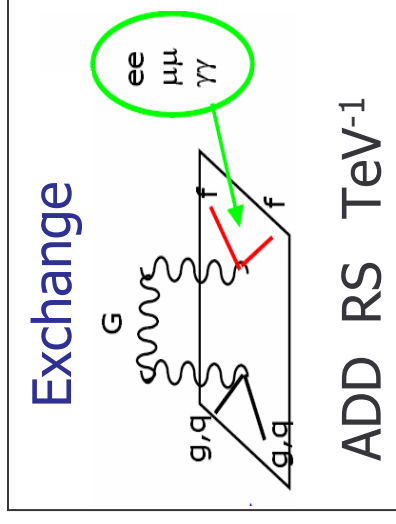
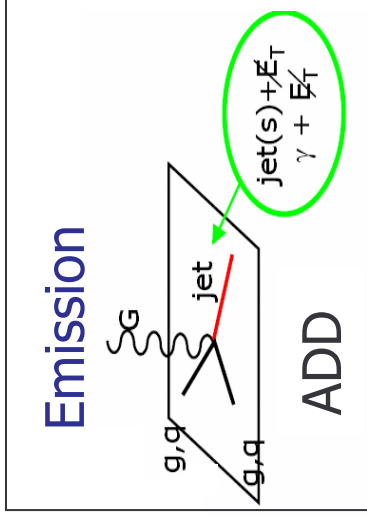


Experimental Signatures of ED



Covered in this talk

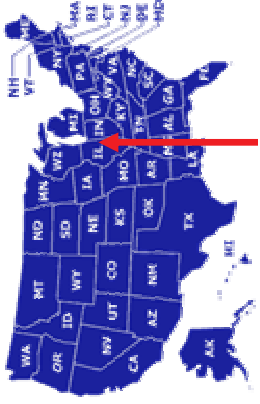
- **Single jets/Single photons + missing E_T**
(direct graviton production in ADD)
- **Di-lepton, di-jet continuum modifications**
(virtual graviton production in ADD)
- **Di-lepton, di-jet and di-photon resonances**
(new particles) in RS1-model (RS1-graviton) and TeV⁻¹ ED model (Z^{KK})
- **Single leptons + missing E_T**
in TeV⁻¹ ED model (W^{KK})
- **Back-to-back energetic jets + missing E_T (UED)**
- **4 jets + 4 leptons + missing E_T (mUED)**





Present/Past ED Search Facilities

Tevatron, Fermilab, USA



Tevatron: Highest energy collider operating in the world!

Run I $\sqrt{s} = 1.8 \text{ TeV}$

Run II $\sqrt{s} = 1.96 \text{ TeV}$

LEP, CERN, Geneva

CERN: world's largest particle physics laboratory



LEP I $\sqrt{s} = 91 \text{ GeV}$

LEP II $\sqrt{s} = 136\text{-}208 \text{ GeV}$

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Future ED Search Facilities!



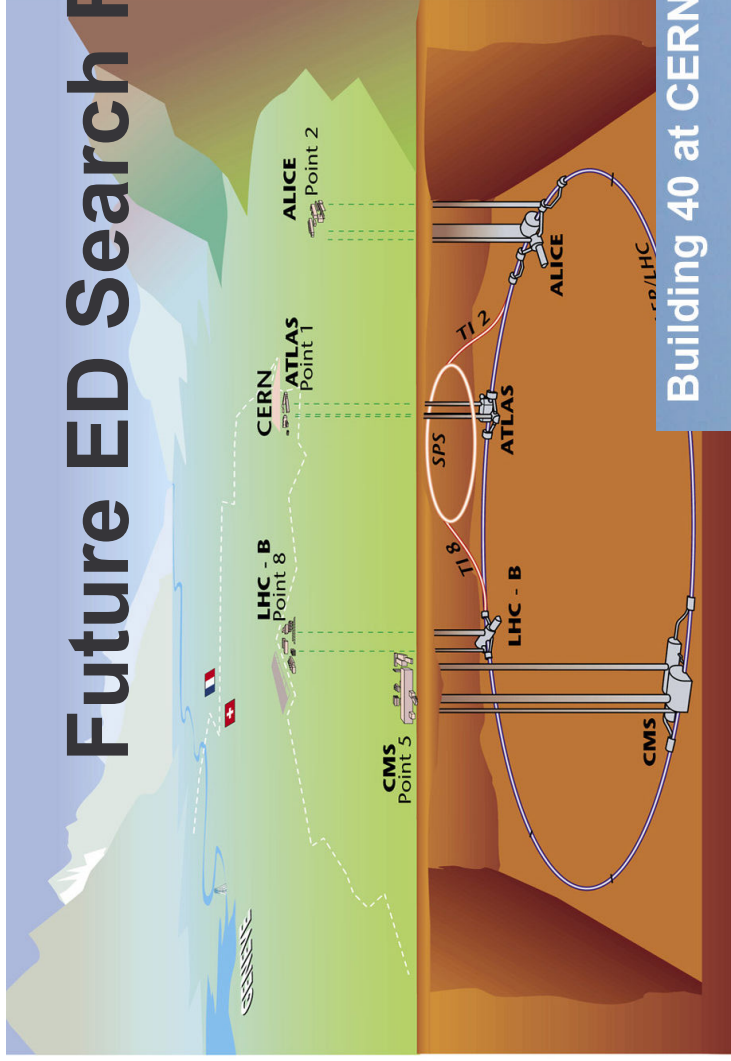
Bigger & Better (!)
Collider & Detectors!!

Higher CM of mass energy
LHC: $\sqrt{s} = 14 \text{ TeV}$

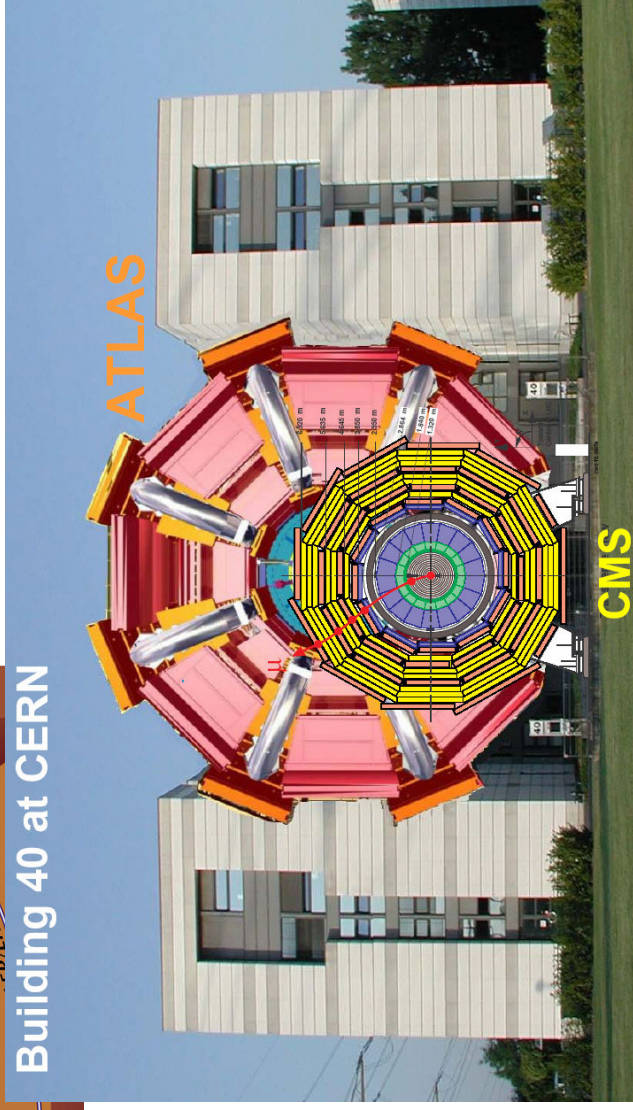


Aim of Tev4LHC
use what has been
learnt at the Tevatron
to fully exploit LHC's
physics potential

Tracey Berry



Building 40 at CERN

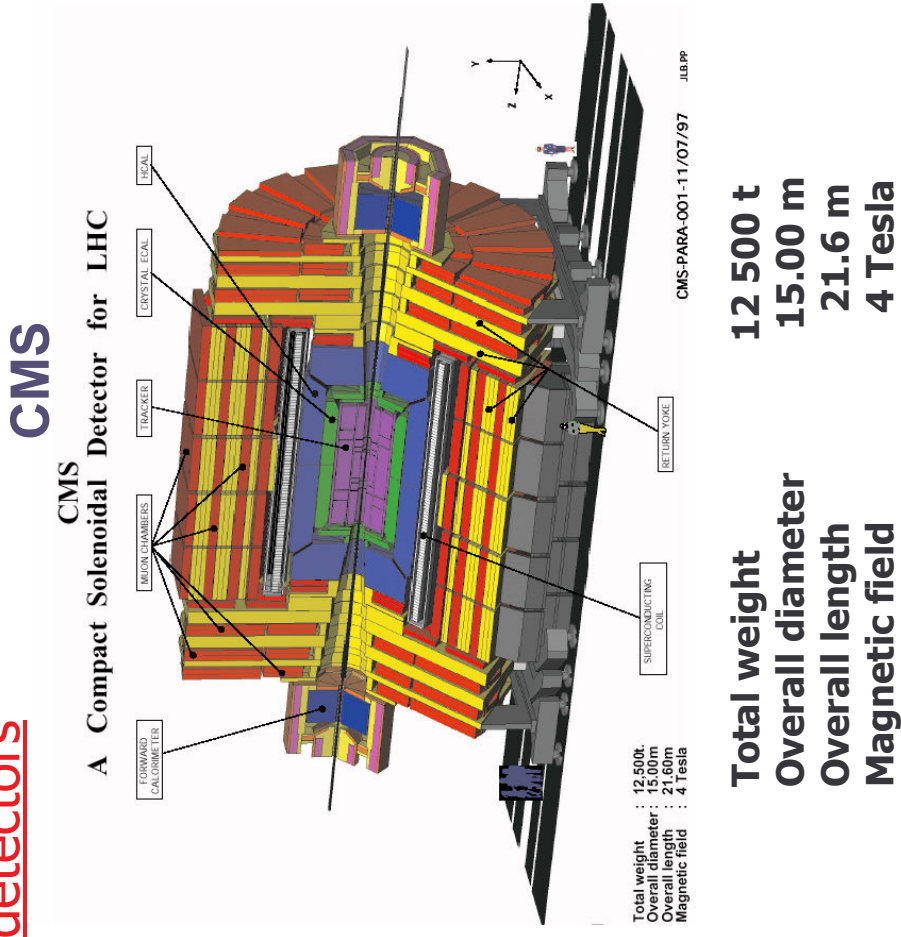
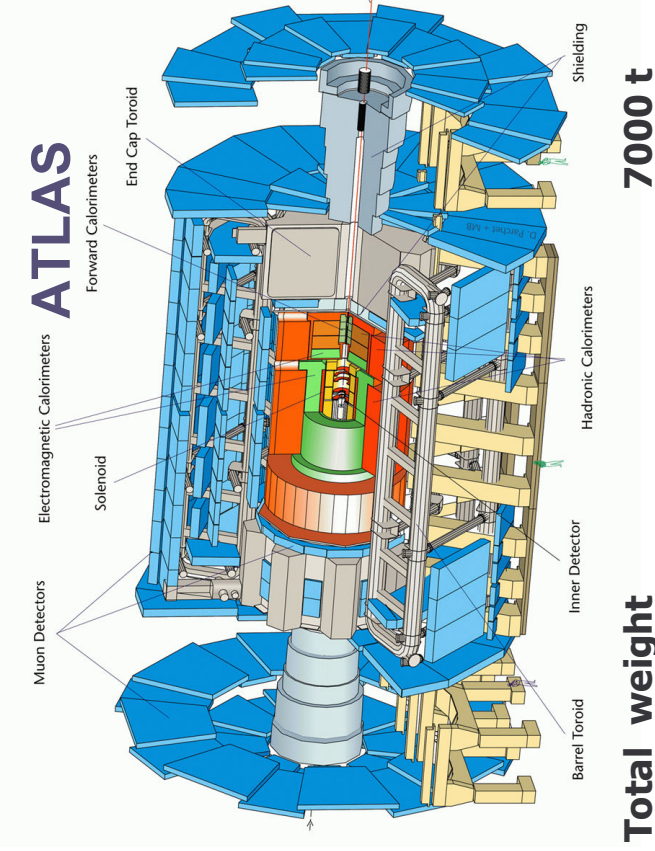


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ATLAS and CMS Experiments



Large general-purpose particle physics detectors



Detector subsystems are designed to measure:
energy and momentum of γ , e , μ , jets, missing E_T up to a few TeV

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Experimental Uncertainties



Systematic uncertainties associated with the detector measurements

- Luminosity
- Energy miscalibration which affects the performance of e/γ /hadron energy reconstruction
- Drift time and drift velocities uncertainties
- Misalignment affects track and vertex reconstruction efficiency → increase of the mass residuals by around 30%
- Magnetic and gravitational field effects → can cause a scale shift in a mass resolution by 5-10%
- Pile-up → mass residuals increase by around 0.1-0.2%
- Trigger and reconstruction acceptance uncertainties
 - Affect the background and signal
- Background uncertainties: variations of the bkgd shape → a drop of about 10-15% in the significance values

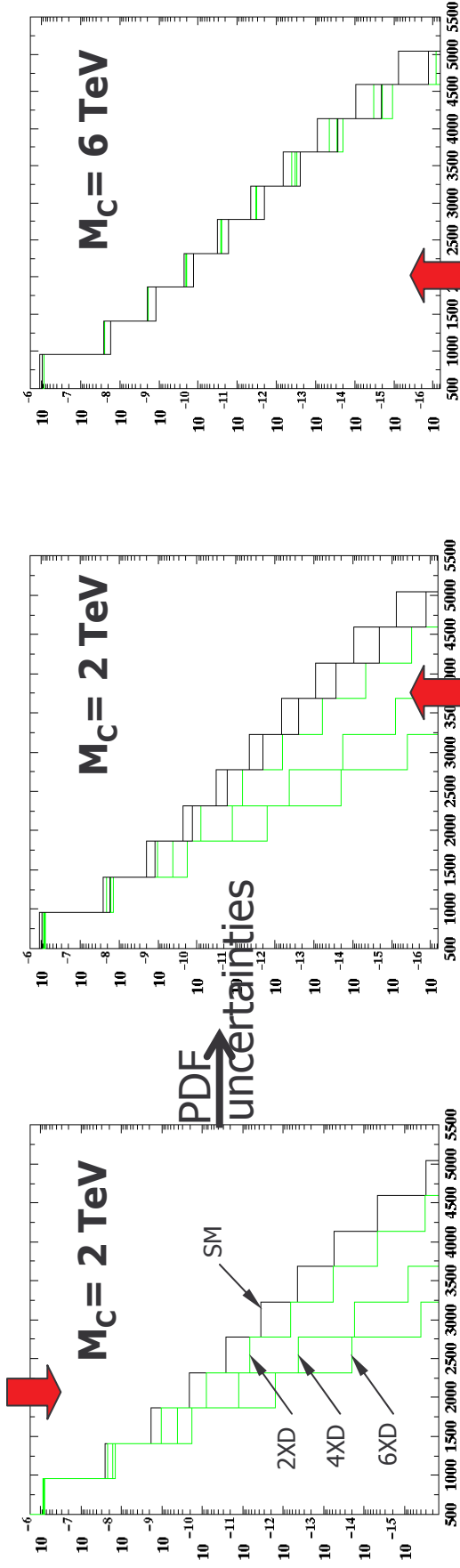
Theoretical Uncertainties



- QCD and EW higher-order corrections (K-factors)
- Parton Distribution Functions (PDF)
- Hard process scale (Q^2)
- Differences between Next-to-Next-to-Leading Order (NNLO), NLO and LO calculations

→ affect signal and background magnitudes,
efficiency of the selection cuts,
significance computation...

- Extra dimensions affect the di-jet cross section through the running of α_s .
 → So could potentially use σ deviation to detect ED
 Parameterised by number of extra dimensions δ and compactification scale M_C .



