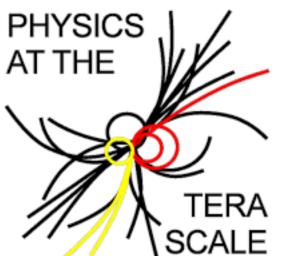
Statistics for Particle Physics Lecture 2



Helmholtz Alliance



Terascale Statistics School https://indico.desy.de/event/43398/

DESY, Hamburg 2-5 April 2024



Glen Cowan Physics Department Royal Holloway, University of London g.cowan@rhul.ac.uk www.pp.rhul.ac.uk/~cowan

Outline

Tuesday 11:05IntroductionProbabilityHypothesis tests, parameter estimation

→ Wednesday 9:15 From tests to ML (finish from Tuesday)
 Confidence limits
 Systematic uncertainties
 General analysis, asymptotics

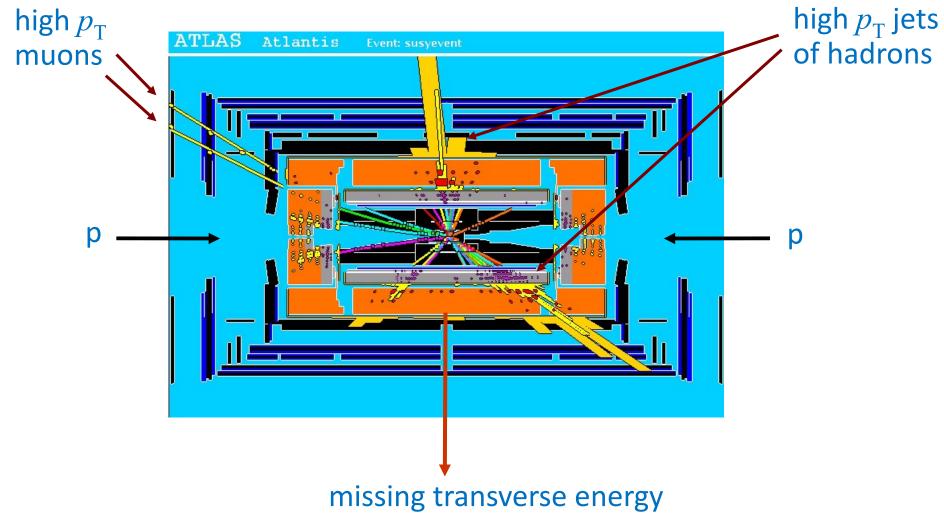
Thursday 16:00 "Errors on errors"

More resources in the University of London course: http://www.pp.rhul.ac.uk/~cowan/stat_course.html

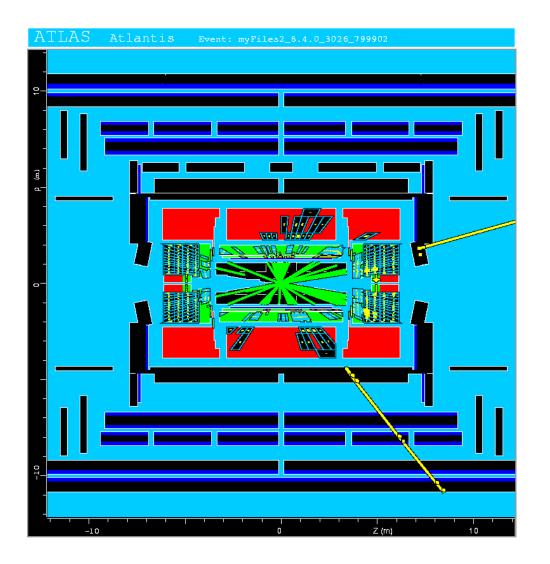
G. Cowan / RHUL Physics

Particle Physics context for a hypothesis test

A simulated SUSY event ("signal"):



Background events



This event from Standard Model ttbar production also has high $p_{\rm T}$ jets and muons, and some missing transverse energy.

→ can easily mimic a signal event.

Classification of proton-proton collisions

Proton-proton collisions can be considered to come in two classes: signal (the kind of event we're looking for, y = 1) background (the kind that mimics signal, y = 0)

For each collision (event), we measure a collection of features:

 $x_1 =$ energy of muon $x_4 =$ missing transverse energy $x_2 =$ angle between jets $x_5 =$ invariant mass of muon pair $x_3 =$ total jet energy $x_6 = \dots$

The real events don't come with true class labels, but computersimulated events do. So we can have a set of simulated events that consist of a feature vector x and true class label y (0 for background, 1 for signal):

$$(x, y)_1, (x, y)_2, ..., (x, y)_N$$

The simulated events are called "training data".

