Statistical Methods for HEP Lecture 3: Discovery and Limits



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# Outline

Lecture 1: Introduction and basic formalism Probability, statistical tests, parameter estimation.

Lecture 2: Multivariate methods General considerations Example of a classifier: Boosted Decision Trees

 Lecture 3: Discovery and Exclusion limits Quantifying significance and sensitivity Systematic uncertainties (nuisance parameters)

### A simple example

For each event we measure two variables,  $x = (x_1, x_2)$ . Suppose that for background events (hypothesis  $H_0$ ),

$$f(\mathbf{x}|H_0) = \frac{1}{\xi_1} e^{-x_1/\xi_1} \frac{1}{\xi_2} e^{-x_2/\xi_2}$$

and for a certain signal model (hypothesis  $H_1$ ) they follow

$$f(\mathbf{x}|H_1) = C \frac{1}{\sqrt{2\pi\sigma_1}} e^{-(x_1 - \mu_1)^2 / 2\sigma_1^2} \frac{1}{\sqrt{2\pi\sigma_2}} e^{-(x_2 - \mu_2)^2 / 2\sigma_2^2}$$

where  $x_1, x_2 \ge 0$  and *C* is a normalization constant.

#### Likelihood ratio as test statistic

In a real-world problem we usually wouldn't have the pdfs  $f(\mathbf{x}|H_0)$  and  $f(\mathbf{x}|H_1)$ , so we wouldn't be able to evaluate the likelihood ratio

$$t(\mathbf{x}) = \frac{f(\mathbf{x}|H_1)}{f(\mathbf{x}|H_0)}$$

for a given observed *x*, hence the need for multivariate methods to approximate this with some other function.

But in this example we can find contours of constant likelihood ratio such as:



#### Event selection using the LR

Using Monte Carlo, we can find the distribution of the likelihood ratio or equivalently of

$$q = \left(\frac{x_1 - \mu_1}{\sigma_1}\right)^2 + \left(\frac{x_2 - \mu_2}{\sigma_2}\right)^2 - \frac{2x_1}{\xi_1} - \frac{2x_2}{\xi_2} = -2\ln t(\mathbf{x}) + C$$



From the Neyman-Pearson lemma we know that by cutting on this variable we would select a signal sample with the highest signal efficiency (test power) for a given background efficiency.

### Search for the signal process

But what if the signal process is not known to exist and we want to search for it. The relevant hypotheses are therefore

- $H_0$ : all events are of the background type
- $H_1$ : the events are a mixture of signal and background

Rejecting  $H_0$  with Z > 5 constitutes "discovering" new physics.

Suppose that for a given integrated luminosity, the expected number of signal events is *s*, and for background *b*.

The observed number of events *n* will follow a Poisson distribution:

$$P(n|b) = \frac{b^n}{n!}e^{-b} \qquad P(n|s+b) = \frac{(s+b)^n}{n!}e^{-(s+b)}$$

### Likelihoods for full experiment

We observe *n* events, and thus measure *n* instances of  $x = (x_1, x_2)$ . The likelihood function for the entire experiment assuming the background-only hypothesis (*H*<sub>0</sub>) is

$$L_b = \frac{b^n}{n!} e^{-b} \prod_{i=1}^n f(\mathbf{x}_i | \mathbf{b})$$

and for the "signal plus background" hypothesis  $(H_1)$  it is

$$L_{s+b} = \frac{(s+b)^n}{n!} e^{-(s+b)} \prod_{i=1}^n \left( \pi_{s} f(\mathbf{x}_i | s) + \pi_{b} f(\mathbf{x}_i | b) \right)$$

where  $\pi_s$  and  $\pi_b$  are the (prior) probabilities for an event to be signal or background, respectively.

## Likelihood ratio for full experiment

We can define a test statistic Q monotonic in the likelihood ratio as

$$Q = -2\ln\frac{L_{s+b}}{L_b} = -s + \sum_{i=1}^n \ln\left(1 + \frac{s}{b}\frac{f(\mathbf{x}_i|\mathbf{s})}{f(\mathbf{x}_i|\mathbf{b})}\right)$$

To compute *p*-values for the b and s+b hypotheses given an observed value of Q we need the distributions f(Q|b) and f(Q|s+b).

Note that the term -s in front is a constant and can be dropped. The rest is a sum of contributions for each event, and each term in the sum has the same distribution.

