Bayesian statistical methods for HEP



Terascale Statistics School DESY, Hamburg

1 October, 2008



Glen Cowan Physics Department Royal Holloway, University of London g.cowan@rhul.ac.uk www.pp.rhul.ac.uk/~cowan

G. Cowan RHUL Physics

Outline

Tuesday:

The Bayesian method Bayesian assessment of uncertainties Bayesian computation: MCMC

Wednesday:

Bayesian limits Bayesian model selection ("discovery") Outlook for Bayesian methods in HEP

G. Cowan RHUL Physics

The Bayesian approach to limits

In Bayesian statistics need to start with 'prior pdf' $\pi(\theta)$, this reflects degree of belief about θ before doing the experiment.

Bayes' theorem tells how our beliefs should be updated in light of the data *x*:

$$p(\theta|x) = \frac{L(x|\theta)\pi(\theta)}{\int L(x|\theta')\pi(\theta') d\theta'} \propto L(x|\theta)\pi(\theta)$$

Integrate posterior pdf $p(\theta | x)$ to give interval with any desired probability content.

For e.g. Poisson parameter 95% CL upper limit from

$$0.95 = \int_{-\infty}^{s_{\rm up}} p(s|n) \, ds$$

G. Cowan RHUL Physics

Analytic formulae for limits

There are a number of papers describing Bayesian limits for a variety of standard scenarios

> Several conventional priors Systematics on efficiency, background Combination of channels

and (semi-)analytic formulae and software are provided.

Joel Heinrich, Bayesian limit software: multi-channel with correlated backgrounds and efficiencies, CDF/MEMO/STATISTICS/PUBLIC/7587 (2005).

Joel Heinrich et al., Interval estimation in the presence of nuisance parameters. 1. Bayesian approach, CDF/MEMO/STATISTICS/PUBLIC/7117, physics/0409129 (2004).

Luc Demortier, A Fully Bayesian Computation of Upper Limits for Poisson Processes, CDF/MEMO/STATISTICS/PUBLIC/5928 (2004).

But for more general cases we need to use numerical methods (e.g. L.D. uses importance sampling).

G. Cowan RHUL Physics Example: Poisson data with background

Count *n* events, e.g., in fixed time or integrated luminosity.

s = expected number of signal events

b = expected number of background events

$$n \sim \text{Poisson}(s+b)$$
: $P(n;s,b) = \frac{(s+b)^n}{n!}e^{-(s+b)}$

Sometimes *b* known, other times it is in some way uncertain.

Goal: measure or place limits on *s*, taking into consideration the uncertainty in *b*.

Widely discussed in HEP community, see e.g. proceedings of PHYSTAT meetings, Durham, Fermilab, CERN workshops...

G. Cowan RHUL Physics

Bayesian prior for Poisson parameter

Include knowledge that $s \ge 0$ by setting prior $\pi(s) = 0$ for s < 0. Often try to reflect 'prior ignorance' with e.g.

$$\pi(s) = egin{cases} 1 & s \geq 0 \ 0 & ext{otherwise} \end{cases}$$

Not normalized but this is OK as long as L(s) dies off for large s.

Not invariant under change of parameter — if we had used instead a flat prior for, say, the mass of the Higgs boson, this would imply a non-flat prior for the expected number of Higgs events.

Doesn't really reflect a reasonable degree of belief, but often used as a point of reference;

or viewed as a recipe for producing an interval whose frequentist properties can be studied (coverage will depend on true *s*). G. Cowan RHUL Physics Bayesian methods for HEP / DESY Terascale School page 6 Bayesian interval with flat prior for s

Solve numerically to find limit s_{up} .

For special case b = 0, Bayesian upper limit with flat prior numerically same as classical case ('coincidence').

Otherwise Bayesian limit is everywhere greater than classical ('conservative').

Never goes negative.

Doesn't depend on *b* if n = 0.