Distributions of the features

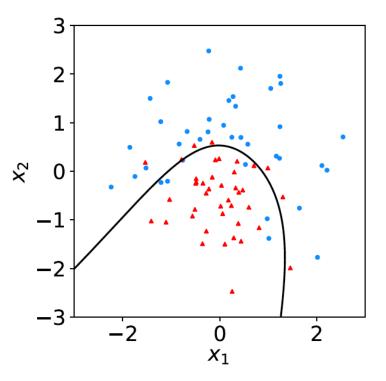
If we consider only two features $x = (x_1, x_2)$, we can display the results in a scatter plot (red: y = 0, blue: y = 1).

For real events, the dots are black (true type is not known).

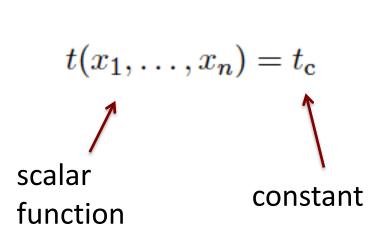
For each real event test the hypothesis that it is background.

(Related to this: test that a sample of events is *all* background.)

The test's critical region is defined by a "decision boundary" – without knowing the event type, we can classify them by seeing where their measured features lie relative to the boundary.



Decision function, test statistic

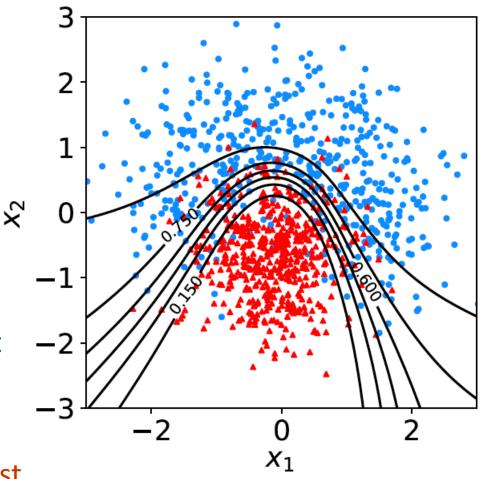


A surface in an *n*-dimensional

space can be described by

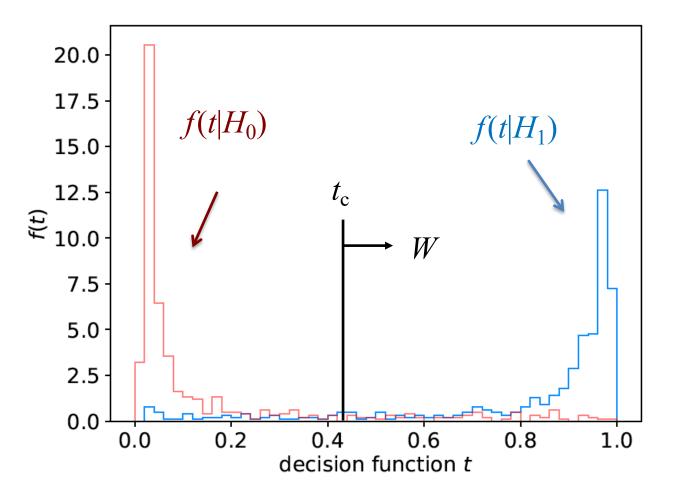
Different values of the constant t_c result in a family of surfaces.

Problem is reduced to finding the best decision function or test statistic t(x).



Distribution of t(x)

By forming a test statistic t(x), the boundary of the critical region in the *n*-dimensional *x*-space is determined by a single single value t_c .

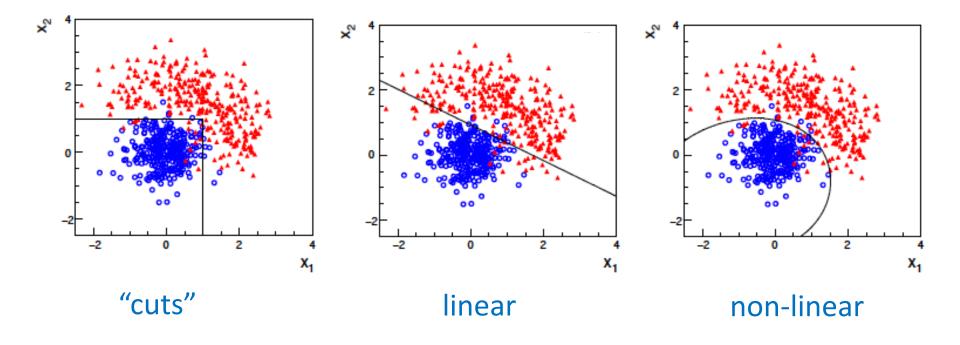


Types of decision boundaries

So what is the optimal boundary for the critical region, i.e., what is the optimal test statistic t(x)?

First find best t(x), later address issue of optimal size of test.

Remember *x*-space can have many dimensions.



Test statistic based on likelihood ratio

How can we choose a test's critical region in an 'optimal way', in particular if the data space is multidimensional?