Can exploit this to relate distribution of *Q* to that of single event terms using (Fast) Fourier Transforms (Hu and Nielsen, physics/9906010).

# Distribution of Q



### Systematic uncertainties

Up to now we assumed all parameters were known exactly.

In practice they have some (systematic) uncertainty.

Suppose e.g. uncertainty in expected number of background events b is characterized by a (Bayesian) pdf  $\pi(b)$ .

Maybe take a Gaussian, i.e.,

$$\pi(b) = \frac{1}{\sqrt{2\pi\sigma_b}} e^{-(b-b_0)^2/2\sigma_b^2}$$

where  $b_0$  is the nominal (measured) value and  $\sigma_b$  is the estimated uncertainty.

In fact for many systematics a Gaussian pdf is hard to defend – more on this later.

## Distribution of Q with systematics

To get the desired *p*-values we need the pdf f(Q), but this depends on *b*, which we don't know exactly.

But we can obtain the Bayesian model average:

$$f(Q) = \int f(Q|b)\pi(b) \, db$$

With Monte Carlo, sample *b* from  $\pi(b)$ , then use this to generate Q from f(Q|b), i.e., a new value of *b* is used to generate the data for every simulation of the experiment.

This broadens the distributions of Q and thus increases the *p*-value (decreases significance Z) for a given  $Q_{obs}$ .

# Distribution of Q with systematics (2)





### Example: ALEPH Higgs search

*p*-value  $(1 - CL_b)$  of background only hypothesis versus tested Higgs mass measured by ALEPH Experiment



### Example: LEP Higgs search

Not seen by the other LEP experiments. Combined analysis gives *p*-value of background-only hypothesis of 0.09 for  $m_{\rm H} = 115$  GeV.



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# Using the likelihood ratio $L(s)/L(\hat{s})$

Instead of the likelihood ratio  $L_{s+b}/L_b$ , suppose we use as a test statistic

$$\lambda(s) = \frac{L(s)}{L(\hat{s})}$$
maximizes  $L(s)$ 

Intuitively this is a good measure of the level of agreement between the data and the hypothesized value of *s*.

low  $\lambda$ : poor agreement high  $\lambda$  : good agreement  $0 \le \lambda \le 1$ 

# $L(s)/L(\hat{s})$ for counting experiment

Consider an experiment where we only count *n* events with  $n \sim \text{Poisson}(s + b)$ . Then  $\hat{s} = n - b$ .

To establish discovery of signal we test the hypothesis s = 0 using

$$\ln \lambda(0) = n \ln(b) - b - n \ln n + n$$

whereas previously we had used

$$\ln \frac{L_{s+b}}{L_b} = n \ln \left(1 + \frac{s}{b}\right) - s$$

which is monotonic in n and thus equivalent to using n as the test statistic.

# $L(s)/L(\hat{s})$ for counting experiment (2)

But if we only consider the possibility of signal being present when n > b, then in this range  $\lambda(0)$  is also monotonic in n, so both likelihood ratios lead to the same test.



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# $L(s)/L(\hat{s})$ for general experiment

If we do not simply count events but also measure for each some set of numbers, then the two likelihood ratios do not necessarily give equivalent tests, but in practice will be very close.

 $\lambda(s)$  has the important advantage that for a sufficiently large event sample, its distribution approaches a well defined form (Wilks' Theorem).

In practice the approach to the asymptotic form is rapid and one obtains a good approximation even for relatively small data samples (but need to check with MC).

This remains true even when we have adjustable nuisance parameters in the problem, i.e., parameters that are needed for a correct description of the data but are otherwise not of interest (key to dealing with systematic uncertainties).

#### Prototype search analysis

*Expected Performance of the ATLAS Experiment: Detector, Trigger and Physics*, arXiv:0901.0512, CERN-OPEN-2008-20.