G. Cowan RHUL Physics

Upper limit versus b



If *n* = 0 observed, should upper limit depend on *b*? Classical: yes Bayesian: no FC: yes

G. Cowan RHUL Physics

Coverage probability of confidence intervals

Because of discreteness of Poisson data, probability for interval to include true value in general > confidence level ('over-coverage')



G. Cowan RHUL Physics

Bayesian limits with uncertainty on b

Uncertainty on b goes into the prior, e.g.,

 $\begin{aligned} \pi(s,b) &= \pi_s(s)\pi_b(b) \quad \text{(or include correlations as appropriate)} \\ \pi_s(s) &= \text{const}, \quad \sim 1/s, \dots \\ \pi_b(b) &= \frac{1}{\sqrt{2\pi}\sigma_b} e^{-(b-b_{\text{meas}})^2/2\sigma_b^2} \quad \text{(or whatever)} \end{aligned}$

Put this into Bayes' theorem,

 $p(s,b|n) \propto L(n|s,b)\pi(s,b)$

Marginalize over *b*, then use p(s|n) to find intervals for *s* with any desired probability content.

Controversial part here is prior for signal $\pi_s(s)$ (treatment of nuisance parameters is easy).

G. Cowan RHUL Physics

Discussion on limits

Different sorts of limits answer different questions.

A frequentist confidence interval does not (necessarily) answer, "What do we believe the parameter's value is?"

Coverage — nice, but crucial?

Look at sensitivity, e.g., $E[s_{up} | s = 0]$.

Consider also:

politics, need for consensus/conventions; convenience and ability to combine results, ...

For any result, consumer will compute (mentally or otherwise):

 $p(\theta|\text{result}) \propto L(\text{result}|\theta)\pi(\theta)$

Need likelihood (or summary thereof).

consumer's prior

G. Cowan RHUL Physics

Frequentist discovery, *p*-values

To discover e.g. the Higgs, try to reject the background-only (null) hypothesis (H_0) .

Define a statistic *t* whose value reflects compatibility of data with H_0 .

p-value = Prob(data with \leq compatibility with H_0 when compared to the data we got $| H_0 \rangle$) For example, if high values of t mean less compatibility,

$$p=\int_t^\infty f(t'|H_0)\,dt'$$
 .

If *p*-value comes out small, then this is evidence against the background-only hypothesis \rightarrow discovery made!

G. Cowan RHUL Physics

Significance from *p*-value

Define significance Z as the number of standard deviations that a Gaussian variable would fluctuate in one direction to give the same p-value.



$$p = \int_Z^\infty rac{1}{\sqrt{2\pi}} e^{-x^2/2}\,dx = 1 - \Phi(Z)$$
 TMath::Prob

 $Z = \Phi^{-1}(1-p)$ TMath::NormQuantile

G. Cowan RHUL Physics

When to publish

HEP folklore is to claim discovery when $p = 2.85 \times 10^{-7}$, corresponding to a significance Z = 5.

This is very subjective and really should depend on the prior probability of the phenomenon in question, e.g.,

phenomenonreasonable p-value for discovery D^0D^0 mixing~0.05Higgs~10^{-7} (?)Life on Mars~10^{-10}Astrology~10^{-20}

Note some groups have defined 5σ to refer to a two-sided fluctuation, i.e., $p = 5.7 \times 10^{-7}$

G. Cowan RHUL Physics

Bayesian model selection ('discovery')

The probability of hypothesis H_0 relative to its complementary alternative H_1 is often given by the posterior odds:



The Bayes factor is regarded as measuring the weight of evidence of the data in support of H_0 over H_1 .

Interchangeably use $B_{10} = 1/B_{01}$

G. Cowan RHUL Physics

Assessing Bayes factors

One can use the Bayes factor much like a *p*-value (or *Z* value). There is an "established" scale, analogous to our 5σ rule:

 B_{10} Evidence against H_0

1 to 3Not worth more than a bare mention3 to 20Positive

- 5 to 20 Positive
- 20 to 150 Strong

RHUL Physics

> 150 Very strong

Kass and Raftery, Bayes Factors, J. Am Stat. Assoc 90 (1995) 773.

11 May 07: Not clear how useful this scale is for HEP.
3 Sept 07: Upon reflection & PHYSTAT07 discussion, seems like an intuitively useful complement to *p*-value.