Neyman-Pearson lemma states:

For a test of H_0 of size α , to get the highest power with respect to the alternative H_1 we need for all x in the critical region W

"likelihood ratio (LR)"
$$\frac{P(\mathbf{x}|H_1)}{P(\mathbf{x}|H_0)} \ge c_{\alpha}$$

inside W and $\leq c_{\alpha}$ outside, where c_{α} is a constant chosen to give a test of the desired size.

Equivalently, optimal scalar test statistic is

$$t(\mathbf{x}) = \frac{P(\mathbf{x}|H_1)}{P(\mathbf{x}|H_0)}$$

N.B. any monotonic function of this is leads to the same test.

G. Cowan / RHUL Physics

Neyman-Pearson doesn't usually help

We usually don't have explicit formulae for the pdfs f(x|s), f(x|b), so for a given x we can't evaluate the likelihood ratio

$$t(\mathbf{x}) = \frac{f(\mathbf{x}|s)}{f(\mathbf{x}|b)}$$

Instead we may have Monte Carlo models for signal and background processes, so we can produce simulated data:

generate
$$\boldsymbol{x} \sim f(\boldsymbol{x}|\mathbf{s}) \rightarrow \boldsymbol{x}_1, \dots, \boldsymbol{x}_N$$

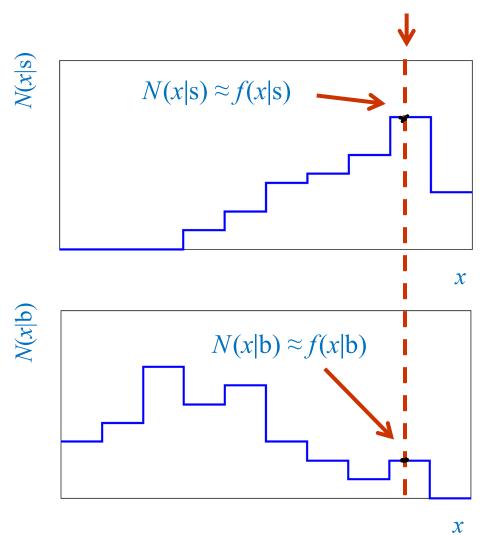
generate
$$\mathbf{x} \sim f(\mathbf{x}|\mathbf{b}) \longrightarrow \mathbf{x}_1, ..., \mathbf{x}_N$$

This gives samples of "training data" with events of known type.

 Use these to construct a statistic that is as close as possible to the optimal likelihood ratio (→ Machine Learning).

Approximate LR from histograms

Want t(x) = f(x|s)/f(x|b) for x here



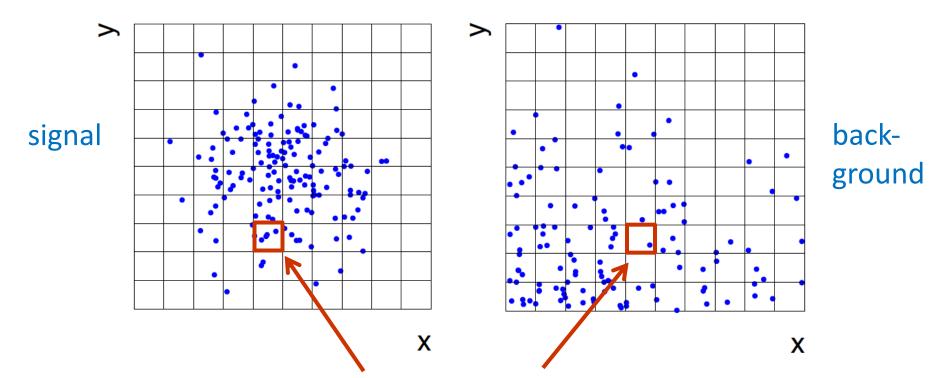
One possibility is to generate MC data and construct histograms for both signal and background.

Use (normalized) histogram values to approximate LR:

$$t(x) \approx \frac{N(x|s)}{N(x|b)}$$

Can work well for single variable.

Approximate LR from 2D-histograms Suppose problem has 2 variables. Try using 2-D histograms:



Approximate pdfs using N(x,y|s), N(x,y|b) in corresponding cells. But if we want M bins for each variable, then in n-dimensions we have M^n cells; can't generate enough training data to populate.

 \rightarrow Histogram method usually not usable for n > 1 dimension.

Strategies for multivariate analysis

Neyman-Pearson lemma gives optimal answer, but cannot be used directly, because we usually don't have $f(\mathbf{x}|\mathbf{s}), f(\mathbf{x}|\mathbf{b})$.

Histogram method with M bins for n variables requires that we estimate M^n parameters (the values of the pdfs in each cell), so this is rarely practical.

A compromise solution is to assume a certain functional form for the test statistic t(x) with fewer parameters; determine them (using MC) to give best separation between signal and background.