- PDF uncertainties (mainly due to high-x gluon) give an uncertainty “zone” on the SM cross sections
- This reduces sensitivity to M_C from 5 TeV to 2 (3) TeV for $\delta=4, 6$ and for $\delta=2$ sensitivity is lost ($M_C < 2$ TeV)

ADD Model



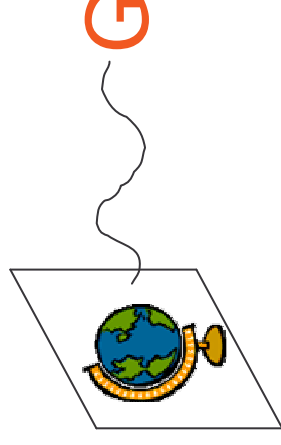
Arkani-Hamed, Dimopoulos, Dvali, Phys Lett B429 (98), Nuc.Phys.B544(1999)

(Many) Large flat Extra-Dimensions (LED),

could be as large as a few μm

the maximum total number of dimensions is $3(\text{our}) + 6(\text{extra})=9$

G can propagate in ED



SM particles restricted to 3D brane

The fundamental scale is not planckian: $M_D = M_{\text{Pl}(4+\delta)} \sim \text{TeV}$

Model parameters are:

- δ = number of ED
- $M_{\text{Pl}(4+\delta)}$ = Planck mass in the $4+\delta$ dimensions

$$M_{\text{Pl}}^2 \sim R^\delta M_{\text{Pl}(4+\delta)}^{(2+\delta)}$$

$$\text{For } M_{\text{Pl}} \sim 10^{19} \text{ GeV and } M_{\text{Pl}(4+\delta)} \sim M_{\text{EW}} \rightarrow R \sim 10^{32/\delta} \times 10^{-17} \text{ cm}$$

Tracey Berry

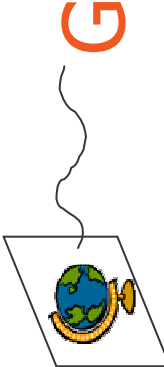
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Present Constraints on the ADD Model



$$M_{\text{Pl}}^2 \sim R^\delta M_{\text{Pl}(4+\delta)}^{(2+\delta)} \rightarrow \boxed{R \sim 10^{32/\delta} \times 10^{-17} \text{ cm}}$$

For $M_{\text{Pl}} \sim 10^{19} \text{ GeV}$ and $M_{\text{Pl}(4+\delta)} \sim M_{\text{EW}} \rightarrow$



- $\delta=1 \rightarrow R \sim 10^{13} \text{ cm}$, ruled out because deviations from Newtonian gravity over solar distances have not been observed
- $\delta=2 \rightarrow R \sim 1 \text{ mm}$, not likely because of cosmological arguments:

In particular graviton emission from Supernova 1987a* implies $M_D > 50 \text{ TeV}$
 Closest allowed $M_{\text{Pl}(4+n)}$ value for $\delta=2$ is $\sim 30 \text{ TeV}$, out of reach at LHC

Can detect at collider detectors via:

- ❖ graviton emission
- ❖ Or graviton exchange

*Cullen, Perelstein
 Phys. Rev. Lett 83,268 (1999)

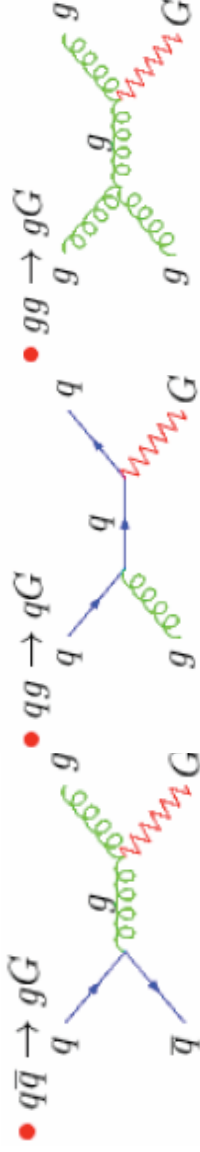
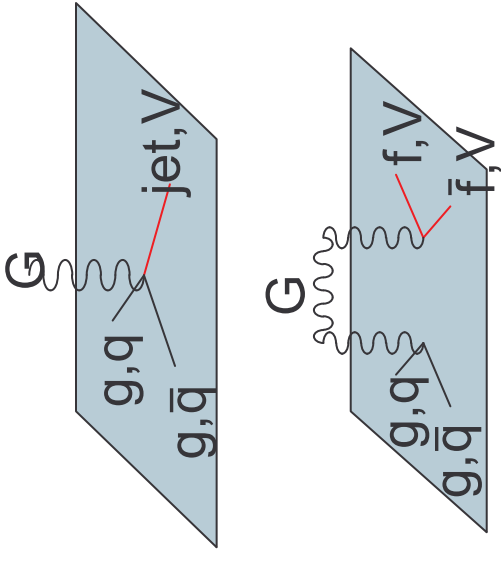
ADD Collider Signatures



➤ Real Graviton emission in association with a vector-boson

Signature: jets + missing E_T , V+missing E_T

σ depends on the number of ED

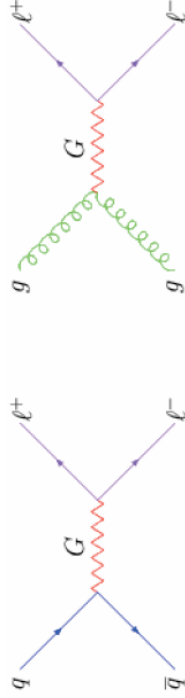


➤ Virtual Graviton exchange

Signature:

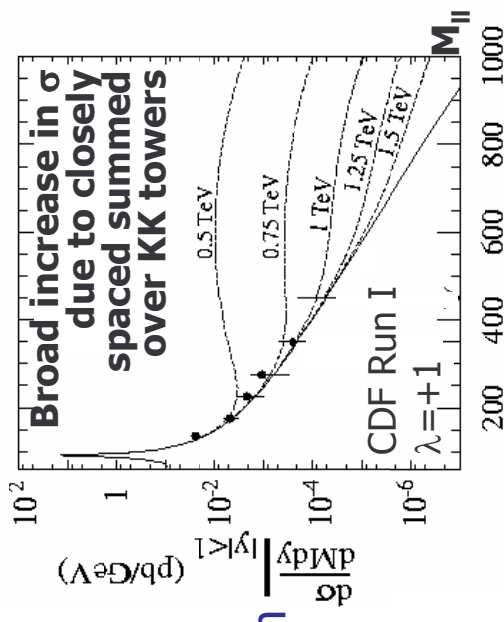
deviations in σ and asymmetries of SM processes
e.g. $q\bar{q} \rightarrow l^+l^-$, $\gamma\gamma$ & new processes e.g. $g\bar{g} \rightarrow l^+l^-$

$q\bar{q} \rightarrow l^+l^-$ $g\bar{g} \rightarrow l^+l^-$
Excess above
di-lepton continuum



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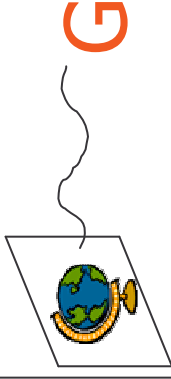


Present Constraints on the ADD Model



$$M_{\text{Pl}}^2 \sim R^\delta M_{\text{Pl}(4+\delta)}^{(2+\delta)}$$

For $M_{\text{Pl}} \sim 10^{19}$ GeV and $M_{\text{Pl}(4+\delta)} \sim M_{\text{EW}} \rightarrow \boxed{R \sim 10^{32/\delta} \times 10^{-17} \text{ cm}}$



➤ $\delta=1 \rightarrow R \sim 10^{13}$ cm, ruled out because deviations from Newtonian gravity over solar distances have not been observed

➤ $\delta=2 \rightarrow R \sim 1$ mm, not likely because of cosmological arguments:

In particular graviton emission from Supernova 1987a* implies $M_D > 50$ TeV
Closest allowed $M_{\text{Pl}(4+n)}$ value for $\delta=2$ is ~ 30 TeV, out of reach at LHC

➤ LEP & Tevatron limits is $M_{\text{Pl}(4+\delta)} \sim > 1 \text{ TeV}$

➤ $\delta > 6$ difficult to probe at LHC since cross-sections are very low

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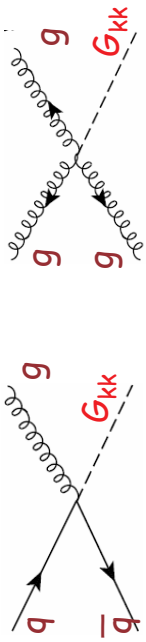
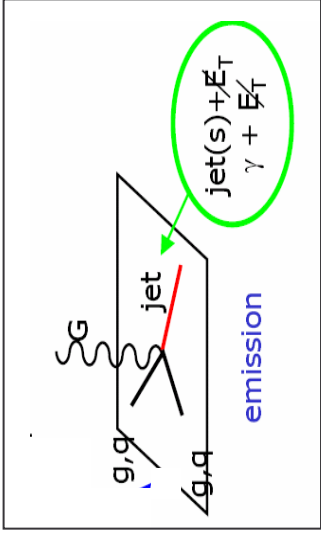
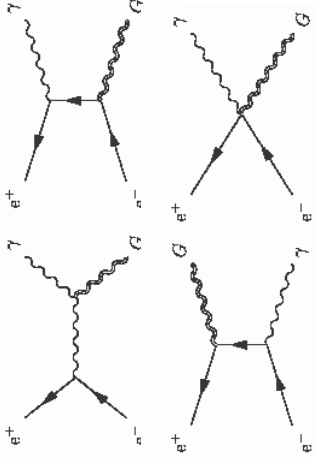
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*Cullen, Perelstein
Phys. Rev. Lett 83,268 (1999)

Present ADD Emission Limits



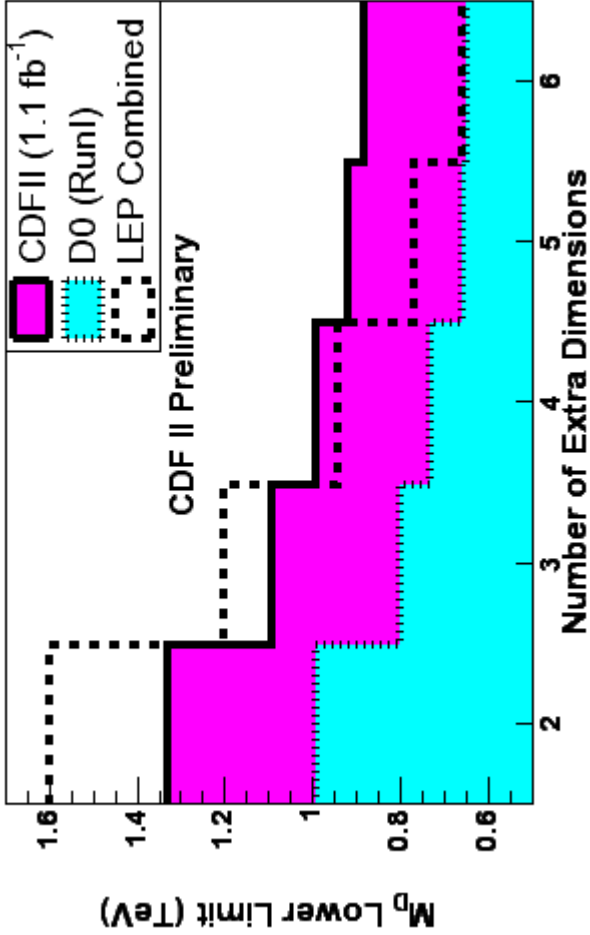
LEP and Tevatron results are complementary



For $n < 4$: LEP limits best $\gamma + ME_T$

For $n > 4$: CDF limits best $jet + ME_T$

For $n < 4$: LEP limits best $\gamma + ME_T$



n	M_D (TeV/c ²) K=1.3	R (mm)
2	> 1.33	< 0.27
3	> 1.09	< 3.1×10^{-6}
4	> 0.99	< 9.9×10^{-9}
5	> 0.92	< 3.2×10^{-10}
6	> 0.88	< 3.1×10^{-11}

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ADD Discovery Limit: $\gamma + G$ Emission



Real graviton production

$$pp \rightarrow \gamma + G^{KK}$$

J. Weng et al. CMS NOTE 2006/129

- $\gamma G \Rightarrow$ high- p_T photon + high missing E_T

At low p_T the bkgd, particularly irreducible $Z\gamma \rightarrow \nu\nu\gamma$ is too large $\Rightarrow$ require $p_T > 400$ GeV

- Main Bkgd: $Z\gamma \rightarrow \nu\nu\gamma$,
Also $W \rightarrow e(\mu, \tau)\nu$, $W\gamma \rightarrow e\nu$,
 $\gamma + \text{jets}$, QCD, di- γ , $Z^0 + \text{jets}$
- Signals generated with PYTHIA (compared to SHERPA)
Bkgds: PYTHIA and compared to SHERPA/CompHEP/Madgraph (B),
Using CTEQ6L
- Full simulation & reconstruction
- Theoretical uncert.

Not considered by CMS analysis: Cosmic Rays at rate of 11 HZ: main background at CDF, also beam halo muons for $p_T > 400$ GeV rate 1 HZ

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Integrated Lum for a 5σ significance discovery

M_D/n	$n=2$	$n=3$	$n=4$	$n=5$	$n=6$
Significance: $S=2(\sqrt{(S+B)}-\sqrt{B}) > 5$					
$M_D = 1.0$ TeV	0.21 fb $^{-1}$	0.16 fb $^{-1}$	0.14 fb $^{-1}$	0.15 fb $^{-1}$	0.15 fb $^{-1}$
$M_D = 1.5$ TeV	0.83 fb $^{-1}$	0.59 fb $^{-1}$	0.56 fb $^{-1}$	0.61 fb $^{-1}$	0.59 fb $^{-1}$
$M_D = 2.0$ TeV	2.8 fb $^{-1}$	2.1 fb $^{-1}$	1.9 fb $^{-1}$	2.1 fb $^{-1}$	2.3 fb $^{-1}$
$M_D = 2.5$ TeV	9.9 fb $^{-1}$	8.2 fb $^{-1}$	8.7 fb $^{-1}$	9.4 fb $^{-1}$	10.9 fb $^{-1}$
$M_D = 3.0$ TeV	47.8 fb $^{-1}$	46.4 fb $^{-1}$	64.4 fb $^{-1}$	100.8 fb $^{-1}$	261.2 fb $^{-1}$
$M_D = 3.5$ TeV	5 σ discovery not possible anymore				

$M_D =$ 1- 1.5 TeV for 1 fb $^{-1}$
2 - 2.5 TeV for 10 fb $^{-1}$
3 - 3.5 TeV for 60 fb $^{-1}$



ADD Discovery Limit: γ +G Emission



ATLAS

$$pp \rightarrow \gamma + G^{KK} : q\bar{q} \rightarrow \gamma G^{KK}$$

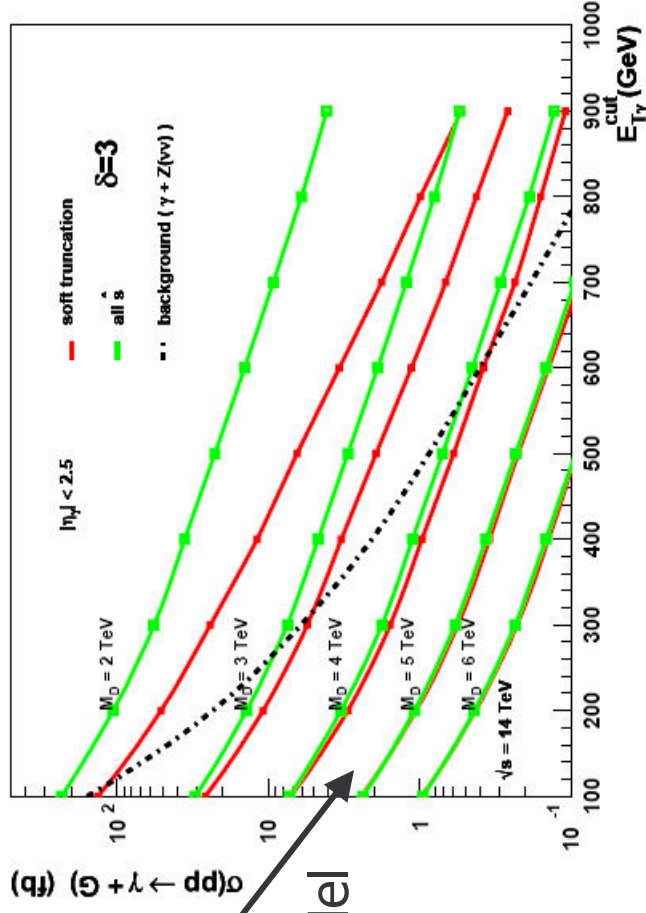
Rates for $M_D \geq 4\text{TeV}$ are very low

$M_D^{\text{MAX}} (\text{TeV})$	$\delta=2$
HL 100fb ⁻¹	4

For $\delta > 2$: No region where the model made independent predictions can be made and where the rate is high enough to observe signal events over the background.

This gets worse as δ increases

- Better limits from the jet+G emission which has a higher production rate
- This signature could be used as confirmation after the discovery in the jet channels



L. Vacavant, I. Hinchcliffe
 ATLAS-PHYS 2000-016
 J. Phys., G 27 (2001) 1839-50

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ADD Discovery Limit: jet+G Emission



Real graviton production

$$pp \rightarrow \text{jet} + G^{KK}$$

$$gg \rightarrow gG, qg \rightarrow qG \text{ \& } q\bar{q} \rightarrow Gg$$



Dominant subprocess

□ Signature: jet + G $\Rightarrow$ jet with high transverse energy ($E_T > 500$ GeV) + high missing E_T ($E_{T\text{miss}} > 500$ GeV),

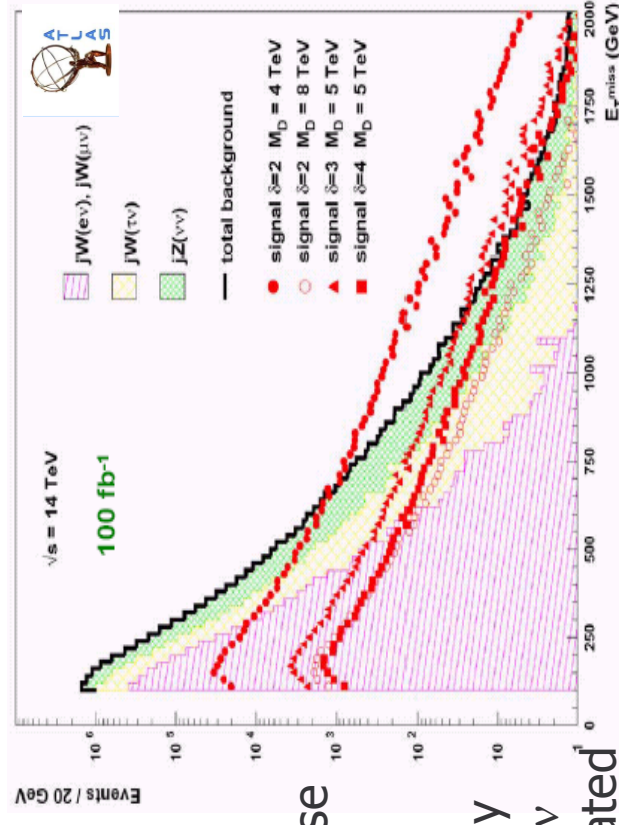
□ vetos leptons: to reduce jet+W bkgd mainly

□ Bkgd.: irreducible jet+Z/W $\rightarrow$ jet+ $\nu\nu$ /jet+l ν

jZ($\nu\nu$) dominant bkgd, can be calibrated using ee and $\mu\mu$ decays of Z.