Rewriting the Bayes factor

Suppose we have models H_i , i = 0, 1, ...,

each with a likelihood $p(x|H_i, \vec{\theta_i})$

and a prior pdf for its internal parameters $\pi_i(\vec{\theta_i})$

so that the full prior is $\pi(H_i, \vec{\theta}_i) = p_i \pi_i(\vec{\theta}_i)$

where $p_i = P(H_i)$ is the overall prior probability for H_i .

The Bayes factor comparing H_i and H_j can be written

$$B_{ij} = \frac{P(H_i|\vec{x})}{P(H_i)} / \frac{P(H_j|\vec{x})}{P(H_j)}$$

G. Cowan RHUL Physics

Bayes factors independent of $P(H_i)$

For B_{ij} we need the posterior probabilities marginalized over all of the internal parameters of the models:

$$P(H_i|\vec{x}) = \int P(H_i, \vec{\theta}_i | \vec{x}) d\vec{\theta}_i \qquad \text{Use Bayes} \\ = \frac{\int L(\vec{x} | H_i, \vec{\theta}_i) p_i \pi_i(\vec{\theta}_i) d\vec{\theta}_i}{P(x)}$$

So therefore the Bayes factor is

Ratio of marginal likelihoods

$$B_{ij} = \frac{\int L(\vec{x}|H_i, \vec{\theta}_i) \pi_i(\vec{\theta}_i) d\vec{\theta}_i}{\int L(\vec{x}|H_j, \vec{\theta}_j) \pi_j(\vec{\theta}_j) d\vec{\theta}_j}$$

The prior probabilities $p_i = P(H_i)$ cancel.

G. Cowan RHUL Physics

Numerical determination of Bayes factors

Both numerator and denominator of B_{ii} are of the form

Various ways to compute these, e.g., using sampling of the posterior pdf (which we can do with MCMC).

Harmonic Mean (and improvements) Importance sampling Parallel tempering (~thermodynamic integration)

See e.g. Kass and Raftery, Bayes Factors, J. Am. Stat. Assoc. 90 (1995) 773-795.

G. Cowan RHUL Physics . . .

Harmonic mean estimator

E.g., consider only one model and write Bayes theorem as:

$$\frac{\pi(\boldsymbol{\theta})}{m} = \frac{p(\boldsymbol{\theta}|\mathbf{x})}{L(\mathbf{x}|\boldsymbol{\theta})}$$

 $\pi(\theta)$ is normalized to unity so integrate both sides,

posterior expectation

$$m^{-1} = \int \frac{1}{L(\mathbf{x}|\boldsymbol{\theta})} p(\boldsymbol{\theta}|\mathbf{x}) d\boldsymbol{\theta} = E_p[1/L]$$

Therefore sample θ from the posterior via MCMC and estimate *m* with one over the average of 1/L (the harmonic mean of *L*).

M.A. Newton and A.E. Raftery, Approximate Bayesian Inference by the Weighted Likelihood Bootstrap, Journal of the Royal Statistical Society B 56 (1994) 3-48.

G. Cowan RHUL Physics

Improvements to harmonic mean estimator

The harmonic mean estimator is numerically very unstable; formally infinite variance (!). Gelfand & Dey propose variant:

Rearrange Bayes thm; multiply both sides by arbitrary $pdf f(\theta)$:

$$\frac{f(\boldsymbol{\theta})p(\boldsymbol{\theta}|\mathbf{x})}{L(\mathbf{x}|\boldsymbol{\theta})\pi(\boldsymbol{\theta})} = \frac{f(\boldsymbol{\theta})}{m}$$

Integrate over
$$\boldsymbol{\theta}$$
: $m^{-1} = \int \frac{f(\boldsymbol{\theta})}{L(\mathbf{x}|\boldsymbol{\theta})\pi(\boldsymbol{\theta})} p(\boldsymbol{\theta}|\mathbf{x}) = E_p \left[\frac{f(\boldsymbol{\theta})}{L(\mathbf{x}|\boldsymbol{\theta})\pi(\boldsymbol{\theta})} \right]$

Improved convergence if tails of $f(\theta)$ fall off faster than $L(x|\theta)\pi(\theta)$ Note harmonic mean estimator is special case $f(\theta) = \pi(\theta)$.