Alternatively, try to estimate the probability densities $f(\mathbf{x}|\mathbf{s})$ and $f(\mathbf{x}|\mathbf{b})$ (with something better than histograms) and use the estimated pdfs to construct an approximate likelihood ratio.

Multivariate methods (Machine Learning)

Many new (and some old) methods:

Fisher discriminant

(Deep) Neural Networks

Kernel density methods

Support Vector Machines

Decision trees

Boosting

Bagging

More in the lectures by Tilman Plehn

Confidence intervals by inverting a test

In addition to a 'point estimate' of a parameter we should report an interval reflecting its statistical uncertainty.

Confidence intervals for a parameter θ can be found by defining a test of the hypothesized value θ (do this for all θ):

Specify values of the data that are 'disfavoured' by θ (critical region) such that $P(\text{data in critical region} | \theta) \le \alpha$ for a prespecified α , e.g., 0.05 or 0.1.

If data observed in the critical region, reject the value θ .

Now invert the test to define a confidence interval as:

set of θ values that are not rejected in a test of size α (confidence level CL is $1 - \alpha$).

Relation between confidence interval and *p*-value

Equivalently we can consider a significance test for each hypothesized value of θ , resulting in a *p*-value, p_{θ} .

If $p_{\theta} \leq \alpha$, then we reject θ .

The confidence interval at $CL = 1 - \alpha$ consists of those values of θ that are not rejected.

E.g. an upper limit on θ is the greatest value for which $p_{\theta} > \alpha$.

In practice find by setting $p_{\theta} = \alpha$ and solve for θ .

For a multidimensional parameter space $\theta = (\theta_1, \dots, \theta_M)$ use same idea – result is a confidence "region" with boundary determined by $p_{\theta} = \alpha$.

Coverage probability of confidence interval

If the true value of θ is rejected, then it's not in the confidence interval. The probability for this is by construction (equality for continuous data):

 $P(\text{reject } \theta | \theta) \leq \alpha = \text{type-I error rate}$

Therefore, the probability for the interval to contain or "cover" θ is

P(conf. interval "covers" $\theta | \theta \ge 1 - \alpha$

This assumes that the set of θ values considered includes the true value, i.e., it assumes the composite hypothesis $P(\mathbf{x}|H,\theta)$.

Frequentist upper limit on Poisson parameter

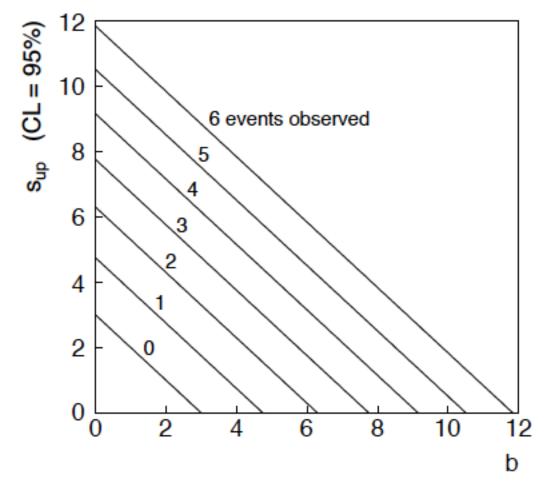
Consider again the case of observing $n \sim \text{Poisson}(s + b)$. Suppose b = 4.5, $n_{\text{obs}} = 5$. Find upper limit on s at 95% CL. Relevant alternative is s = 0 (critical region at low n) p-value of hypothesized s is $P(n \le n_{\text{obs}}; s, b)$ Upper limit s_{up} at $\text{CL} = 1 - \alpha$ found from

$$\begin{aligned} \alpha &= P(n \le n_{\text{obs}}; s_{\text{up}}, b) = \sum_{n=0}^{n_{\text{obs}}} \frac{(s_{\text{up}} + b)^n}{n!} e^{-(s_{\text{up}} + b)} \\ s_{\text{up}} &= \frac{1}{2} F_{\chi^2}^{-1} (1 - \alpha; 2(n_{\text{obs}} + 1)) - b \\ &= \frac{1}{2} F_{\chi^2}^{-1} (0.95; 2(5 + 1)) - 4.5 = 6.0 \end{aligned}$$

G. Cowan / RHUL Physics

$n \sim \text{Poisson}(s+b)$: frequentist upper limit on s

For low fluctuation of *n*, formula can give negative result for s_{up} ; i.e. confidence interval is empty; all values of $s \ge 0$ have $p_s \le \alpha$.



Limits near a boundary of the parameter space

Suppose e.g. b = 2.5 and we observe n = 0.

If we choose CL = 0.9, we find from the formula for s_{up}

$$s_{\rm up} = -0.197$$
 (CL = 0.90)

Physicist:

We already knew $s \ge 0$ before we started; can't use negative upper limit to report result of expensive experiment!