□ ISAJET with CTEQ3L

□ Fast simulation/reco



Discovery limits

$M_{Pl(4+d)}^{MAX}(\text{TeV})$	$\delta=2$	$\delta=3$	$\delta=4$
LL 30fb ⁻¹	7.7	6.2	5.2
HL 100fb ⁻¹	9.1	7.0	6.0

Tracey Berry

L.Vacavant, I.Hinchcliffe, ATLAS-PHYS 2000-016

Getting ready for the LHC, Madrid J. Phys., G 27 (2001) 1839-50

Oct 23-27, 2006

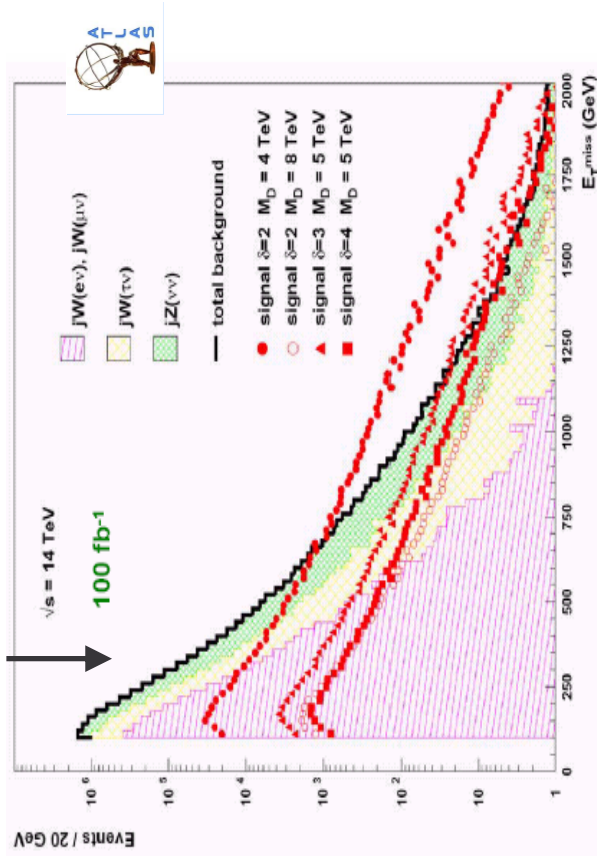


ADD Parameters: jet+G Emission



To characterise the model need to measure M_D and δ

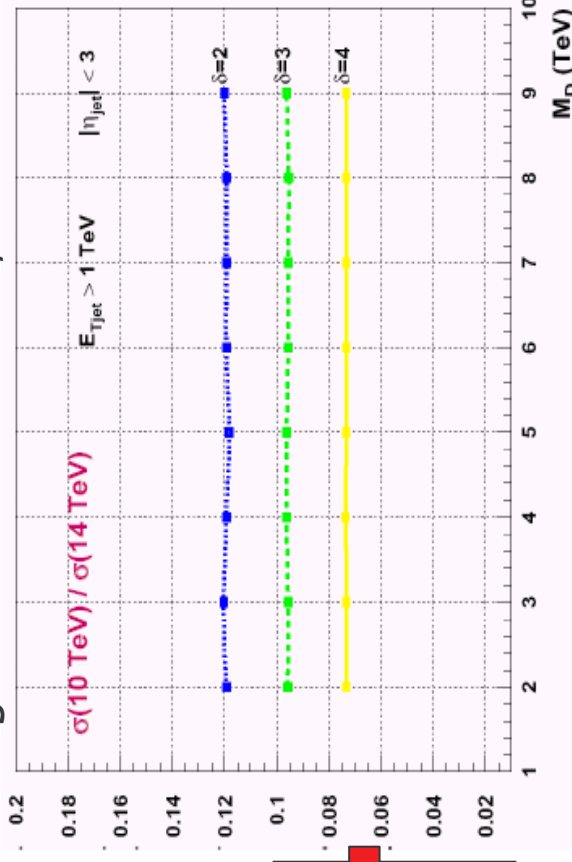
Measuring $\sigma(pp \rightarrow \text{jet} + G^{KK})$ gives ambiguous results



Use variation of σ on $\sqrt{s}$
 σ at different $\sqrt{s}$ almost
independent of M_D , varies with δ

Run at two different $\sqrt{s}$

e.g. 10 TeV and 14 TeV, need 50 fb⁻¹



Rates at 14 TeV of $\delta=2, M_D=6$ TeV very
similar to $\delta=3, M_D=5$ TeV whereas
Rates at 10 TeV of ($\delta=2, M_D=6$ TeV) and
($\delta=3, M_D=5$ TeV) differ by $\sim$ factor of 2

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L. Vacavant, I. Hinchcliffe, ATLAS-PHYS 2000-016
J. Phys., G 27 (2001) 1839-50



Tevatron ADD Exchange Limits



Both D0 and CDF have observed no significant excess

95% CL lower limits on fundamental Planck scale (M_s) in

TeV, using different formalisms:

most stringent
collider limits on
LED to date!

collider limits on LED to date!

	GRW	HLZ for n=						Hewett
		2	3	4	5	6	7	$\lambda = +1/-1$
D0 Run II: $\mu\mu$								
D0 Run II: $ee+\gamma\gamma$	1.09	1.00	1.29	1.09	0.98	0.91	0.86	0.97/0.95
D0 Run I+II: $ee+\gamma\gamma$	1.36	1.56	1.61	1.36	1.23	1.14	1.08	1.22/1.10
D0 Run I+II: $ee+\gamma\gamma$	1.43	1.61	1.70	1.43	1.29	1.20	1.14	1.28/NA
CDF Run II: ee 200pb ⁻¹	1.11	1.32 1.11 1.00 0.93 0.88						0.96/0.99

D0 perform a 2D search in invariant mass & angular distribution

And to maximise reconstruction efficiency they perform combined $ee+\gamma\gamma$ (diEM) search: reduces inefficiencies from

- γ ID requires no track, but γ converts ($\rightarrow ee$)
- e ID requires a track, but loose track due to imperfect track reconstruction/crack





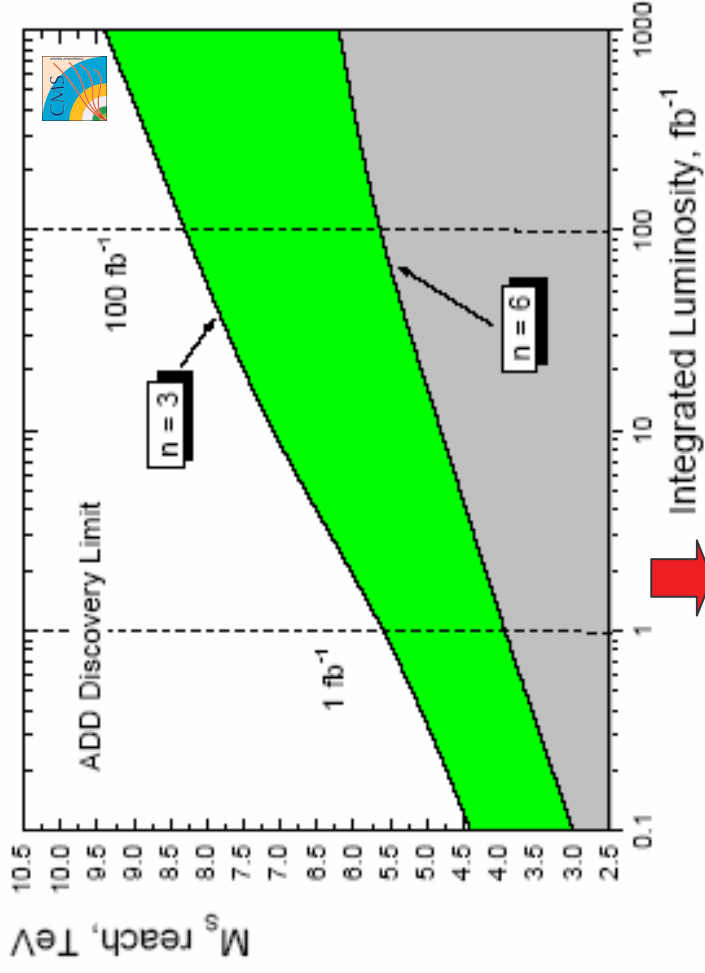
ADD Discovery Limit: G Exchange



Virtual graviton production

$$pp \rightarrow G^{KK} \rightarrow \mu\mu$$

- ☐ Two opposite sign muons in the final state with $M_{\mu\mu} > 1 \text{ TeV}$
- ☐ Irreducible background from Drell-Yan, also ZZ, WW, tt (suppressed after selection cuts)
- ☐ PYTHIA with ISR/FSR + CTEQ6L, LO + K=1.38
- ☐ Full (GEANT-4) simulation/reco + L1 + HLT(riger)
- ☐ Theoretical uncert.
- ☐ μ and tracker misalignment, trigger and off-line recon. inefficiency, acceptance due to PDF



Confidence limits for

1 fb ⁻¹ :	3.9-5.5 TeV for n=6..3
10 fb ⁻¹ :	4.8-7.2 TeV for n=6..3
100 fb ⁻¹ :	5.7-8.3 TeV for n=6..3
300 fb ⁻¹ :	5.9-8.8 TeV for n=6..3

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Belotelov et al.,
CMS NOTE 2006/076, CMS PTDR 2006
22



ADD Discovery Limit: G Exchange



Virtual graviton production



channel	n					
		luminosity	2	3	4	5
$\gamma\gamma$	10 fb ⁻¹	M_S^{max} (TeV)	6.3	5.6	5.1	4.9
		S/B	36/18	36/18	39/25	34/13
	100 fb ⁻¹	M_S^{max} (TeV)	7.9	7.3	6.7	6.3
		S/B	50/53	62/96	55/72	51/53
l^+l^-	10 fb ⁻¹	M_S^{max} (TeV)	6.6	5.9	5.4	5.1
		S/B	33/14	31/8	30/6	30/6
	100 fb ⁻¹	M_S^{max} (TeV)	7.9	7.5	7.0	6.6
		S/B	49/48	38/21	36/16	29/6
$\gamma\gamma + l^+l^-$	10 fb ⁻¹	M_S^{max} (TeV)	7.0	6.3	5.7	5.4
	100 fb ⁻¹	M_S^{max} (TeV)	8.1	7.9	7.4	7.0

V. Kabachenko et al.
ATL-PHYS-2001-012

Fast MC

CMS Confidence limits:

1 fb⁻¹: 3.9-5.5 TeV for n=6..3
10 fb⁻¹: 4.8-7.2 TeV for n=6..3
100 fb⁻¹: 5.7-8.3 TeV for n=6..3
300 fb⁻¹: 5.9-8.8 TeV for n=6..3

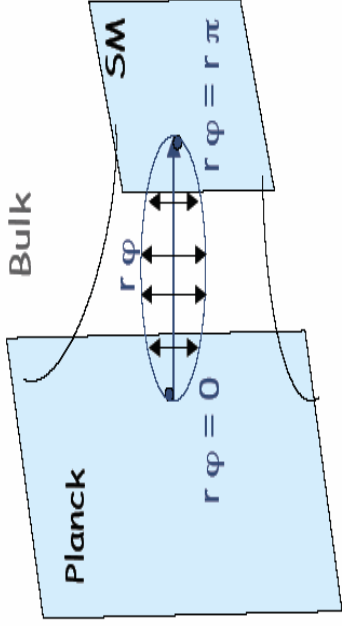


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Belotelov et al.,
CMS NOTE 2006/076, CMS PTDR 2006
23

Experimental Signature for RS Model



Signature:

Narrow, high-mass resonance states
in dilepton/dijet/diboson channels

$$q\bar{q}, gg \rightarrow G_{KK} \rightarrow e^+e^-, \mu^+\mu^-, \gamma\gamma, jet + jet$$

At the LHC only the 1st excitations are likely to be seen at the LHC, since the other modes are suppressed by the falling parton distribution functions.

Allenach et al, JHEP 9 19 (2000), JHEP 0212 39 (2002)

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Davoudiasl, Hewett, Rizzo
hep-ph/0006041

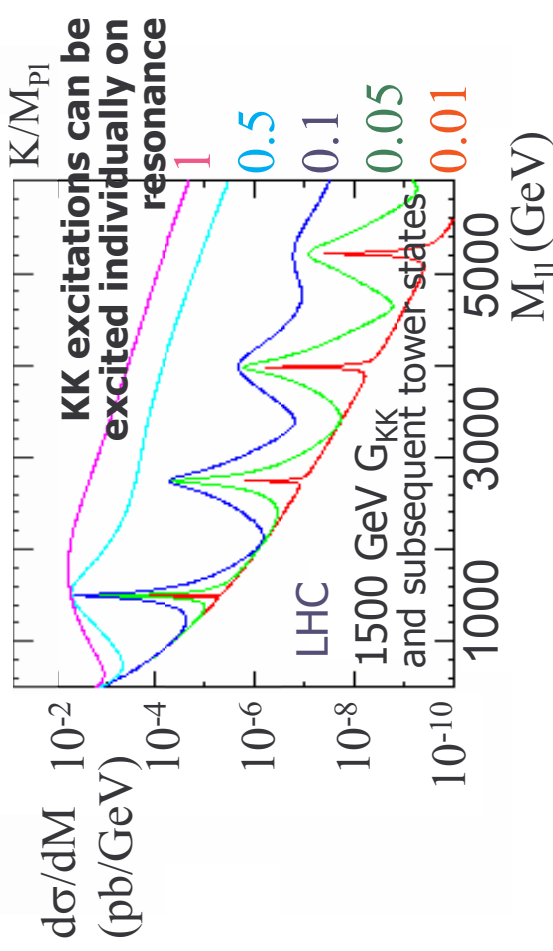
24

Model parameters:

- Gravity Scale: $\Lambda_\pi = \overline{M}_{pl} e^{-kR_c \pi}$ **Resonance**
1st graviton excitation mass: $m_1 \rightarrow$ position
- $\Lambda_\pi = m_1 M_{pl} / k x_1$, & $m_n = k x_n e^{k r c \pi} (J_1(x_n) = 0)$
- Coupling constant: $c = k / M_{pl}$

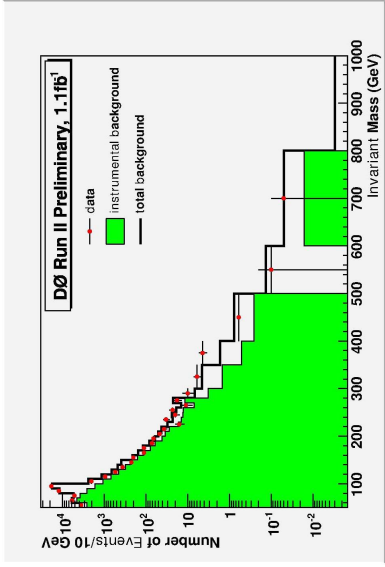
$$\Gamma_1 = \rho m_1 x_1^2 (k / M_{pl})^2 \rightarrow \text{width}$$

k = curvature, R = compactification radius

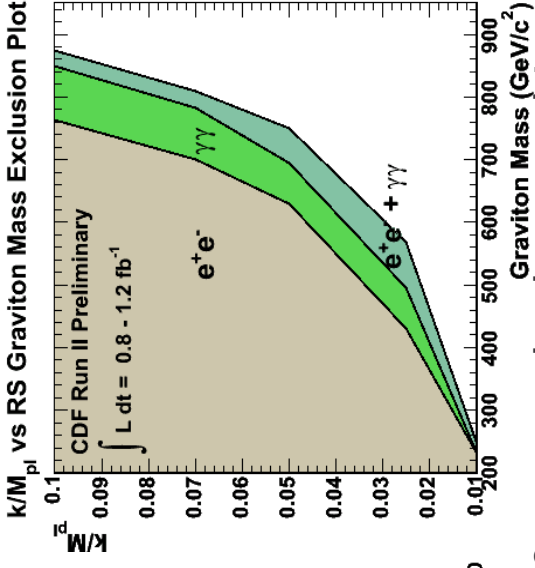
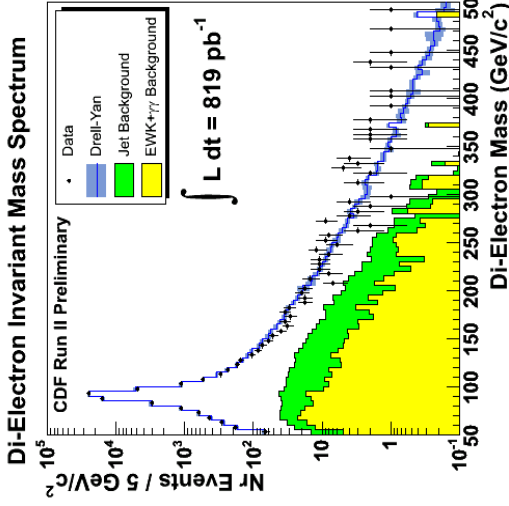




Present RS Constraints

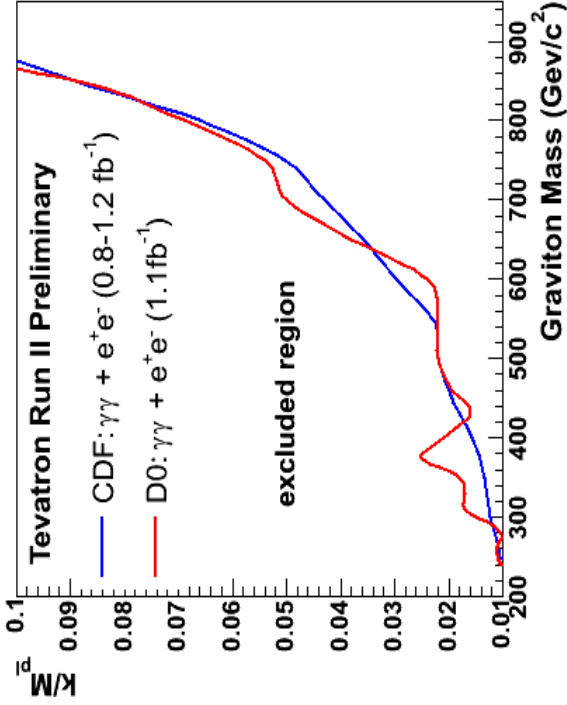


D0 performed combined $ee + \gamma\gamma$ (diem search)