A.E. Gelfand and D.K. Dey, *Bayesian model choice: asymptotics and exact calculations*, Journal of the Royal Statistical Society B 56 (1994) 501-514.

G. Cowan RHUL Physics

Importance sampling

Need pdf $f(\theta)$ which we can evaluate at arbitrary θ and also sample with MC.

The marginal likelihood can be written

$$m = \int \frac{L(\mathbf{x}|\boldsymbol{\theta})\pi(\boldsymbol{\theta})}{f(\boldsymbol{\theta})} f(\boldsymbol{\theta}) d\boldsymbol{\theta} = E_f \left[\frac{L(\mathbf{x}|\boldsymbol{\theta})\pi(\boldsymbol{\theta})}{f(\boldsymbol{\theta})} \right]$$

Best convergence when $f(\theta)$ approximates shape of $L(x|\theta)\pi(\theta)$.

Use for $f(\theta)$ e.g. multivariate Gaussian with mean and covariance estimated from posterior (e.g. with MINUIT).

G. Cowan RHUL Physics

Bayes factor computation discussion

Also can use method of parallel tempering; see e.g.

Phil Gregory, Bayesian Logical Data Analysis for the Physical Sciences, Cambridge University Press, 2005.

Harmonic mean OK for very rough estimate.

I had trouble with all of the methods based on posterior sampling.

Importance sampling worked best, but may not scale well to higher dimensions.

Lots of discussion of this problem in the literature, e.g.,

Cong Han and Bradley Carlin, Markov Chain Monte Carlo Methods for Computing Bayes Factors: A Comparative Review, J. Am. Stat. Assoc. 96 (2001) 1122-1132.

G. Cowan RHUL Physics

Bayesian Higgs analysis

N independent channels, count n_i events in search regions:

Constrain expected background
$$b_i$$
 with sideband measurements:

$$P(\mathbf{n}|\mathbf{s}, \mathbf{b}) = \prod_{i=1}^{N} \frac{(s_i + b_i)^{n_i}}{n_i!} e^{-(s_i + b_i)}$$

$$P(\mathbf{m}|\mathbf{b}, oldsymbol{ au}) = \prod_{i=1}^{N} rac{(au_i b_i)^{m_i}}{m_i!} \, e^{- au_i b_i}$$

page 24

Expected number of signal events: $s_i = \mu \sigma_{SM} \mathcal{B}_i \varepsilon_{s,i} L_i \equiv \mu \varphi_i$ (μ is global parameter, $\mu = 1$ for SM).

Consider a fixed Higgs mass and assume SM branching ratios B_i .

Suggested method: constrain μ with limit μ_{up} ; consider m_{H} excluded if upper limit $\mu_{up} < 1.0$.

For discovery, compute Bayes factor for $H_0: \mu = 0$ vs. $H_1: \mu = 1$ G. Cowan
RHUL PhysicsBayesian methods for HEP / DESY Terascale School

Parameters of Higgs analysis

E.g. combine cross section, branching ratio, luminosity, efficiency into a single factor ϕ :

$$s_i = \mu \sigma_{\rm SM} \mathcal{B}_i \varepsilon_{s,i} L_i \equiv \mu \varphi_i$$

Systematics in any of the factors can be described by a prior for ϕ , use e.g. Gamma distribution. For now ignore correlations, but these would be present e.g. for luminosity error:

$$\pi_{oldsymbol{arphi}}(oldsymbol{arphi}) = \prod_{i=1}^{N} rac{a_i (a_i arphi_i)^{b_i - 1} e^{-a_i arphi_i}}{\Gamma(b_i)}$$

 a_{i} , b_{i} from nominal value $\phi_{i,0}$ and relative error $r_{i} = \sigma_{\phi,i} / \phi_{i,0}$:

$$a \;\; = \;\; rac{1}{arphi_0 r_arphi^2} \,, \quad b \;\; = \;\; rac{1}{r_arphi^2} \,.$$