Statistician:

The interval is designed to cover the true value only 90% of the time — this was clearly not one of those times.

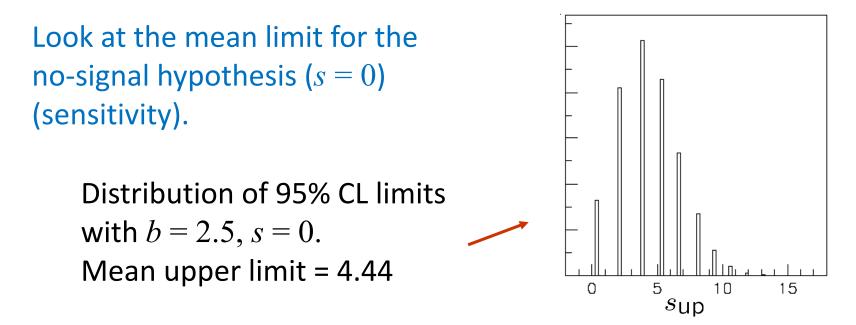
Not uncommon dilemma when testing parameter values for which one has very little experimental sensitivity, e.g., very small *s*.

Expected limit for s = 0

Physicist: I should have used CL = 0.95 — then $s_{up} = 0.496$

Even better: for CL = 0.917923 we get $s_{up} = 10^{-4}$!

Reality check: with b = 2.5, typical Poisson fluctuation in n is at least $\sqrt{2.5} = 1.6$. How can the limit be so low?



Approximate confidence intervals/regions from the likelihood function

Suppose we test parameter value(s) $\theta = (\theta_1, ..., \theta_n)$ using the ratio

$$\lambda(\theta) = \frac{L(\theta)}{L(\hat{\theta})} \qquad \qquad 0 \le \lambda(\theta) \le 1$$

Lower $\lambda(\theta)$ means worse agreement between data and hypothesized θ . Equivalently, usually define

$$t_{\theta} = -2\ln\lambda(\theta)$$

so higher t_{θ} means worse agreement between θ and the data.

p-value of θ therefore

$$p_{\theta} = \int_{t_{\theta,\text{obs}}}^{\infty} f(t_{\theta}|\theta) \, dt_{\theta}$$

G. Cowan / RHUL Physics

Confidence region from Wilks' theorem

Wilks' theorem says (in large-sample limit and provided certain conditions hold...)

 $f(t_{\theta}|\theta) \sim \chi_n^2 \qquad \begin{array}{l} \text{chi-square dist. with $\#$ d.o.f. =} \\ \# \text{ of components in $\theta = (\theta_1, ..., \theta_n)$.} \end{array}$

Assuming this holds, the *p*-value is

$$p_{m{ heta}} = 1 - F_{\chi^2_n}(t_{m{ heta}}) \quad \leftarrow \text{set equal to } lpha$$

To find boundary of confidence region set $p_{\theta} = \alpha$ and solve for t_{θ} :

$$t_{\theta} = F_{\chi_n^2}^{-1}(1-\alpha)$$

Recall also

$$t_{\theta} = -2\ln\frac{L(\theta)}{L(\hat{\theta})}$$

G. Cowan / RHUL Physics

Confidence region from Wilks' theorem (cont.)

i.e., boundary of confidence region in θ space is where

$$\ln L(\theta) = \ln L(\hat{\theta}) - \frac{1}{2}F_{\chi_n^2}^{-1}(1-\alpha)$$

For example, for $1 - \alpha = 68.3\%$ and n = 1 parameter,

$$F_{\chi_1^2}^{-1}(0.683) = 1$$

and so the 68.3% confidence level interval is determined by

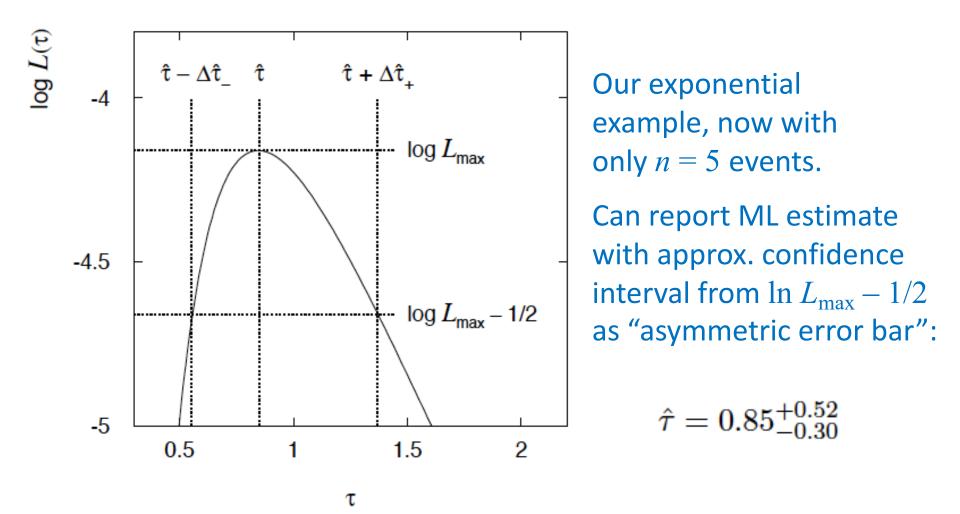
$$\ln L(\theta) = \ln L(\hat{\theta}) - \frac{1}{2}$$