CDF performed ee & $\gamma\gamma$ search, then combine

k/M_{pl} vs RS Graviton Mass Exclusion



Present Experimental Limits

Theoretical Constraints

- $c > 0.1$ disfavoured as bulk curvature becomes to large (larger than the 5-dim Planck scale)
- Theoretically preferred $\Lambda_\pi < 10 \text{ TeV}$ assures no new hierarchy appears between m_{EW} and Λ_π

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RS1 Discovery Limit



- Best channels to search in are $G(1) \rightarrow e+e-$ and $G(1) \rightarrow \gamma\gamma$ due to the energy and angular resolutions of the LHC detectors
- $G(1) \rightarrow e+e-$ best chance of discovery due to relatively small bkgd, from Drell-Yan*

Di-electron

- HERWIG
- Main Bkgd: Drell-Yan
- Model-independent analysis
- RS model with $k/M_{Pl}=0.01$ as a reference (pessimistic scenario)
- Fast Simulation

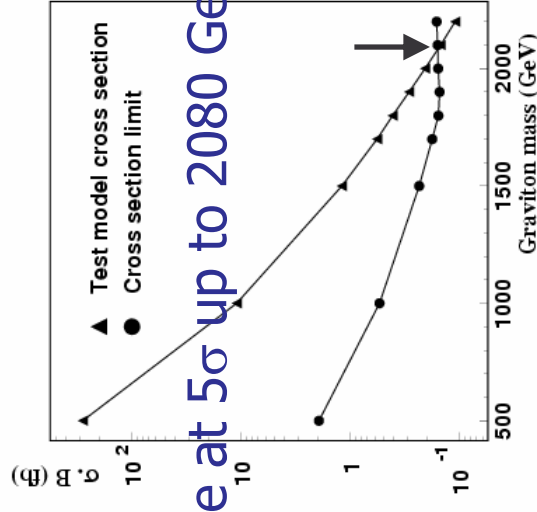
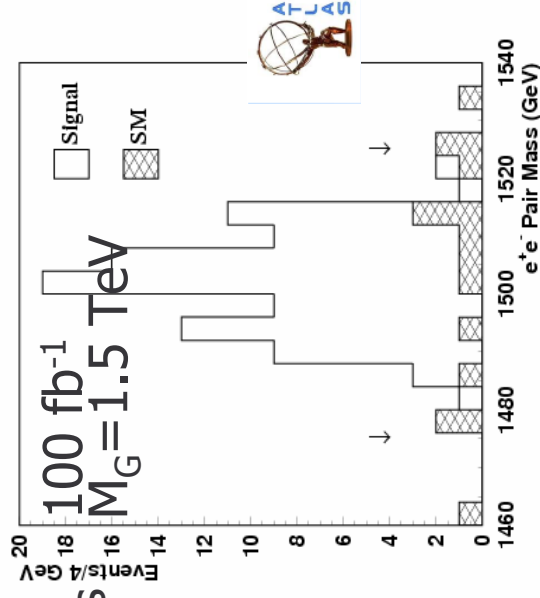
*Reach goes up to 3.5 TeV for $c=0.1$ for a 20% measurement of the coupling.

Allenach et al, hep-ph0006114

*Allenach et al, hep-ph0211205

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Sensitive at 5σ up to 2080 GeV



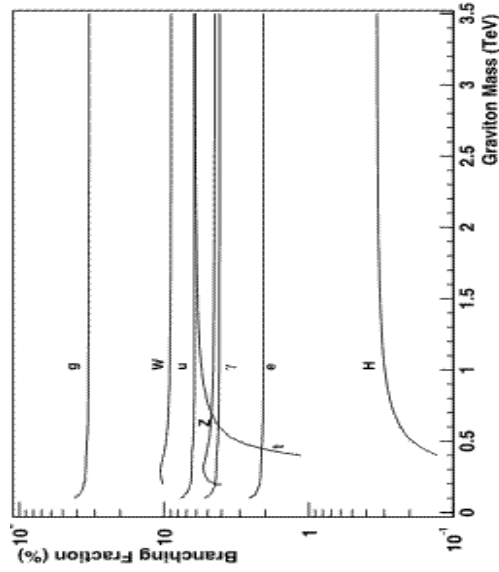
RS1 Model Parameters



A resonance could be seen in many other channels: $\mu\mu$, $\gamma\gamma$, $j\bar{j}$, $b\bar{b}$, $t\bar{t}$, WZ , ZZ , hence allowing to check universality of its couplings:

Channel	Point m_G, Λ_π (TeV)					
	1,10	1,20	1,30	2,10	2,20	2,30
e^+e^-	1.6	3.3	5.3	5.4	11.0	17.1
$\mu^+\mu^-$	1.9	4.5	8.2	6.2	15.2	28.2
$\gamma\gamma$	1.2	2.9	5.2	3.9	8.8	15.2
WW	11.6	44.9	-	38.2	-	-
ZZ	13.7	50.1	-	52.7	-	-
$j\bar{j}$	19.0	77.0	-	31.0	-	59.0

Relative precision achievable (in %) for measurements of σ_B in each channel for fixed points in the M_G/Λ_π plane. Points with errors above 100% are not shown.



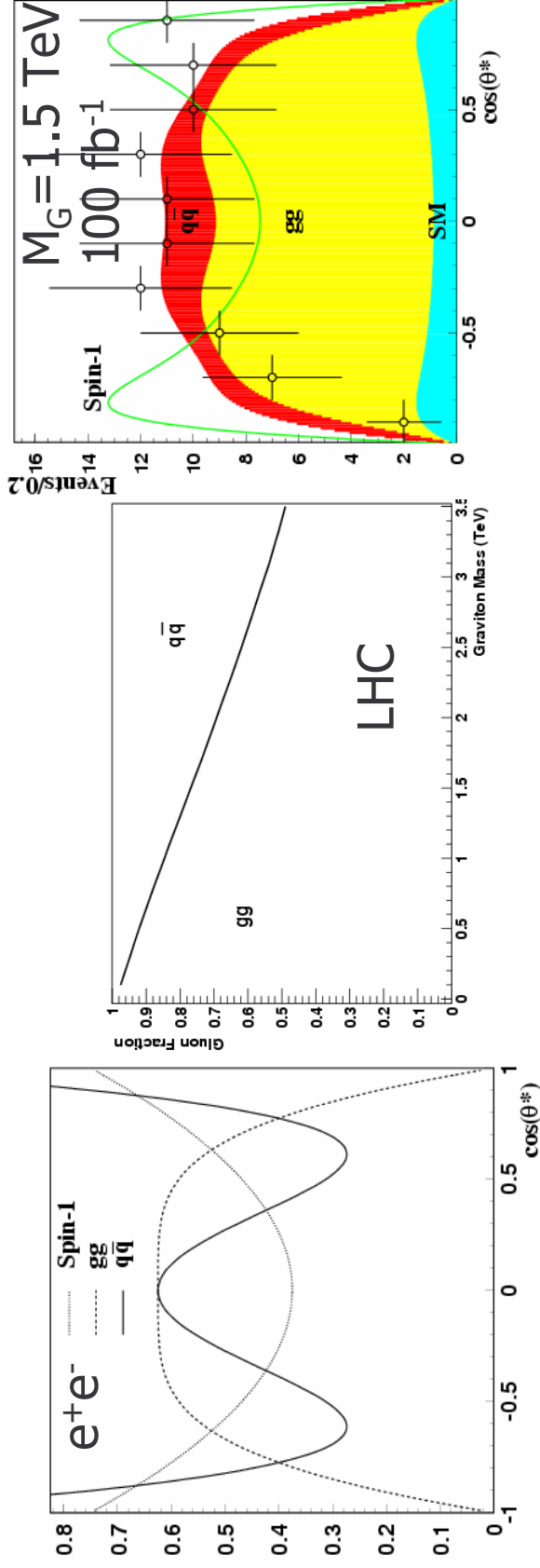
Also the size (R) of the ED could also be estimated from mass and cross-section measurements.



RS1 Model Determination



Spin-2 could be determined (spin-1 ruled out) with 90% C.L. up to $M_G = 1720$ GeV



Note: acceptance at large pseudo-rapidities is essential for spin discrimination ($1.5 < |\eta| < 2.5$)

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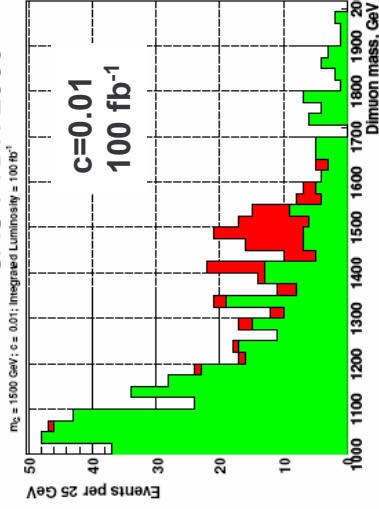
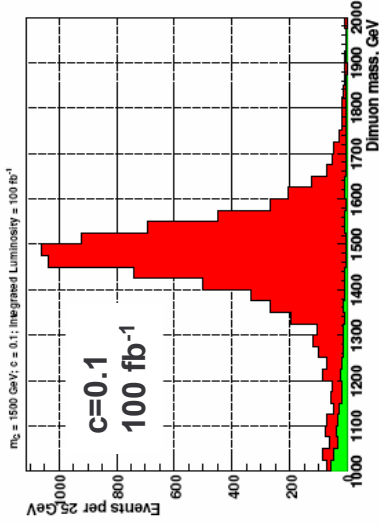
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Allanach et al, hep-ph 0006114

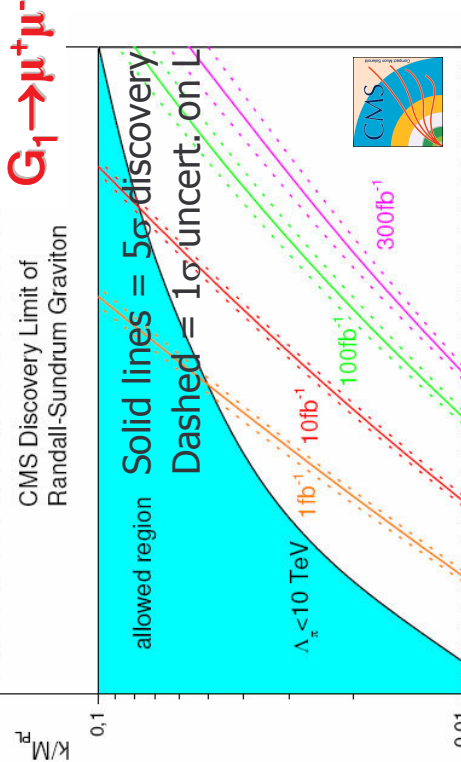


RS1 Discovery Limit

Di-lepton states



I. Belotelov et al.
 CMS NOTE 2006/104
 CMS PTDR 2006



- Two muons/electrons in the final state

- Bckg: Drell-Yan/ZZ/WW/ZW/ttbar

- PYTHIA/CTEQ6L

- LO + K=1.30 both for signal and DY

- Full (GEANT-4) and fast simulation/reco

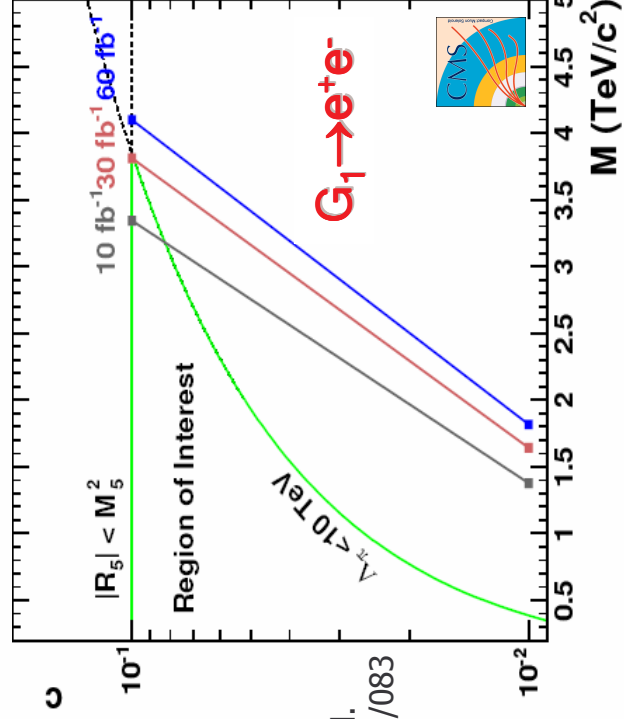
- Viable L1 + HLT(rigger) cuts

- Theoretical uncert.

- Misalignment, trigger and off-line reco inefficiency, pile-up

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B. Clerbaux et al.
 CMS NOTE 2006/083
 CMS PTDR 2006



RS1 Discovery Limit



Di-photon states

- Two photons in the final state
- Bckg: prompt di-photons, QCD hadronic jets and gamma+jet events, Drell-Yan e^+e^-
- PYTHIA/CTEQ5L
- LO for signal, LO + K-factors for bckg.
- Fast simulation/reco + a few points with full GEANT-4 MC
- Viable L1 + HLT(rigger) cuts
- Theoretical uncert.
- Preselection inefficiency

M.-C. Lemaire et al.
CMS NOTE 2006/051
CMS PTDR 2006

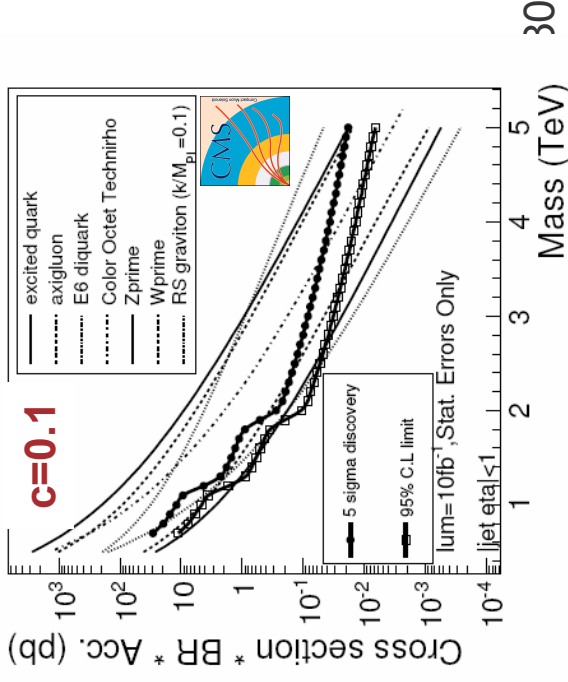
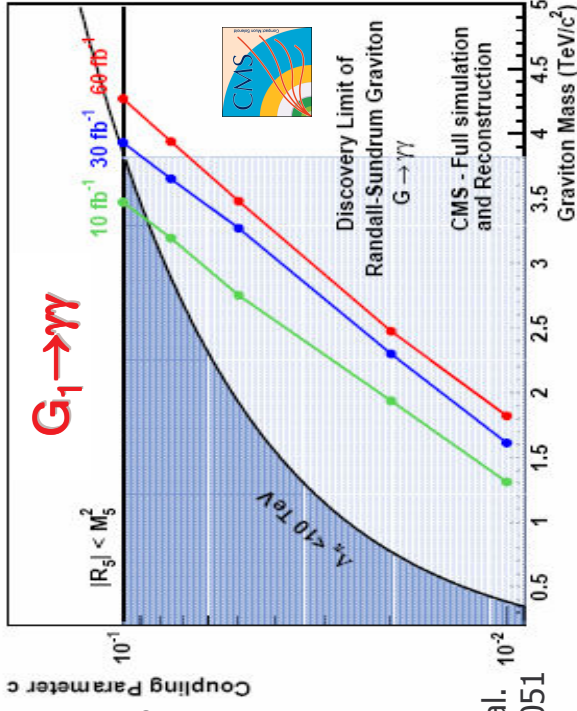
Di-jet states

- Bckg: QCD hadronic jets
- L1 + HLT(rigger) cuts

5 σ Discovered Mass: 0.7-0.8 TeV/ c^2

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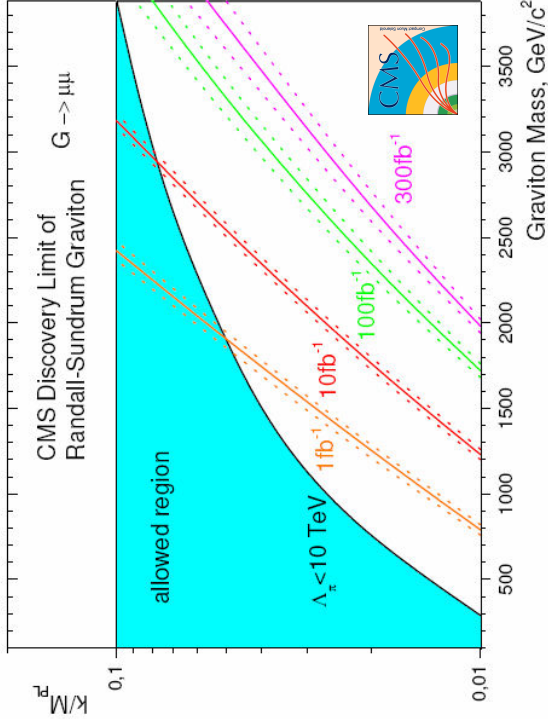