Search for signal in a region of phase space; result is histogram of some variable *x* giving numbers:

$$\mathbf{n} = (n_1, \ldots, n_N)$$

Assume the  $n_i$  are Poisson distributed with expectation values

$$E[n_i] = \mu s_i + b_i$$
strength parameter

where

$$s_{i} = s_{\text{tot}} \int_{\text{bin } i} f_{s}(x; \boldsymbol{\theta}_{s}) \, dx \,, \quad b_{i} = b_{\text{tot}} \int_{\text{bin } i} f_{b}(x; \boldsymbol{\theta}_{b}) \, dx \,.$$
signal
background

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# Prototype analysis (II)

Often also have a subsidiary measurement that constrains some of the background and/or shape parameters:

$$\mathbf{m} = (m_1, \ldots, m_M)$$

Assume the  $m_i$  are Poisson distributed with expectation values

$$E[m_i] = u_i(\boldsymbol{\theta})$$
  
nuisance parameters ( $\boldsymbol{\theta}_{s}, \boldsymbol{\theta}_{b}, b_{tot}$ )

Likelihood function is

$$L(\mu, \theta) = \prod_{j=1}^{N} \frac{(\mu s_j + b_j)^{n_j}}{n_j!} e^{-(\mu s_j + b_j)} \prod_{k=1}^{M} \frac{u_k^{m_k}}{m_k!} e^{-u_k}$$

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# The profile likelihood ratio

Base significance test on the profile likelihood ratio:



The likelihood ratio of point hypotheses gives optimum test (Neyman-Pearson lemma).

The profile LR hould be near-optimal in present analysis with variable  $\mu$  and nuisance parameters  $\theta$ .

### Test statistic for discovery

Try to reject background-only ( $\mu = 0$ ) hypothesis using

$$q_0 = \begin{cases} -2\ln\lambda(0) & \hat{\mu} \ge 0\\ 0 & \hat{\mu} < 0 \end{cases}$$

i.e. only regard upward fluctuation of data as evidence against the background-only hypothesis.

### *p*-value for discovery

Large  $q_0$  means increasing incompatibility between the data and hypothesis, therefore *p*-value for an observed  $q_{0,obs}$  is

$$p_0 = \int_{q_{0,\text{obs}}}^{\infty} f(q_0|0) \, dq_0$$

will get formula for this later



From *p*-value get equivalent significance,

$$Z = \Phi^{-1}(1-p)$$

# Expected (or median) significance / sensitivity

When planning the experiment, we want to quantify how sensitive we are to a potential discovery, e.g., by given median significance assuming some nonzero strength parameter  $\mu'$ .



So for *p*-value, need  $f(q_0|0)$ , for sensitivity, will need  $f(q_0|\mu')$ ,

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### Wald approximation for profile likelihood ratio

Glen Cowan, Kyle Cranmer, Eilam Gross and Ofer Vitells, Using the Profile Likelihood in Searches for New Physics, arXiv:1007:1727

To find *p*-values, we need:  $f(q_0|0), f(q_\mu|\mu)$ 

For median significance under alternative, need:  $f(q_{\mu}|\mu')$ 

Use approximation due to Wald (1943)

$$-2\ln\lambda(\mu) = \frac{(\mu - \hat{\mu})^2}{\sigma^2} + \mathcal{O}(1/\sqrt{N})$$
$$\hat{\mu} \sim \text{Gaussian}(\mu', \sigma) \qquad \qquad \text{sample size}$$

i.e., 
$$E[\hat{\mu}] = \mu'$$

$$\begin{split} \sigma \text{ from covariance matrix } V, \text{ use, e.g.,} \\ V^{-1} &= -E\left[\frac{\partial^2 \ln L}{\partial \theta_i \partial \theta_j}\right] \end{split}$$

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#### Noncentral chi-square for $-2\ln\lambda(\mu)$

If we can neglect the  $O(1/\sqrt{N})$  term,  $-2\ln\lambda(\mu)$  follows a noncentral chi-square distribution for one degree of freedom with noncentrality parameter

$$\Lambda = \frac{(\mu - \mu')^2}{\sigma^2}$$

As a special case, if  $\mu' = \mu$  then  $\Lambda = 0$  and  $-2\ln\lambda(\mu)$  follows a chi-square distribution for one degree of freedom (Wilks).

# Distribution of $q_0$

Assuming the Wald approximation, we can write down the full distribution of  $q_0$  as

$$f(q_0|\mu') = \Phi\left(\frac{\mu'}{\sigma}\right)\delta(q_0) + \frac{1}{2}\frac{1}{\sqrt{2\pi}}\frac{1}{\sqrt{q_0}}\exp\left[-\frac{1}{2}\left(\sqrt{q_0} - \frac{\mu'}{\sigma}\right)^2\right]$$

The special case  $\mu' = 0$  is a "half chi-square" distribution:

$$f(q_0|0) = \frac{1}{2}\delta(q_0) + \frac{1}{2}\frac{1}{\sqrt{2\pi}}\frac{1}{\sqrt{q_0}}e^{-q_0/2}$$