Same as recipe for finding the estimator's standard deviation, i.e.,

 $[\hat{\theta} - \sigma_{\hat{\theta}}, \hat{\theta} + \sigma_{\hat{\theta}}]$ is a 68.3% CL confidence interval.

Example of interval from $\ln L(\theta)$

For n=1 parameter, CL = 0.683, $Q_{\alpha} = 1$.



Multiparameter case

For increasing number of parameters, $CL = 1 - \alpha$ decreases for confidence region determined by a given

$$Q_{\alpha} = F_{\chi_n^2}^{-1}(1-\alpha)$$

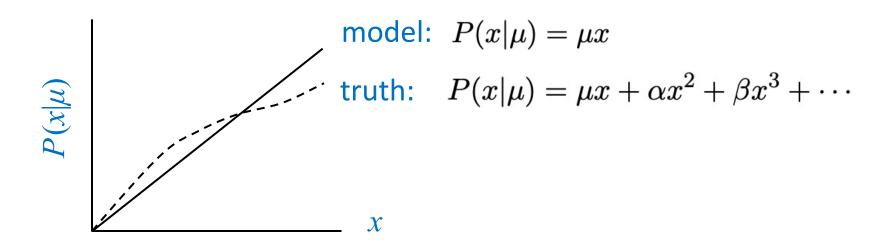
Q_{lpha}	1-lpha					
	n = 1	n = 2	n = 3	n = 4	n = 5	
1.0	0.683	0.393	0.199	0.090	0.037	
2.0	0.843	0.632	0.428	0.264	0.151	
4.0	0.954	0.865	0.739	0.594	0.451	
9.0	0.997	0.989	0.971	0.939	0.891	

Multiparameter case (cont.)

Equivalently, Q_{α} increases with *n* for a given $CL = 1 - \alpha$.

$1 - \alpha$	\widehat{Q}_{lpha}					
	n = 1	n = 2	n = 3	n = 4	n = 5	
0.683	1.00	2.30	3.53	4.72	5.89	
0.90	2.71	4.61	6.25	7.78	9.24	
0.95	3.84	5.99	7.82	9.49	11.1	
0.99	6.63	9.21	11.3	13.3	15.1	

Systematic uncertainties and nuisance parameters In general, our model of the data is not perfect:



Can improve model by including additional adjustable parameters.

$$P(x|\mu) \to P(x|\mu, \theta)$$

Nuisance parameter ↔ systematic uncertainty. Some point in the parameter space of the enlarged model should be "true".

Presence of nuisance parameter decreases sensitivity of analysis to the parameter of interest (e.g., increases variance of estimate).

Profile Likelihood

Suppose we have a likelihood $L(\mu, \theta) = P(x|\mu, \theta)$ with Nparameters of interest $\mu = (\mu_1, ..., \mu_N)$ and M nuisance parameters $\theta = (\theta_1, ..., \theta_M)$. The "profiled" (or "constrained") values of θ are:

$$\hat{\hat{\boldsymbol{\theta}}}(\boldsymbol{\mu}) = \operatorname*{argmax}_{\boldsymbol{\theta}} L(\boldsymbol{\mu}, \boldsymbol{\theta})$$

and the profile likelihood is: $L_{\rm p}(\boldsymbol{\mu}) = L(\boldsymbol{\mu}, \hat{\boldsymbol{\theta}})$

The profile likelihood depends only on the parameters of interest; the nuisance parameters are replaced by their profiled values.

The profile likelihood can be used to obtain confidence intervals/regions for the parameters of interest in the same way as one would for all of the parameters from the full likelihood.

Profile Likelihood Ratio – Wilks theorem

Goal is to test/reject regions of μ space (param. of interest).

Rejecting a point μ should mean $p_{\mu} \leq \alpha$ for all possible values of the nuisance parameters θ .