CMS RS Discovery Limits



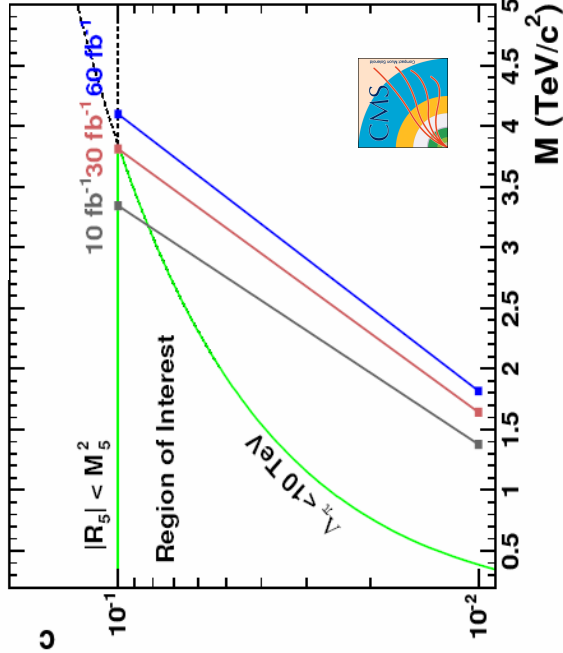
$G_1 \rightarrow \mu^+ \mu^-$



$c > 0.1$ disfavoured as bulk curvature becomes to large (larger than the 5-dim Planck scale)

Theoretically preferred $\Lambda_\pi < 10 \text{ TeV}$

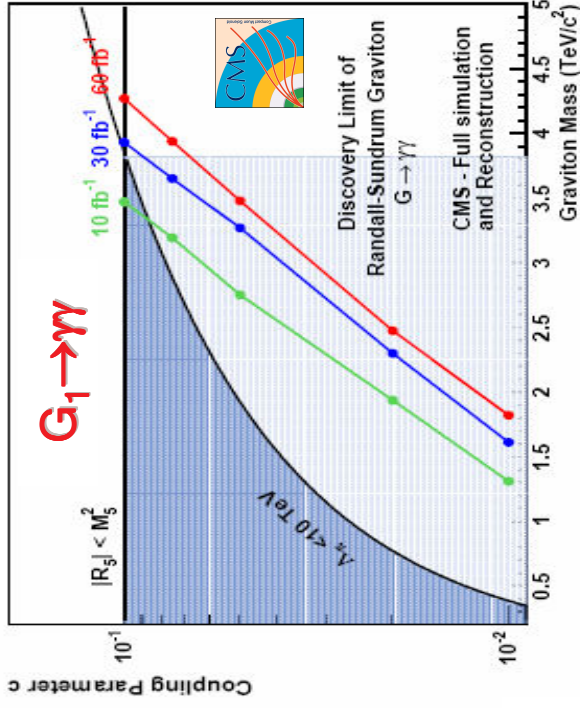
$G_1 \rightarrow e^+ e^-$



LHC completely covers the region of interest

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Getting



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TeV⁻¹ Extra Dimension Model



I. Antoniadis, PLB246 377 (1990)

- Multi-dimensional space with orbifolding (5D in the simplest case, $n=1$)
- The fundamental scale is not planckian: $M_D \sim \text{TeV}$
- Gauge bosons can travel in the bulk
 $\Rightarrow$ Search for KK excitations of $Z, \gamma, \dots$

New Parameters

$R = M_C^{-1}$: size of the compact dimension

M_C : corresponding compactification scale

M_0 : mass of the SM gauge boson

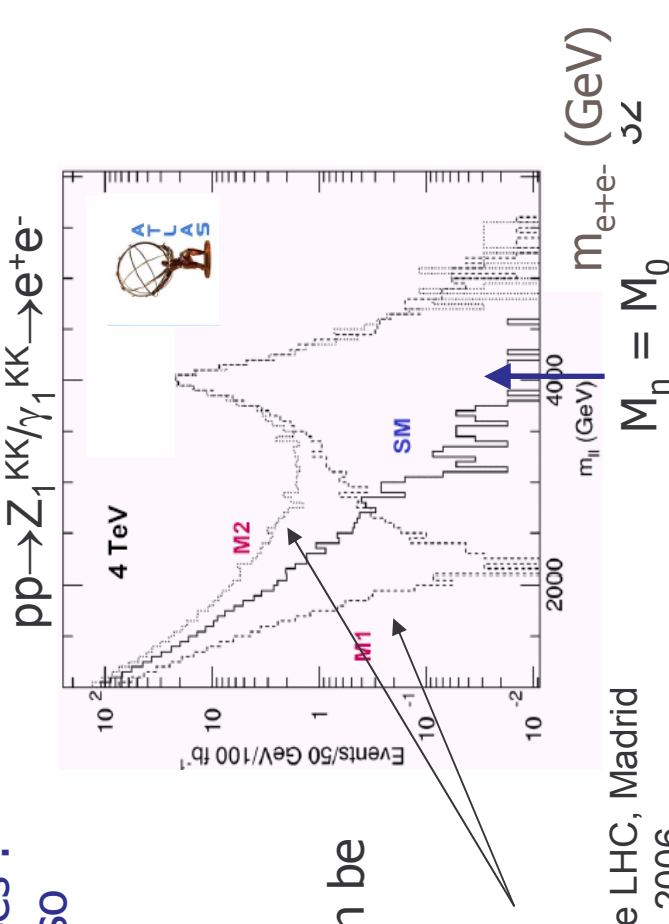
Characteristic Signature: KK excitations of the gauge bosons appearing as resonances with masses :
 $M_n = \sqrt{(M_0^2 + n^2/R^2)}$ where ($n=1, 2, \dots$) & also interference effects!

- Fermion-gauge boson couplings can be exponentially suppressed for higher KK-modes
- Fundamental fermions (quarks/leptons) can be localized at the same (M1) or opposite (M2) points of orbifold
 $\Rightarrow$ destructive (M1) or constructive (M2) interference of the KK excitations with SM model gauge bosons

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G. Azuelos, G. Polesello
 EPJ Direct 10.1140 (2004)





Present Constraints on TeV^{-1} ED



DØ performed the first dedicated experimental search for TeV^{-1} ED at a collider

Search for effects of virtual exchanges of the KK states of the Z and γ

Search Signature: Signal has 2 distinct features:

- enhancement at large masses (like LED)
- negative interference between the 1st KK state of the Z/ γ and the SM Drell-Yan in between the Z mass and M_C

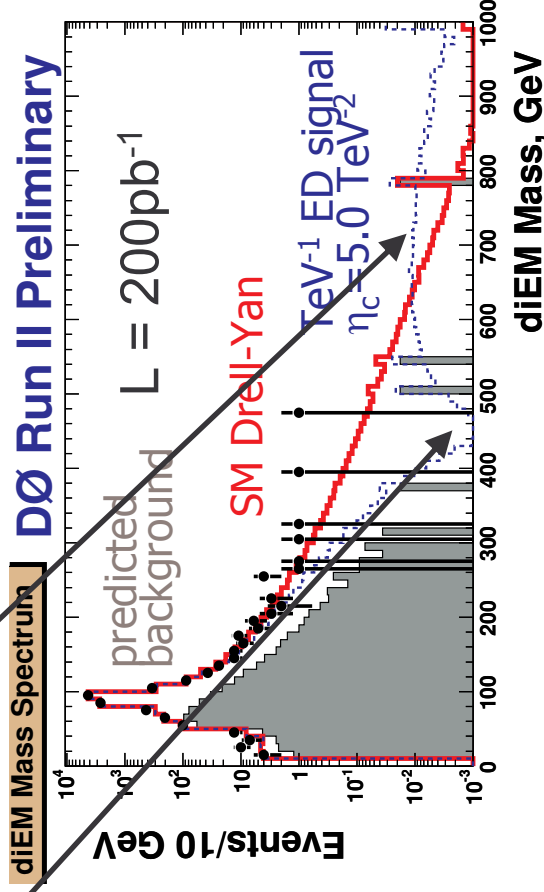
diEM search 200 pb^{-1}

Lower limit on the compactification scale of the longitudinal ED:

$M_C > 1.12 \text{ TeV}$ at 95% C.L.

Better Limit: from precision electroweak data $M_C \geq 4 \text{ GeV}$

World Combined Limit $M_C > 6.8 \text{ TeV}$ at 95% C.L, dominated by LEP2 measurements





TeV⁻¹ ED Discovery Limits



Di-electron states (Z_{KK} decays)

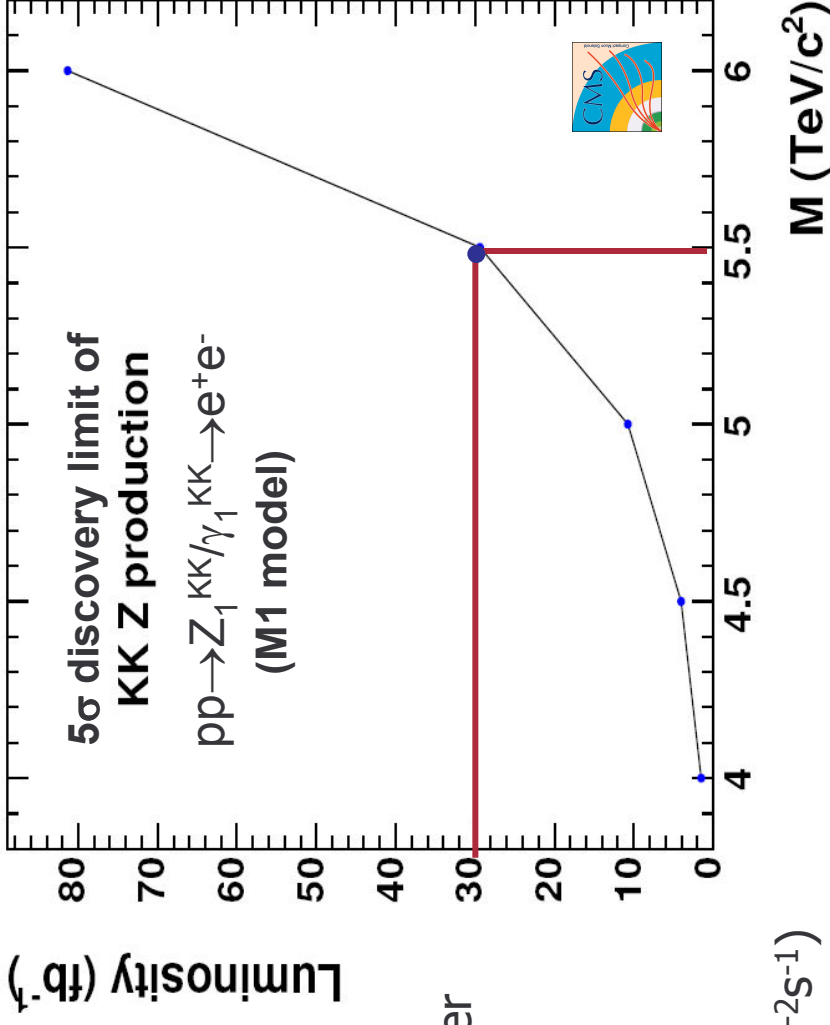
- Two high p_T isolated electrons in the final state
- Bckg: irreducible: Drell-Yan
Also ZZ/WW/ZW/ttabr

- Signal and Bkgd: PYTHIA, CTEQ61M, PHOTOS used for inner bremsstrahlung production

- LO + K=1.30 for signals,
LO + K-factors for bckg.
- Full (GEANT-4) simulation/reco with pile-up at low lum. ($\sim 10^{33} \text{cm}^{-2} \text{s}^{-1}$)

- L1 + HLTrigger cuts

- Theoretical uncert.



With $\mathcal{L}=30/80 \text{ fb}^{-1}$ CMS will be able to detect a peak in the e^+e^- invar. mass distribution if $M_C < 5.5/6 \text{ TeV}$.

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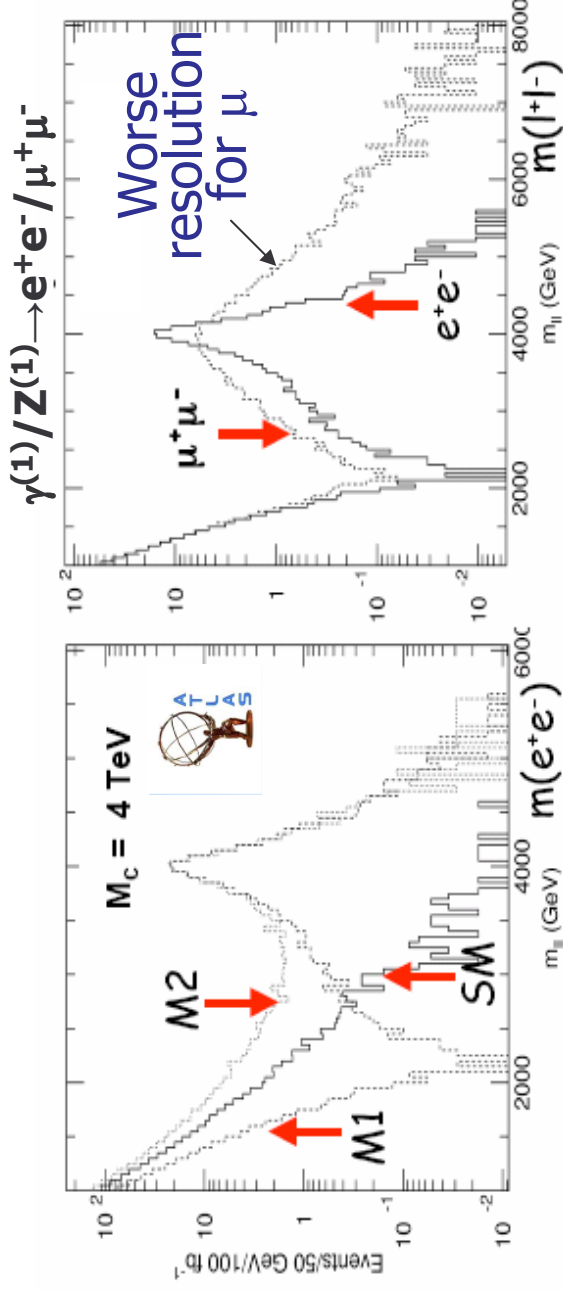
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B. Clerbaux et al.
CMS NOTE 2006/083
CMS PTDR 2006b

TeV⁻¹ ED Discovery Limits

ATLAS expectations for e and μ :

PYTHIA + Fast simu/paramaterized reco + Theor. uncert.



In ee channel experimental resolution is smaller than the natural width of the $Z^{(1)}$, in $\mu\mu$ channel exp. momentum resol. dominates the width

2 TeV e in ATLFAS:

$$\Delta E/E \sim 0.7 \%$$

$$\sim 20\% \text{ for } \mu$$

- Requiring >10 events above a given m_{ll} and $S = (N - N_B)/\sqrt{N_B} > 5$
With 100 fb⁻¹ ATLAS will be able to detect resonance if $M_C (R^{-1}) < 5.8 \text{ TeV} (ee + \mu\mu)$

Even for lowest resonances of M_C (4 TeV), no events would be observed for the $n=2$ resonances of Z and γ at 8 TeV ($M_n = \sqrt{M_0^2 + n^2/R^2}$), which would have been the most striking signature for this kind of model.

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G. Azuelos, G. Polesello
EPJ Direct 10.1140 (2004) 35

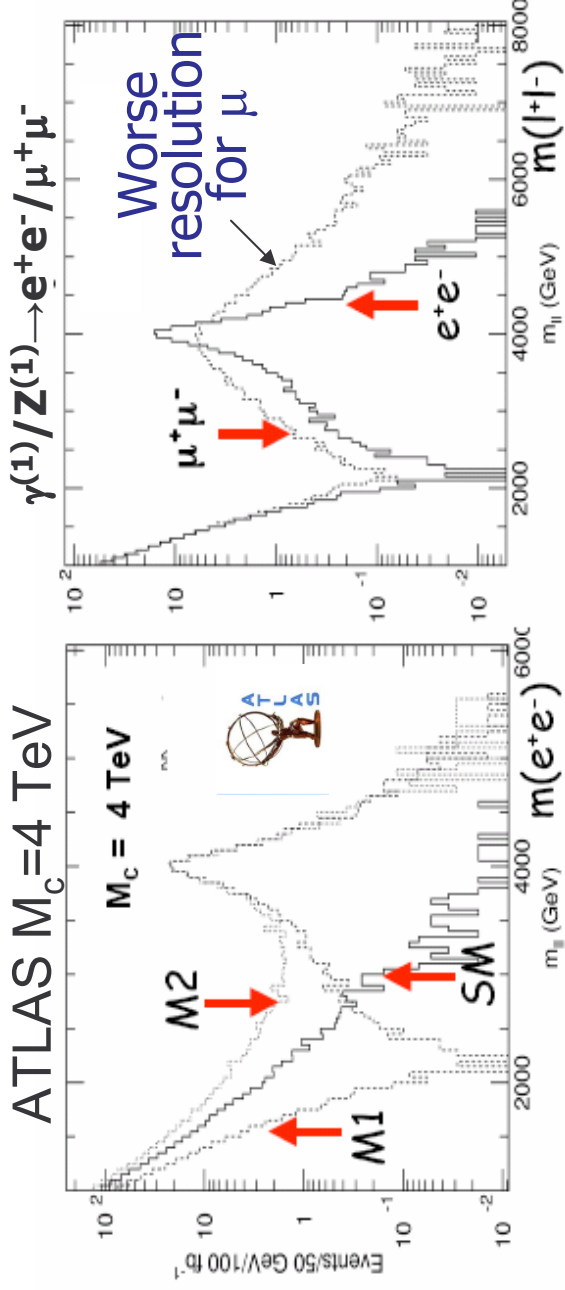
Oct 23-27 2006

G. Azuelos, G. Polesello (Les Houches 2001 Workshop Proceedings), Physics at TeV Colliders, 210-228 (2001)

TeV⁻¹ ED Discovery Limits

ATLAS expectations for e and μ :

PYTHIA + Fast simu/paramaterized reco + Theor. uncert.