Cumulative distribution of  $q_0$ , significance

From the pdf, the cumulative distribution of  $q_0$  is found to be

$$F(q_0|\mu') = \Phi\left(\sqrt{q_0} - \frac{\mu'}{\sigma}\right)$$

The special case  $\mu' = 0$  is

$$F(q_0|0) = \Phi\left(\sqrt{q_0}\right)$$

The *p*-value of the  $\mu = 0$  hypothesis is

$$p_0 = 1 - F(q_0|0)$$

Therefore the discovery significance Z is simply

$$Z = \Phi^{-1}(1 - p_0) = \sqrt{q_0}$$

#### The Asimov data set

To estimate median value of  $-2\ln\lambda(\mu)$ , consider special data set where all statistical fluctuations suppressed and  $n_i$ ,  $m_i$  are replaced by their expectation values (the "Asimov" data set):

$$n_{i} = \mu' s_{i} + b_{i}$$

$$m_{i} = u_{i}$$

$$\rightarrow \qquad \hat{\mu} = \mu' \qquad \hat{\theta} = \theta$$

$$\lambda_{A}(\mu) = \frac{L_{A}(\mu, \hat{\theta})}{L_{A}(\hat{\mu}, \hat{\theta})} = \frac{L_{A}(\mu, \hat{\theta})}{L_{A}(\mu', \theta)} \qquad Asimov value of$$

$$-2 \ln \lambda_{A}(\mu) = \frac{(\mu - \mu')^{2}}{\sigma^{2}} = \Lambda \qquad Asimov value of$$

$$-2 \ln \lambda_{A}(\mu) = \frac{(\mu - \mu')^{2}}{\sigma^{2}} = \Lambda \qquad or equivalently, \sigma$$

### Relation between test statistics and $\hat{\mu}$

Assuming Wald approximation, the relation between  $q_0$  and  $\hat{\mu}$  is



Monotonic, therefore quantiles of  $\hat{\mu}$  map one-to-one onto those of  $q_0$ , e.g.,

$$med[q_0] = q_0(med[\hat{\mu}]) = q_0(\mu') = \frac{{\mu'}^2}{{\sigma}^2} = -2\ln\lambda_A(0)$$
$$med[Z_0] = \sqrt{-2\ln\lambda_A(0)}$$

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# Higgs search with profile likelihood

Combination of Higgs boson search channels (ATLAS) Expected Performance of the ATLAS Experiment: Detector, Trigger and Physics, arXiv:0901.0512, CERN-OPEN-2008-20.

Standard Model Higgs channels considered (more to be used later):

$$\begin{split} H &\to \gamma \gamma \\ H &\to WW^{(*)} \to e\nu\mu\nu \\ H &\to ZZ^{(*)} \to 4l \ (l = e, \mu) \\ H &\to \tau^{+}\tau^{-} \to ll, lh \end{split}$$

Used profile likelihood method for systematic uncertainties: background rates, signal & background shapes.

# An example: ATLAS Higgs search (ATLAS Collab., CERN-OPEN-2008-020)

Statistical Combination of Several Important Standard Model Higgs Boson Search Channels.



Figure 12: The distribution of the test statistic  $q_0$  for  $H \rightarrow \tau^+ \tau^-$  under the null background-only hypothesis, for  $m_H = 130 \text{ GeV}$  with an integrated luminosity of 2 (a) and 10 (b) fb<sup>-1</sup>. A  $\frac{1}{2}\chi_1^2$  distribution is superimposed. Figures (c) and (d) show  $1 - F(q_0)$  where  $F(q_0)$  is the corresponding cumulative distribution. The small excess of events at high  $q_0$  is statistically compatible with the expected curves, as can be seen by comparison with the dotted histograms that show the 68.3% central confidence intervals for  $p = 1 - F(q_0|0)$ . The lower dotted line at  $2.87 \times 10^{-7}$  shows the 5 $\sigma$  discovery threshold.

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### Cumulative distributions of $q_0$

To validate to  $5\sigma$  level, need distribution out to  $q_0 = 25$ , i.e., around 10<sup>8</sup> simulated experiments.

Will do this if we really see something like a discovery.



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# Combination of channels

For a set of independent decay channels, full likelihood function is product of the individual ones:

$$L(\mu, \theta) = \prod_i L_i(\mu, \theta_i)$$

For combination need to form the full function and maximize to find estimators of  $\mu$ ,  $\theta$ .