G. Cowan RHUL Physics

Bayes factors for Higgs analysis

The Bayes factor B_{10} is

$$B_{10} = \frac{\int \int \int L(\mathbf{n}, \mathbf{m} | \mu, \mathbf{b}, \varphi) \pi_{\mu}(\mu) \pi_{\varphi}(\varphi) \pi_{\mathbf{b}}(\mathbf{b}) d\mu d\varphi d\mathbf{b}}{\int \int L(\mathbf{n}, \mathbf{m} | \mu = 0, \mathbf{b}, \varphi) \pi_{\varphi}(\varphi) \pi_{\mathbf{b}}(\mathbf{b}) d\varphi d\mathbf{b}}$$

Compute this using a fixed μ for H_1 , i.e., $\pi_{\mu}(\mu) = \delta(\mu - \mu')$, then do this as a function of μ' . Look in particular at $\mu = 1$. Take numbers from VBF paper for 10 fb⁻¹, $m_{\rm H} = 130$ GeV:

| Channel | s | <i>b</i> | |
|--------------------------------|------|----------|---|
| $H \to WW^* \to e \mu + X$ | 12.3 | 9.2 | $\sim lvjj$ was for 30 fb ⁻¹ , |
| $H \to WW^* \to ee/\mu\mu + X$ | 11.7 | 10.1 | in paper: divided by 3 |
| $H \to WW^* \to l\nu jj + X$ | 1.5 | 2.0 | in puper, arriaea ey e |

S. Asai et al., Prospects for the Search for a Standard Model Higgs Boson in ATLAS using Vector Boson Fusion, Eur. Phys. J. C32S2 (2004) 19-54; hep-ph/0402254.

G. Cowan RHUL Physics

Bayes factors for Higgs analysis: results (1)

Create data set by hand with $n_i \sim$ nearest integer $(\phi_i + b_i)$, i.e., $\mu = 1$: $n_1 = 22, n_2 = 22, n_3 = 4.$

For the sideband measurements m_i , choose desired σ_b/b , use this to set size of sideband (i.e. $\sigma_b/b = 0.1 \rightarrow m = 100$).



 B_{10} for $\sigma_{\phi}/\phi = 0.1$, different values of σ_{b}/b ., as a function of μ .

G. Cowan RHUL Physics

Bayes factors for Higgs analysis: results (2)



Effect of uncertainty in ϕ_i (e.g., in the efficiency): $\mu = 1$ no longer gives a fixed s_i , but a smeared out distribution.

 \rightarrow lower peak value of B_{10} .

G. Cowan RHUL Physics

Bayes factors for Higgs analysis: results (3)

Or try data set with $n_i \sim$ nearest integer b_i , i.e., $\mu = 0$: $n_1 = 9, n_2 = 10, n_3 = 2$. Used $\sigma_b/b = 0.1, \sigma_{\phi}/\phi, = 0.1$.



Here the SM $\mu = 1$ is clearly disfavoured, so we set a limit on μ .

G. Cowan RHUL Physics

Posterior pdf for μ , upper limits (1)

Here done with (improper) uniform prior, $\mu > 0$. (Can/should also vary prior.)



Figure 4: The posterior distribution of μ with a data set compatible with background only $(n_1 = 9, n_2 = 10, n_3 = 2)$ for $\sigma_{\varphi}/\varphi = 0.1$ and several values of σ_b/b .

G. Cowan RHUL Physics

Posterior pdf for μ , upper limits (2)



Figure 5: The posterior distribution of μ with a data set compatible with background only $(n_1 = 9, n_2 = 10, n_3 = 2)$ for $\sigma_b/b = 0.1$ and several values of σ_{φ}/φ .

G. Cowan RHUL Physics

Outlook for Bayesian methods in HEP Bayesian methods allow (indeed require) prior information about the parameters being fitted.

This type of prior information can be difficult to incorporate into a frequentist analysis

This will be particularly relevant when estimating uncertainties on predictions of LHC observables that may stem from theoretical uncertainties, parton densities based on inconsistent data, etc.

Prior ignorance is not well defined. If that's what you've got, don't expect Bayesian methods to provide a unique solution. Try a reasonable variation of priors -- if that yields large variations in the posterior, you don't have much information coming in from the data.

You do not have to be exclusively a Bayesian or a Frequentist Use the right tool for the right job

G. Cowan RHUL Physics