In ee channel experimental resolution is smaller than the natural width of the $Z^{(1)}$, in $\mu\mu$ channel exp. momentum dominates the width

2 TeV e in ATLFASST:
 $\Delta E/E \sim 0.7\%$
 $\sim 20\%$ for μ

- Requiring >10 events above a given m_{ll} and $S = (N - N_B)/\sqrt{N_B} > 5$
 With 100 fb⁻¹ ATLAS will be able to detect resonance if $M_C (R^{-1}) < 5.8$ TeV ($ee + \mu\mu$)
- Using a maximum likelihood method to fit the reconstructed distributions describing the kinematics of the event:
With 300 fb⁻¹ can reach 13.5 TeV ($ee + \mu\mu$)

G. Azuelos, G. Polesello
 EPJ Direct 10.1140 (2004)

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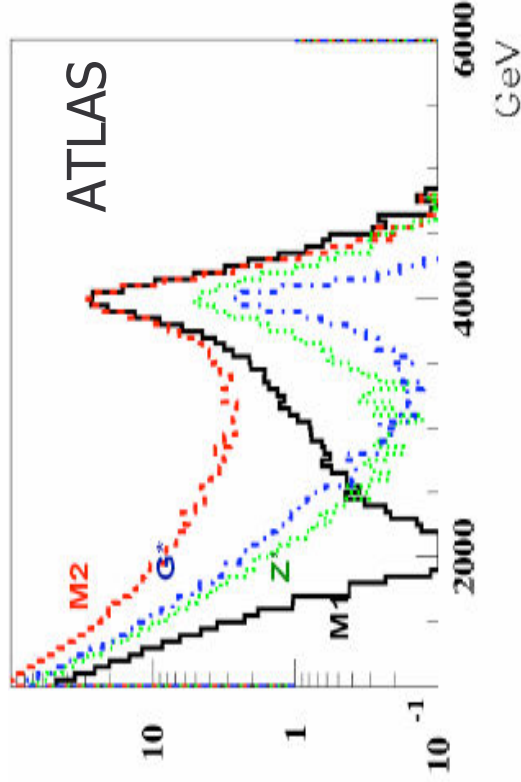
G. Azuelos, G. Polesello (Les Houches 2001 Workshop Proceedings), Physics at TeV Colliders, 210-228 (2001)

Distinguishing $Z^{(1)}$ from Z' , RS G

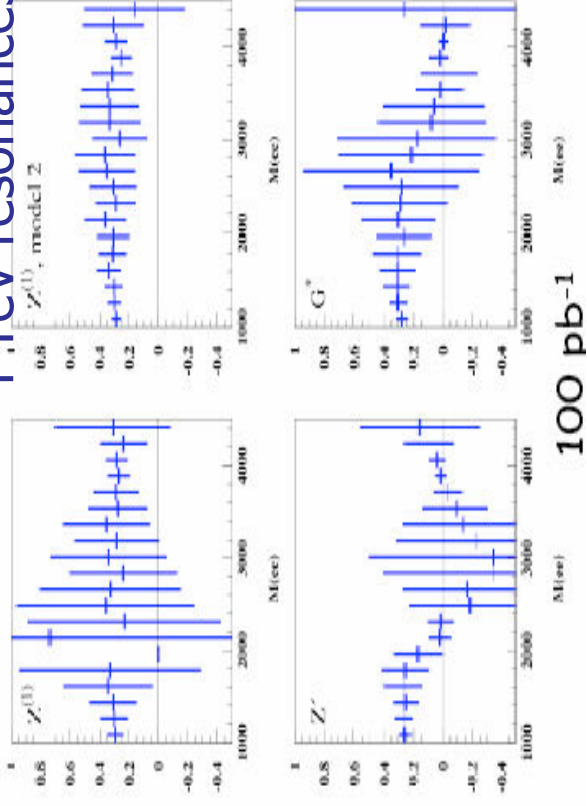
- Spin 1 $Z^{(1)}$ signal can be distinguished from a spin-2 narrow graviton resonance using the angular distribution of its decay products.
- $Z^{(1)}$ can also be distinguished from a Z' with SM-like couplings using the distribution of the forward-backward asymmetry: due to contributions of the higher lying states, the interference terms and the additional $\sqrt{2}$ factor in its coupling to SM fermions.

The $Z^{(1)}$ can be discriminated for masses up to about 5 TeV with $L=300\text{fb}^{-1}$.

$Z^{(1)}$ or Z' or RS Graviton?



Forward-backward asymmetries:
4 TeV resonances



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TeV⁻¹ ED Discovery Limits

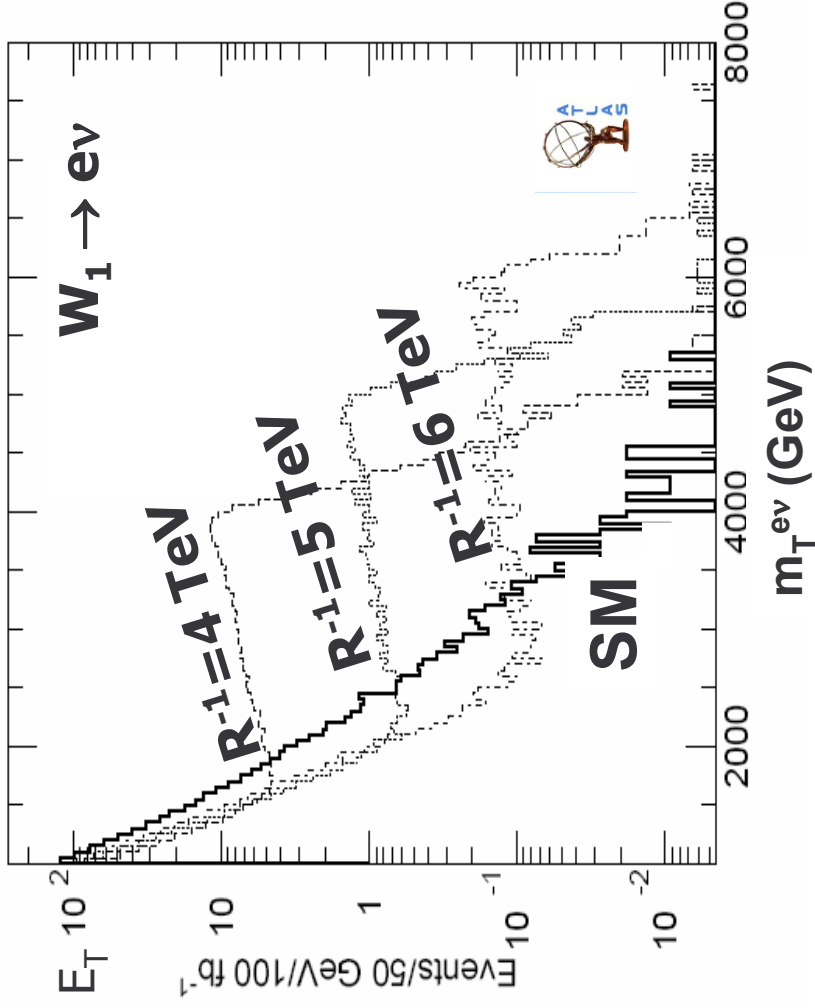


W_{KK} decays

- ☐ Isolated high-p_T lepton + missing E_T
- ☐ Invmass (l,ν) > 1 TeV, veto jets
- ☐ Bkg: irreducible bkgd: W→eν,
Also pairs: WW, WZ, ZZ, ttbar
- ☐ Fast simulation/reco

Sum over 2 lepton flavours

For L=100 fb⁻¹ a **peak** in the lepton-neutrino transverse invariant mass (m_T^{lv}) will be detected if the compactification scale (**M_C= R⁻¹**) is **< 6 TeV**



If a peak is detected, a measurement of the couplings of the boson to the leptons and quarks can be performed for M_C up to ~ 5 TeV.

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G. Polesello, M. Patra
EPJ Direct, ATLAS 2003-023
G. Polesello, M. Patra
EPJ Direct C 32 Sup.2 (2004) pp.55-67



TeV⁻¹ ED Discovery Limits

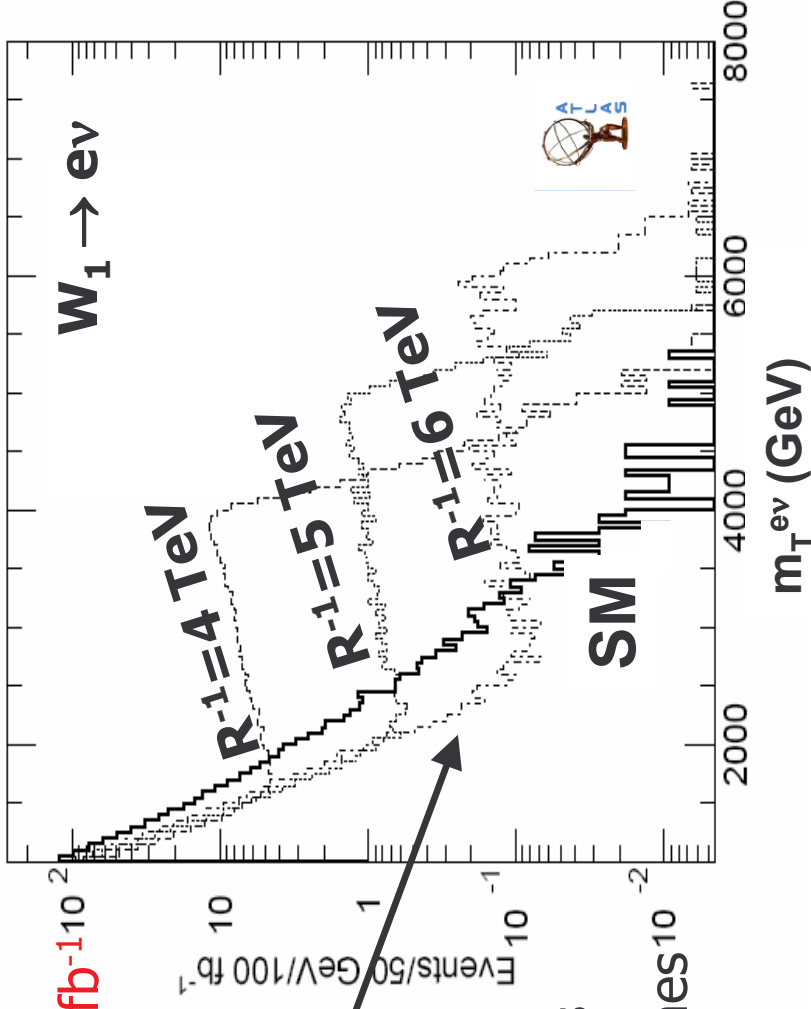


W_{KK} decays

If no signal is observed with **100 fb⁻¹10²** a limit of **M_C > 11.7 TeV** can be obtained from studying the m_T^{ev} distribution **below** the peak:

Here: suppression in σ

- due to -ve interference (M1)
- between SM gauge bosons and the whole tower of KK excitations
- sizable even for M_C above the ones¹⁰ accessible to a direct detection of the mass peak.



- Can't get such a limit with $W \rightarrow \mu\nu$ since momentum spread - can't do optimised fit which uses peak edge

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EPJ Direct, ATLAS 2003-023
G. Polesello, M. Patra
EPJ Direct C 32 Sup.2 (2004) pp.55-67



TeV⁻¹ ED g* Discovery Limits



This is more challenging than Z/W which have leptonic decay modes

Detect KK gluon excitations (g^*) by reconstructing their hadronic decays (no leptonic decays).

Detect g^* by (1) deviation in dijet σ
 (2) analysing its decays into heavy quarks

Coupling of g^* to quarks = $\sqrt{2} * \text{SM couplings}$
 $\Rightarrow g^* \rightarrow$ wide resonances decaying into pairs of quarks



TeV⁻¹ ED g* Discovery Limits



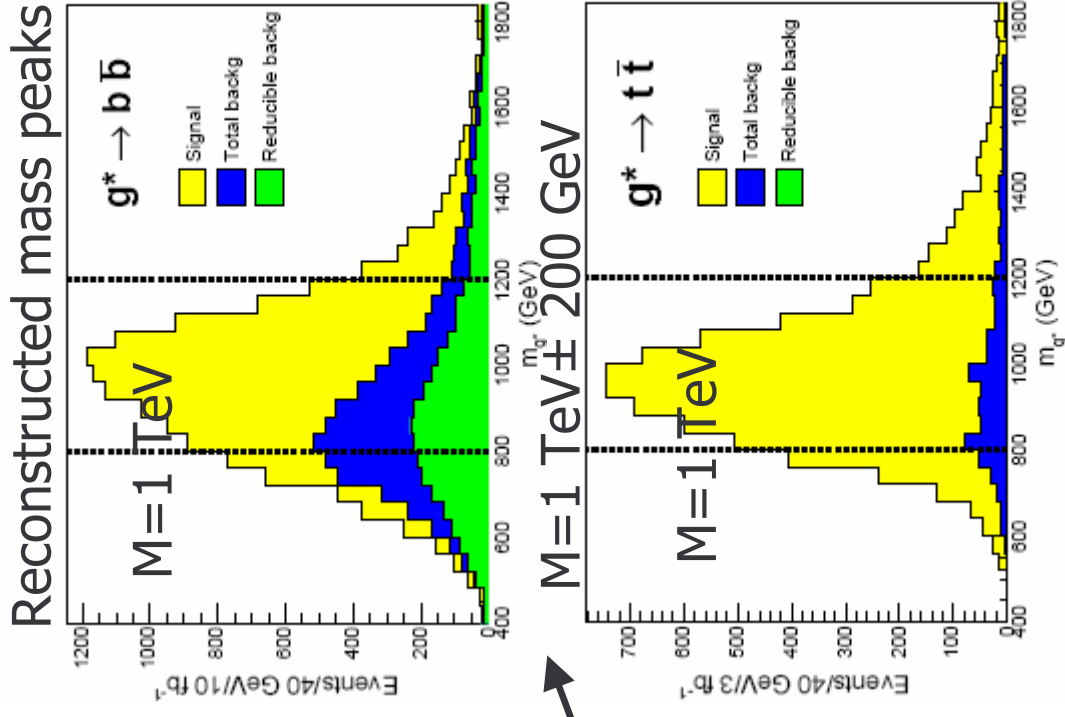
Gluon excitation decays

$$q\bar{q} \rightarrow g^* \rightarrow b\bar{b}, q\bar{q} \rightarrow g^* \rightarrow t\bar{t}$$

- ☐ bbar or ttbar jets
- ☐ For ttbar one t is forced to decay leptonically
- ☐ Bckg: SM continuum bbar, ttbar, 2 jets, W +jets
- ☐ PYTHIA
- ☐ Fast simulation/reco

Width expected to be
 $\Gamma(g^*) = 2 \alpha_s M$ where $M=g^*$ mass
 $\Rightarrow \Gamma(g^*) \sim 200 \text{ GeV}$ for $M=1 \text{ TeV}$

For $M=1 \text{ TeV}$ natural width $\sim$ experimental effects (fragmentation and detector resolution)