→ ongoing ATLAS/CMS effort with RooStats framework

Trick for median significance: estimator for  $\mu$  is equal to the Asimov value  $\mu'$  for all channels separately, so for combination,

$$\lambda_{A}(\mu) = \prod_{i} \lambda_{A,i}(\mu)$$
 where  $\lambda_{A,i}(\mu) = \frac{L_{i}(\mu, \hat{\theta})}{L_{i}(\mu', \theta)}$ 

# Combined median significance



N.B. illustrates statistical method, but study did not include all usable Higgs channels.

# Discovery significance for $n \sim \text{Poisson}(s + b)$

Consider again the case where we observe n events, model as following Poisson distribution with mean s + b(assume b is known).

- 1) For an observed *n*, what is the significance  $Z_0$  with which we would reject the s = 0 hypothesis?
- 2) What is the expected (or more precisely, median )  $Z_0$  if the true value of the signal rate is *s*?

Gaussian approximation for Poisson significance For large s + b,  $n \rightarrow x \sim \text{Gaussian}(\mu, \sigma)$ ,  $\mu = s + b$ ,  $\sigma = \sqrt{(s + b)}$ . For observed value  $x_{\text{obs}}$ , *p*-value of s = 0 is  $\text{Prob}(x > x_{\text{obs}} | s = 0)$ ,:

$$p_0 = 1 - \Phi\left(\frac{x_{\rm obs} - b}{\sqrt{b}}\right)$$

Significance for rejecting s = 0 is therefore

$$Z_0 = \Phi^{-1}(1 - p_0) = \frac{x_{\text{obs}} - b}{\sqrt{b}}$$

Expected (median) significance assuming signal rate s is

$$\mathrm{median}[Z_0|s+b] = \frac{s}{\sqrt{b}}$$

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Better approximation for Poisson significance

Likelihood function for parameter *s* is

$$L(s) = \frac{(s+b)^n}{n!}e^{-(s+b)}$$

or equivalently the log-likelihood is

$$\ln L(s) = n \ln(s+b) - (s+b) - \ln n!$$

Find the maximum by setting  $\frac{\partial \ln L}{\partial s} = 0$ 

gives the estimator for s: 
$$\hat{s} = n - b$$

Approximate Poisson significance (continued) The likelihood ratio statistic for testing s = 0 is

$$q_0 = -2\ln\frac{L(0)}{L(\hat{s})} = 2\left(n\ln\frac{n}{b} + b - n\right) \quad \text{for } n > b, \ 0 \text{ otherwise}$$

For sufficiently large s + b, (use Wilks' theorem),

$$Z_0 \approx \sqrt{q_0} = \sqrt{2\left(n\ln\frac{n}{b} + b - n\right)}$$
 for  $n > b$ , 0 otherwise

To find median[ $Z_0|s+b$ ], let  $n \rightarrow s+b$ ,

$$\mathrm{median}[Z_0|s+b] \approx \sqrt{2\left((s+b)\ln(1+s/b)-s\right)}$$

This reduces to  $s/\sqrt{b}$  for s << b.

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# Wrapping up lecture 3

To establish discovery, we need (at least) to reject the background-only hypothesis (usually with Z > 5).

For two simple hypotheses, best test of  $H_0$  is from likelihood ratio  $L(H_1)/L(H_0)$ .

Alternative likelihood ratio  $\lambda(\mu) = L(\mu) / L(\hat{\mu})$  will lead to similar test; can use Wilks' theorem to obtain sampling distribution.

Methods for incorporating systematics somewhat different for the two approaches (in practice lead to similar results).

#### Extra slides

### Monte Carlo test of asymptotic formula

 $n \sim \text{Poisson}(\mu s + b)$  $m \sim \text{Poisson}(\tau b)$ 

Here take  $\tau = 1$ .

Asymptotic formula is good approximation to  $5\sigma$ level ( $q_0 = 25$ ) already for  $b \sim 20$ .



Monte Carlo test of asymptotic formulae

Significance from asymptotic formula, here  $Z_0 = \sqrt{q_0} = 4$ , compared to MC (true) value.

For very low b, asymptotic formula underestimates  $Z_0$ . Then slight overshoot before rapidly converging to MC value.



Monte Carlo test of asymptotic formulae Asymptotic  $f(q_0|1)$  good already for fairly small samples. Median[ $q_0|1$ ] from Asimov data set; good agreement with MC.



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