Test $\boldsymbol{\mu}$ using the "profile likelihood ratio": $\lambda(\boldsymbol{\mu}) = \frac{L(\boldsymbol{\mu}, \hat{\boldsymbol{\theta}})}{L(\hat{\boldsymbol{\mu}}, \hat{\boldsymbol{\theta}})}$

Let $t_{\mu} = -2 \ln \lambda(\mu)$. Wilks' theorem says in large-sample limit: $t_{\mu} \sim \text{chi-square}(N)$

where the number of degrees of freedom is the number of parameters of interest (components of μ). So *p*-value for μ is

$$p_{\boldsymbol{\mu}} = \int_{t_{\boldsymbol{\mu},\text{obs}}}^{\infty} f(t_{\boldsymbol{\mu}} | \boldsymbol{\mu}, \boldsymbol{\theta}) \, dt_{\boldsymbol{\mu}} = 1 - F_{\chi_N^2}(t_{\boldsymbol{\mu},\text{obs}})$$

G. Cowan / RHUL Physics

Profile Likelihood Ratio – Wilks theorem (2)

If we have a large enough data sample to justify use of the asymptotic chi-square pdf, then if μ is rejected, it is rejected for any values of the nuisance parameters.

The recipe to get confidence regions/intervals for the parameters of interest at $CL = 1 - \alpha$ is thus the same as before, simply use the profile likelihood:

$$\ln L_{\rm p}(\boldsymbol{\mu}) = \ln L_{\rm max} - \frac{1}{2} F_{\chi_N^2}^{-1} (1 - \alpha)$$

where the number of degrees of freedom N for the chi-square quantile is equal to the number of parameters of interest.

If the large-sample limit is not justified, then use e.g. Monte Carlo to get distribution of t_{μ} .

Prototype search analysis

Search for signal in a region of phase space; result is histogram of some variable x giving numbers:

$$\mathbf{n}=(n_1,\ldots,n_N)$$

Assume the n_i are Poisson distributed with expectation values

$$E[n_i] = \mu s_i + b_i$$
strength parameter

where

G. Cowan / RHUL Physics

Prototype analysis (II)

Often also have a subsidiary measurement that constrains some of the background and/or shape parameters:

$$\mathbf{m} = (m_1, \ldots, m_M)$$

Assume the m_i are Poisson distributed with expectation values

$$E[m_i] = u_i(\boldsymbol{\theta})$$
nuisance parameters ($\boldsymbol{\theta}_{s}, \boldsymbol{\theta}_{b}, b_{tot}$)

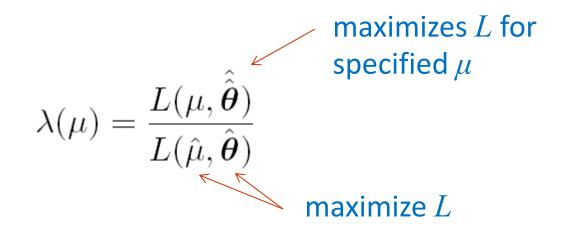
Likelihood function is

$$L(\mu, \theta) = \prod_{j=1}^{N} \frac{(\mu s_j + b_j)^{n_j}}{n_j!} e^{-(\mu s_j + b_j)} \quad \prod_{k=1}^{M} \frac{u_k^{m_k}}{m_k!} e^{-u_k}$$

G. Cowan / RHUL Physics

The profile likelihood ratio

Base significance test on the profile likelihood ratio:



Define critical region of test of μ by the region of data space that gives the lowest values of $\lambda(\mu)$.

Important advantage of profile LR is that its distribution becomes independent of nuisance parameters in large sample limit.

Test statistic for discovery

Suppose relevant alternative to background-only ($\mu = 0$) is $\mu \ge 0$. So take critical region for test of $\mu = 0$ corresponding to high q_0

and $\hat{\mu} > 0$ (data characteristic for $\mu \ge 0$).

That is, to test background-only hypothesis define statistic

$$q_0 = \begin{cases} -2\ln\lambda(0) & \hat{\mu} \ge 0\\ 0 & \hat{\mu} < 0 \end{cases}$$

i.e. here only large (positive) observed signal strength is evidence against the background-only hypothesis.

Note that even though here physically $\mu \ge 0$, we allow $\hat{\mu}$ to be negative. In large sample limit its distribution becomes Gaussian, and this will allow us to write down simple expressions for distributions of our test statistics.

Distribution of q_0 in large-sample limit

Assuming approximations valid in the large sample (asymptotic) limit, we can write down the full distribution of q_0 as

$$f(q_0|\mu') = \left(1 - \Phi\left(\frac{\mu'}{\sigma}\right)\right)\delta(q_0) + \frac{1}{2}\frac{1}{\sqrt{2\pi}}\frac{1}{\sqrt{q_0}}\exp\left[-\frac{1}{2}\left(\sqrt{q_0} - \frac{\mu'}{\sigma}\right)^2\right]$$

The special case $\mu' = 0$ is a "half chi-square" distribution:

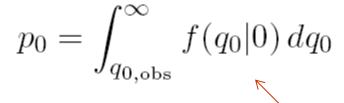
$$f(q_0|0) = \frac{1}{2}\delta(q_0) + \frac{1}{2}\frac{1}{\sqrt{2\pi}}\frac{1}{\sqrt{q_0}}e^{-q_0/2}$$

In large sample limit, $f(q_0|0)$ independent of nuisance parameters; $f(q_0|\mu')$ depends on nuisance parameters through σ .