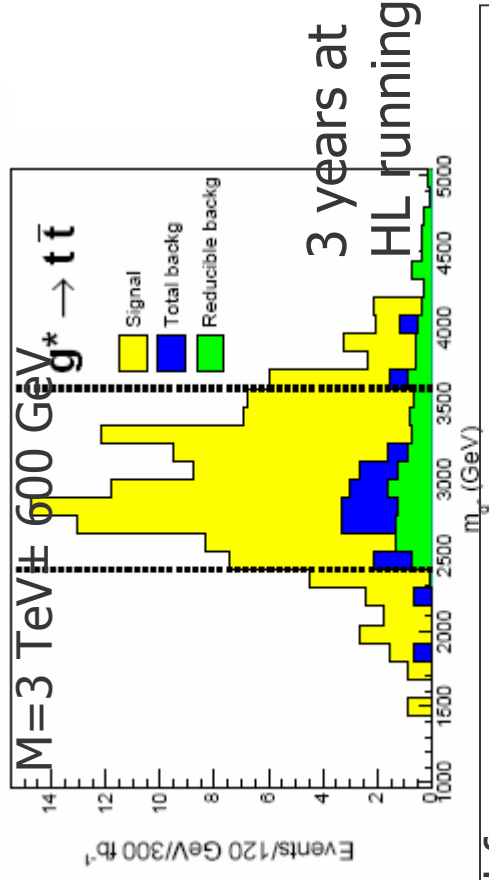
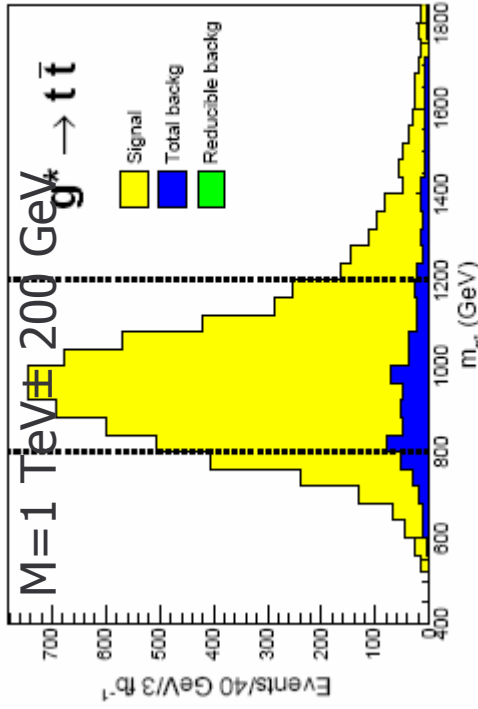
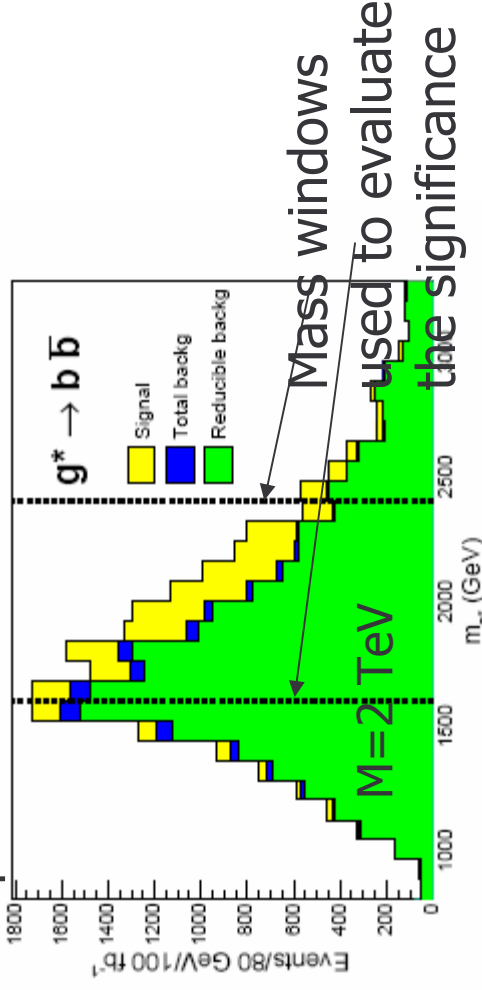
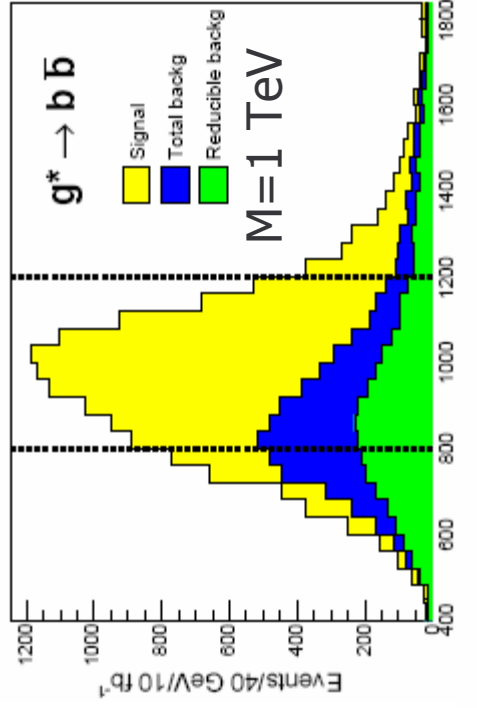




TeV⁻¹ ED g* Discovery Limits



Reconstructed mass peaks



With 300 fb⁻¹ Significance of 5 achieved for:

bbar channel: R⁻¹ = 2.7 TeV

t tbar channel: R⁻¹ = 3.3 TeV

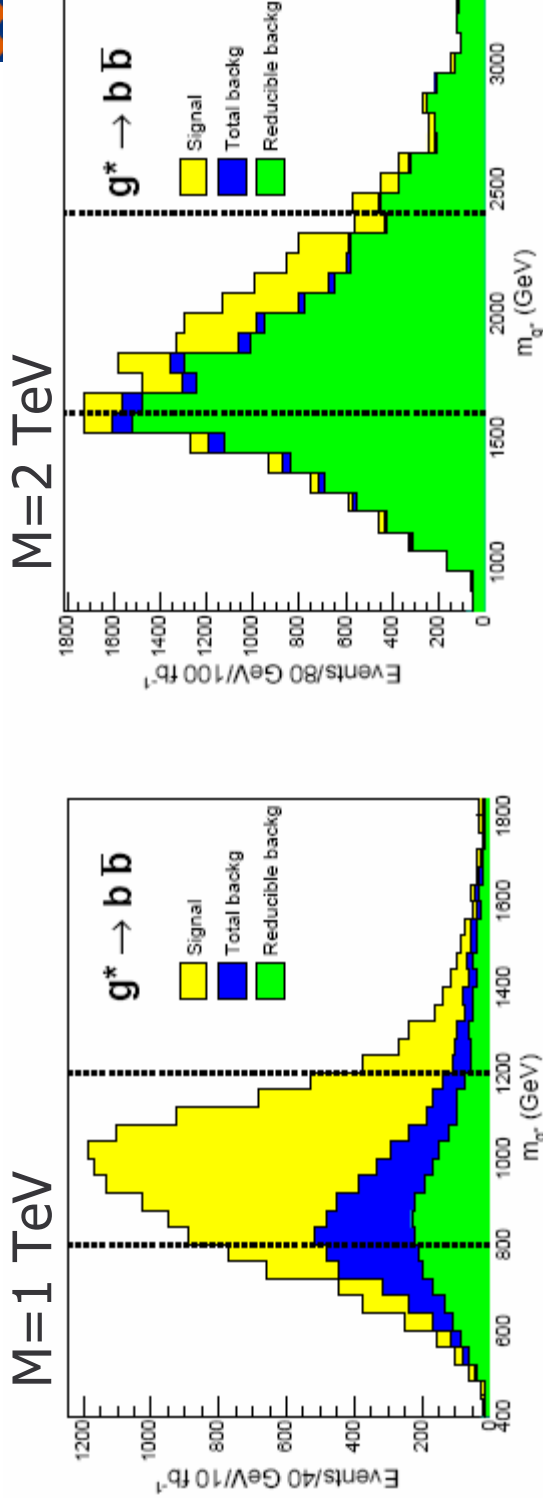
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L. March, E. Ros, B. Salvachua, 42
ATL-PHYS-PUB-2006-002



TeV⁻¹ ED g* Discovery Limits



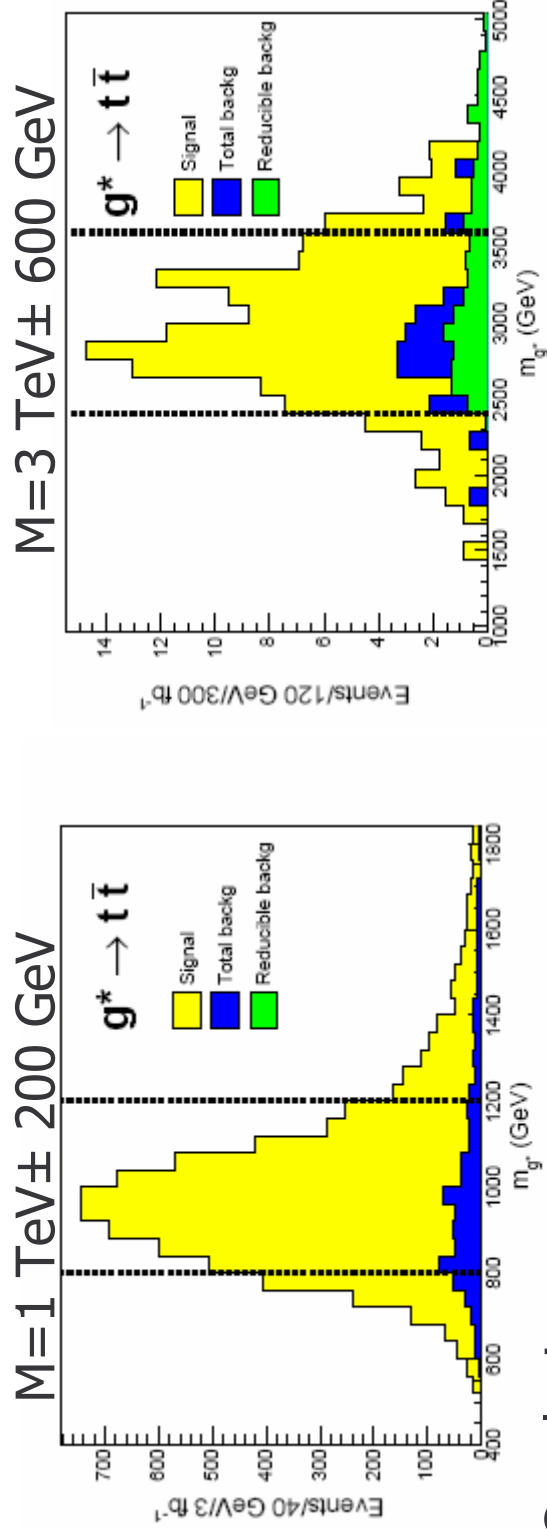
Although with 300 fb⁻¹ Significance of 5 achieved for:

bbar channel: $R^{-1} = 2.7$ TeV

However, it is not in general possible to obtain a mass peak well separated from the bkdg. $\Rightarrow$ it is unlikely that an excess of events in the $g^* \rightarrow b\bar{b}$ channel could be used as evidence of the g^* resonance, since there are large uncertainties in the calculations of the bkdgs. For $M=1$ TeV the peak displacements could be used as evidence for new physics if the b-jet energy scale can be accurately computed.

TeV⁻¹ ED g* Discovery Limits

But in $g^* \rightarrow t\bar{t}$, the bkgd is mainly irreducible and not so large.
 $\Rightarrow$ **g* resonance can be detected** in this decay channel if the $t\bar{t}$ -bar σ can be computed in a reliable way.



Conclusion:

g^* decays into b-quarks are difficult to detect, decays into t-quarks might yield a significant signal for g^* mass below 3.3 TeV.
 This could be used to confirm the presence of g^* in the case that an excess in the dijet σ is observed.

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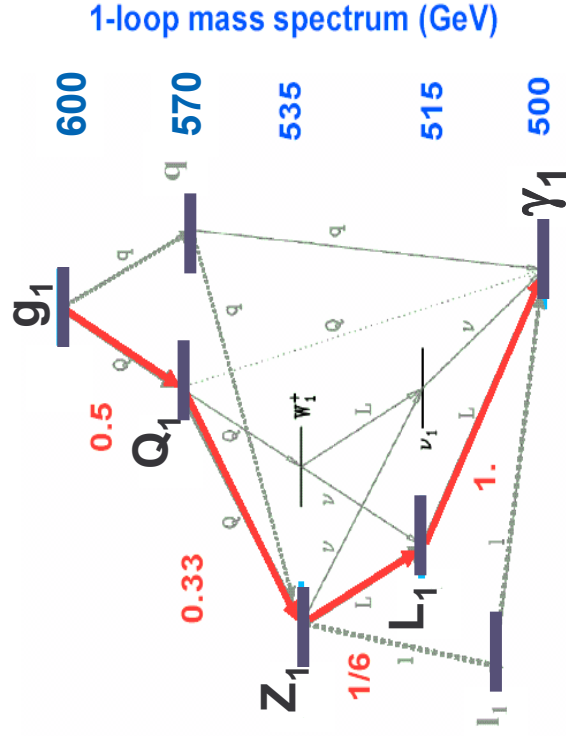
Universal Extra Dimensions



Standard/Minimal UED

- All particles can travel into the bulk, so each SM particle has an infinite tower of KK partners
- Spin of the KK particles is the same as their SM partners
- In minimal UED: 1 ED compactified in an orbifold (S^1/Z_2) of size R
 - KK parity conservation $\rightarrow$ the lightest massive KK particle (LKP) is stable (dark matter candidate).
 - Level one KK states must be pair produced
- Mass degeneration except if radiative corrections included

The model parameters: compactification radius R , cut-off scale Λ , m_h



Thick/Fat brane

- SM brane is endowed with a finite thickness in the ED
- Gravity-matter interactions break KK number conservation:
 - 1st level KK states decay to G+SM.
 - If radiative corrections $\rightarrow$ mass degeneracy is broken and γ and leptons are produced.

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Present Constraints on UED



Bounds to the compactification scale:

- Precision EWK data measurements set a lower bound of $R^{-1} > 300 \text{ GeV}$

Phys. Rev. D64, 035002 (2001) Appelquist, Cheung, Dobrescu

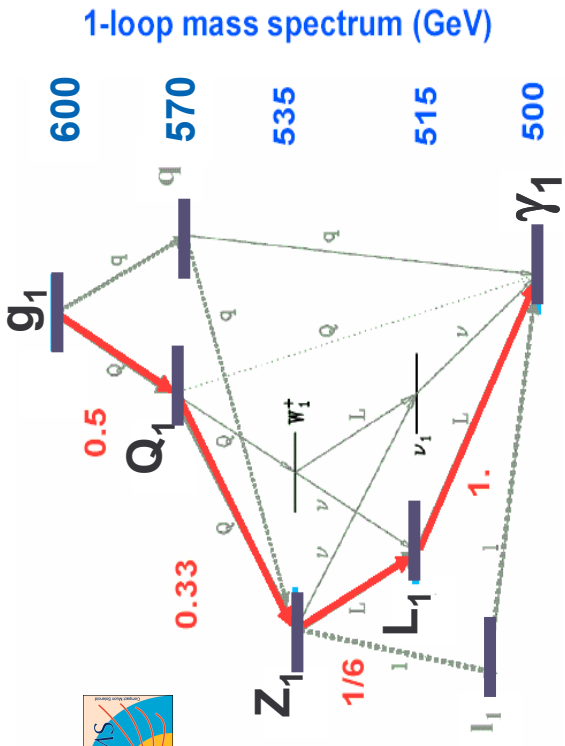
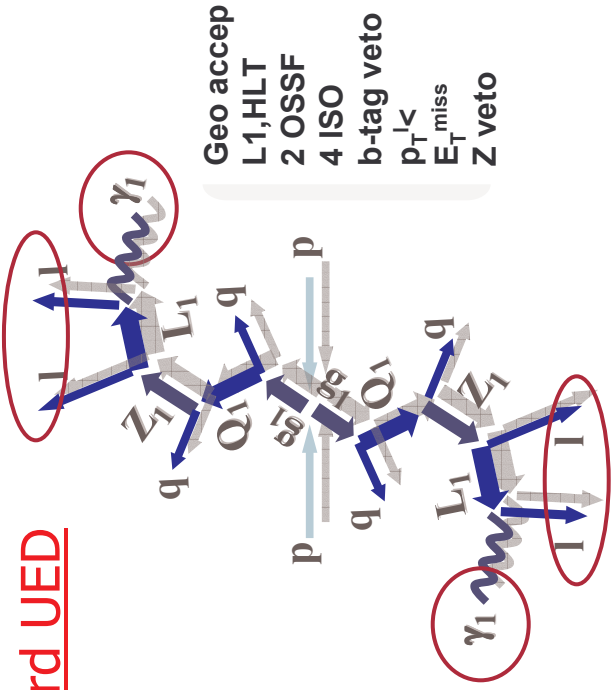
- Dark matter constraints imply that $600 < R^{-1} < 1050 \text{ GeV}$
Servant, Tait, Nucl. Phys. B650,391 (2003)



UED Discovery Limit



Standard UED



$$pp \rightarrow g_1 g_1 \rightarrow 4l + 4q + 2LKP \rightarrow 4l + 4 \text{ jets} + P_T$$

$$pp \rightarrow g_1 Q_1 \rightarrow 4l + 3q + 2LKP \rightarrow 4l + 3 \text{ jets} + P_T$$

$$pp \rightarrow Q_1 Q_1 \rightarrow 4l + 2q + 2LKP \rightarrow 4l + 2 \text{ jets} + P_T$$

Final State:

4 low- p_T isolated leptons (2 pairs of opposite sign, same flavour leptons)
+ n jets + missing E_T (from 2 undetected γ_1)

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Standard UED

UED Discovery Limit



$$pp \rightarrow g_1 g_1 \rightarrow 4l + 4q + 2LKP \rightarrow 4l + 4jets + P_T$$

$$pp \rightarrow g_1 Q_1 \rightarrow 4l + 3q + 2LKP \rightarrow 4l + 3jets + P_T$$

$$pp \rightarrow Q_1 Q_1 \rightarrow 4l + 2q + 2LKP \rightarrow 4l + 2jets + P_T$$

□ 4 leptons in the final state + missing p_T

□ Irreducible Bckg: $t\bar{t}$ + n jets

($n = 0, 1, 2$), 4 b-quarks, ZZ, Zbbar

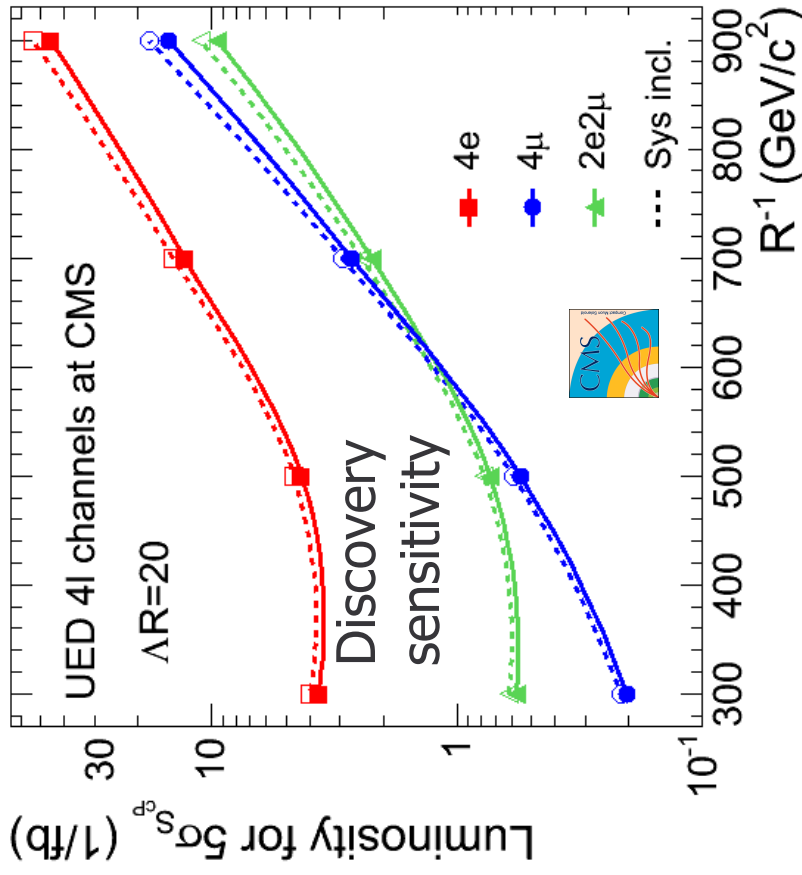
□ To improve bkgd rejection over signal: apply b-tagging and Z-tagging vetoes

□ CompHEP for signal and CompHEP, PYTHIA, Alpgen for bckg. with CTEQ5L

□ Full simulation/reco + L1 + HLT(rigger) cuts

□ Theoretical and experimental uncert.

Studied for low lum run $\sim 2 \times 10^{33} \text{cm}^{-2} \text{s}^{-1}$



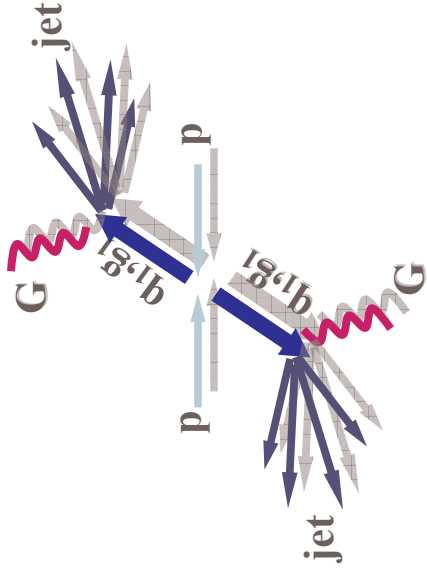


UED Discovery Limit



Thick brane in UED with TeV^{-1} size

$$pp \rightarrow g_1 g_1 / q_1 q_1 \rightarrow 2 \text{ jets} + E_T$$

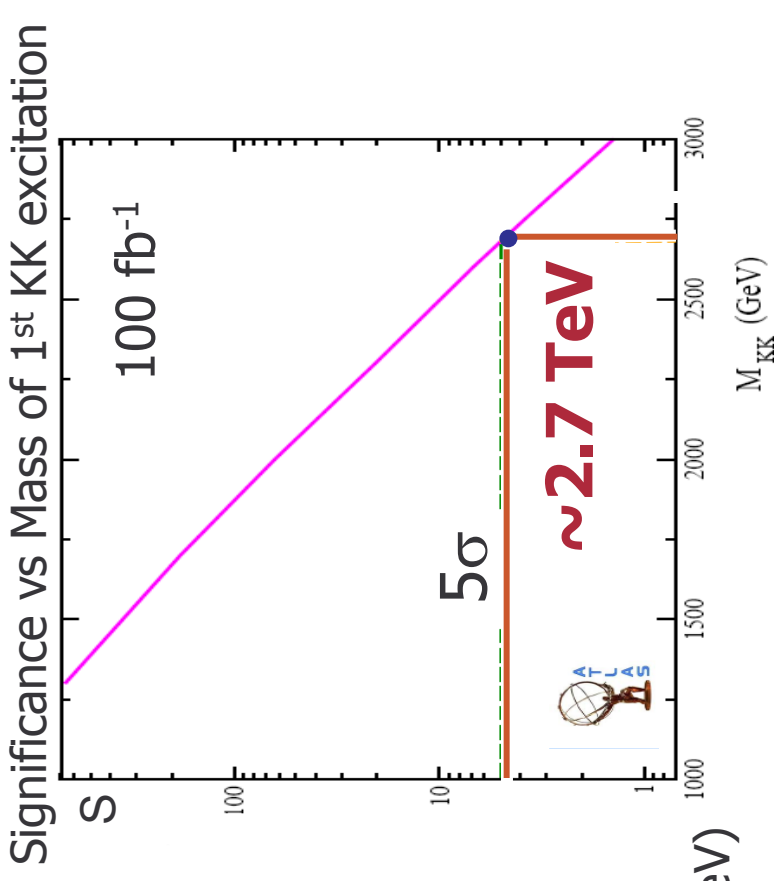


□ 2 back-to back jets + missing $E_T (> 775 \text{ GeV})$

□ Irreducible Bckg: $Z(\rightarrow \nu\nu)$ jj , $W(\rightarrow l\nu)$ jj

□ PYTHIA/CTEQ5L + SHERPA for bckgr.

□ Fast simulation/reco



5σ discovery possible at
ATLAS with 100 fb⁻¹ if first KK
excitation mass < 2.7 TeV

LHC Start-up Expectations



Model	Mass reach	Integrated Luminosity (fb ⁻¹)	Systematic uncertainties
ADD Direct G_{KK}	$M_D \sim 1.5\text{-}1.0\text{ TeV}, n = 3\text{-}6$	1	Theor.
ADD Virtual G_{KK}	$M_D \sim 4.3\text{ - }3\text{ TeV}, n = 3\text{-}6$	0.1	Theor.+Exp.
	$M_D \sim 5\text{ - }4\text{ TeV}, n = 3\text{-}6$	1	
RS1 di-electrons di-photons di-muons di-jets			
	$M_{G1} \sim 1.35\text{- }3.3\text{ TeV}, c=0.01\text{-}0.1$	10	Theor.+Exp. (only stat. for di-jets)
	$M_{G1} \sim 1.31\text{- }3.47\text{ TeV}, c=0.01\text{-}0.1$	10	
	$M_{G1} \sim 0.8\text{- }2.3\text{ TeV}, c=0.01\text{-}0.1$	1	
	$M_{G1} \sim 0.7\text{- }0.8\text{ TeV}, c=0.1$	0.1	
TeV-1 ($Z_{KK}^{(1)}$)	$M_{z1} < 5\text{ TeV}$	1	Theor.
UED 4 leptons	$R^{-1} \sim 600\text{ GeV}$	1.0	Theor.+Exp.
Thick brane	$R^{-1} = 1.3\text{ TeV}$	6 pb ⁻¹	

Conclusions



The discovery potential of both experiments makes it possible to investigate if extra dimensions really exist within various ED scenarios at a few TeV scale:

- Large Extra-Dimensions (ADD model)
- Randall-Sundrum (RS1)
- TeV⁻¹ Extra dimension Model
- Universal Extra Dimensions

Reaches in different channels depend on the performance of detector systems: proper energy, momentum, angular reconstruction for high-energy leptons and jets, E_t measurement, b-tagging and identification of prompt photons

New results have been predicted with data collected in the start-up LHC weeks (integrated luminosity $< 1 \text{ fb}^{-1}$)

For methods to disentangle new physics from SM physics, see lectures by M. Mangano.
For Gravitaton/Black Holes see talk by: B. Webber.

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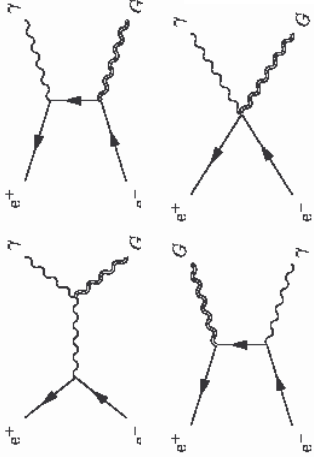
The End!

Backup slides...

Present ADD Emission Limits

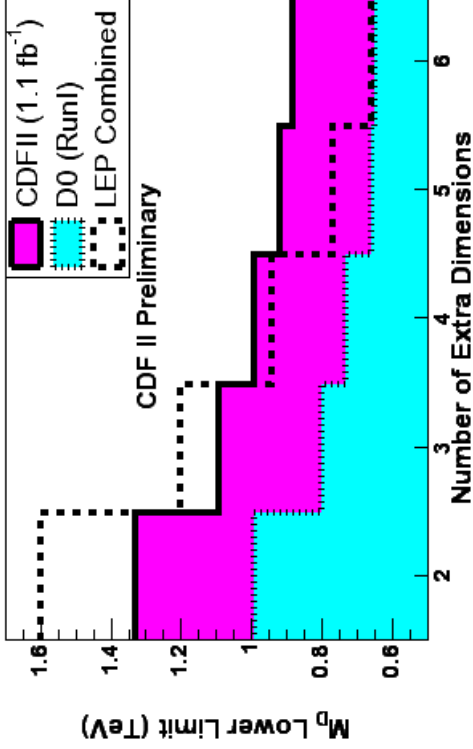


LEP and Tevatron results are complementary



For $n < 4$: LEP limits best
 $\gamma + \text{ME}_T$

For $n > 4$: CDF limits best
 $\text{jet} + \text{ME}_T$



$\gamma + \text{ME}_T$ at LEP is cleaner & has lower backgrounds than $\text{jet} + \text{ME}_T$ (Tevatron), so the precision of their experiments wins out for lower values of n

Tevatron better at large values of n , because of the higher energy, which is a bigger effect at larger values of n .

$\sigma \propto$ total number of possible modes in the KK tower N_{KK}
 $\sigma \propto N_{KK} \propto \sqrt{(s-hat)}$
 But this is true for each ED, so $\sigma \propto (\sqrt{(s-hat)})^n$

$\Rightarrow$ the difference in energy is a bigger effect for $n=6$ than $n=2$

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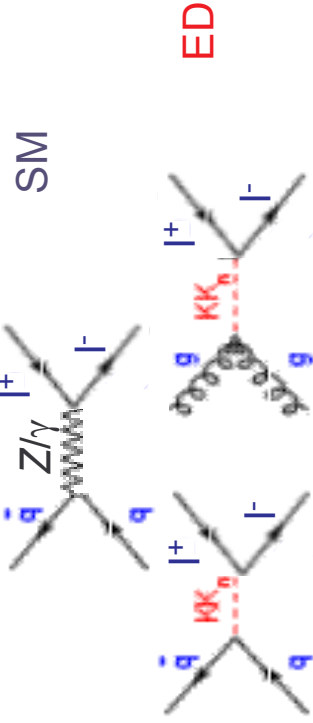
ADD: G Exchange?



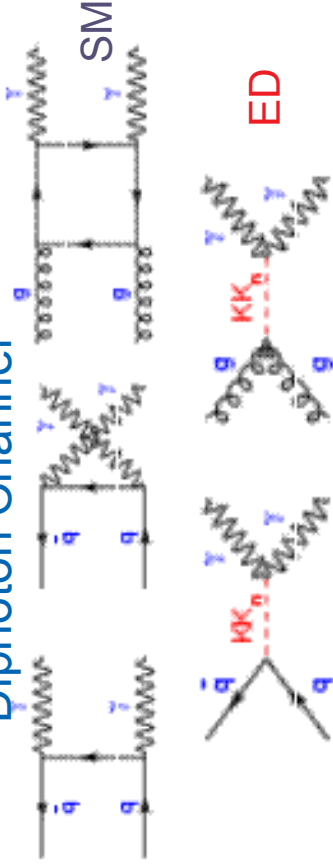
Search for spin-2 broad σ change

⇒ study deviations in invariant mass & angular distribution from SM processes

Dilepton Channel



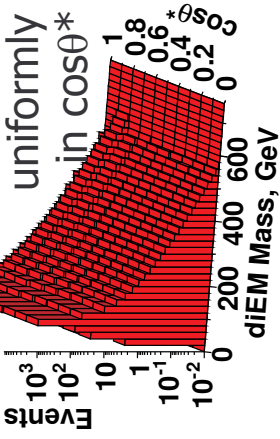
Diphoton Channel



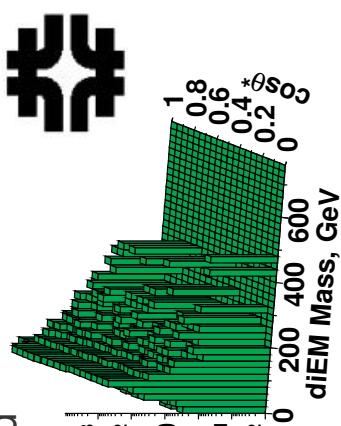
SM Prediction

DØ Run II Preliminary

SM events expected to be distributed uniformly in $\cos\theta^*$

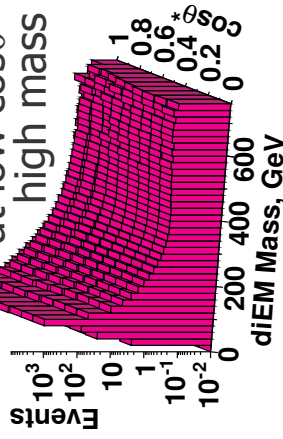


Data



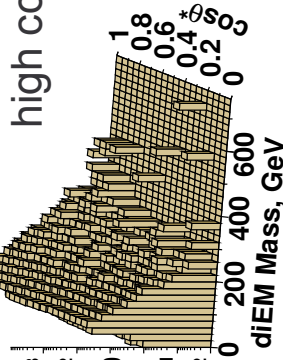
ED Signal

Signal events are accumulated at low $\cos\theta^*$ & high mass



QCD Background

low mass, high $\cos\theta^*$



Spin-1/Spin-2 Discrimination



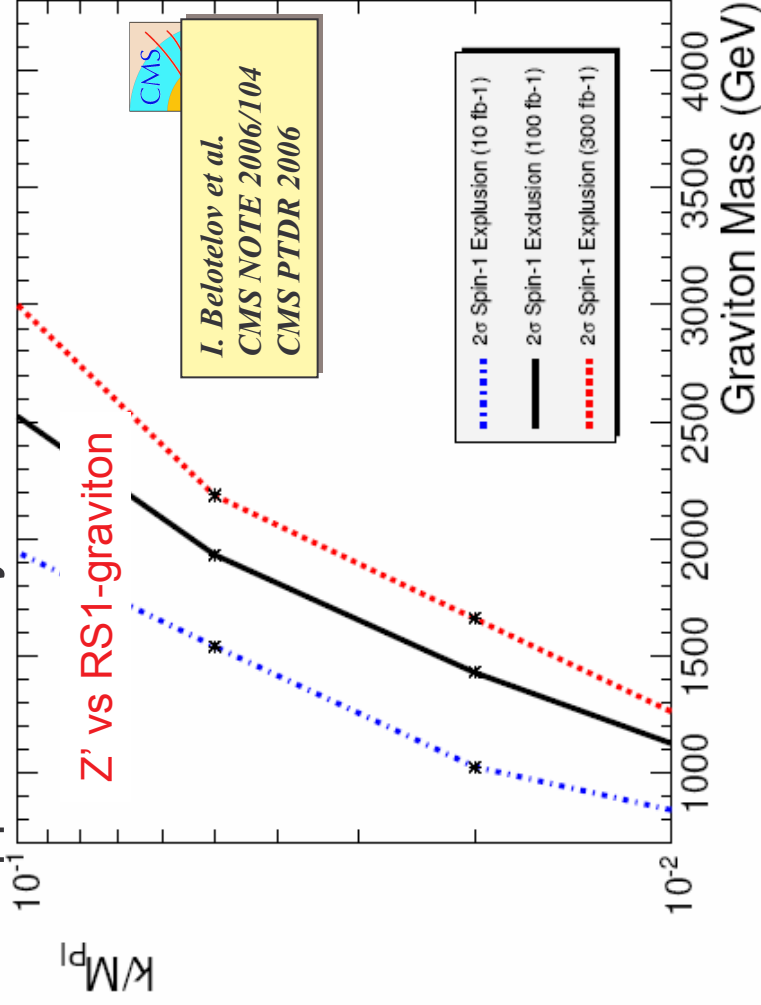
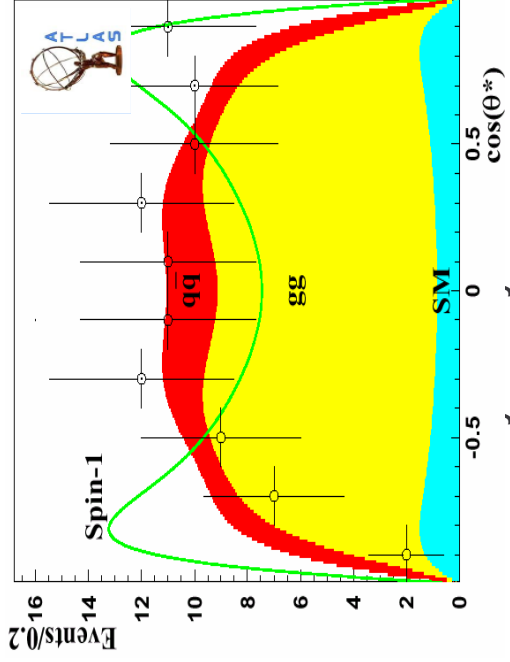
Spin-1 States: Z' from extended gauge models, Z_{KK}

Spin-2 States: RS1-graviton

Method: unbinned likelihood ratio statistics incorporating the angles in of the decay products the Collins-Soper frame (R.Cousins et al. JHEP11 (2005) 046). The statistical technique has been applied to fully simu/reco events.

Angular distributions

- $qq \rightarrow G \rightarrow ff: 1 - 3 \cos^2 \theta + 4 \cos^4 \theta$
- $gg \rightarrow G \rightarrow ff: 1 - \cos^4 \theta$
- $qq \rightarrow G \rightarrow VV: 1 - \cos^4 \theta$
- $gg \rightarrow G \rightarrow VV: 1 + 6 \cos^2 \theta + \cos^4 \theta$
- DY background: $1 + \cos^2 \theta$



Older results on spin discrimination from ATLAS can be

found
Getting I

B.C. Allanach et al, JHEP 09 (2000) 019; ATL-PHYS-2000-029

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