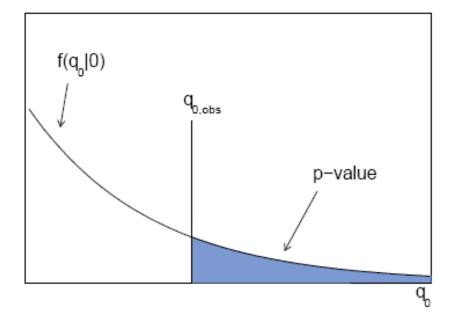
G. Cowan / RHUL Physics

p-value for discovery

Large q_0 means increasing incompatibility between the data and hypothesis, therefore *p*-value for an observed $q_{0,obs}$ is



use e.g. asymptotic formula



From *p*-value get equivalent significance,

$$Z = \Phi^{-1}(1-p)$$

Cumulative distribution of q_0 , significance

From the pdf, the cumulative distribution of q_0 is found to be

$$F(q_0|\mu') = \Phi\left(\sqrt{q_0} - \frac{\mu'}{\sigma}\right)$$

The special case $\mu' = 0$ is

$$F(q_0|0) = \Phi\left(\sqrt{q_0}\right)$$

The *p*-value of the $\mu = 0$ hypothesis is

$$p_0 = 1 - F(q_0|0)$$

Therefore the discovery significance Z is simply

$$Z = \Phi^{-1}(1 - p_0) = \sqrt{q_0}$$

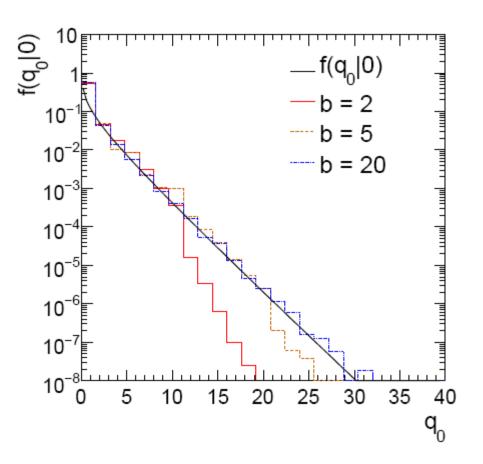
G. Cowan / RHUL Physics

Monte Carlo test of asymptotic formula

- $n \sim \text{Poisson}(\mu s + b)$
- $m \sim \text{Poisson}(\tau b)$
- $\mu =$ param. of interest
- *b* = nuisance parameter

Here take *s* known, $\tau = 1$.

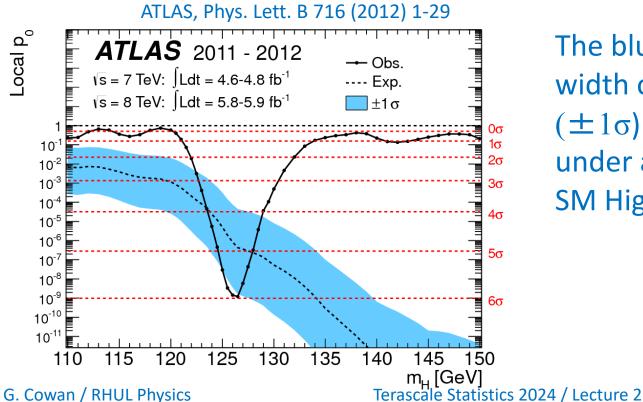
Asymptotic formula is good approximation to 5σ level ($q_0 = 25$) already for $b \sim 20$.



How to read the p_0 plot

The "local" p_0 means the *p*-value of the background-only hypothesis obtained from the test of $\mu = 0$ at each individual $m_{\rm H}$, without any correct for the Look-Elsewhere Effect.

The "Expected" (dashed) curve gives the median p_0 under assumption of the SM Higgs (μ = 1) at each $m_{\rm H}$.



The blue band gives the width of the distribution $(\pm 1\sigma)$ of significances under assumption of the SM Higgs.

Test statistic for upper limits

For purposes of setting an upper limit on μ use

$$q_{\mu} = \begin{cases} -2\ln\lambda(\mu) & \hat{\mu} \leq \mu \\ 0 & \hat{\mu} > \mu \end{cases} \quad \text{where} \quad \lambda(\mu) = \frac{L(\mu, \hat{\hat{\theta}})}{L(\hat{\mu}, \hat{\theta})}$$

I.e. when setting an upper limit, an upwards fluctuation of the data is not taken to mean incompatibility with the hypothesized μ :

From observed q_{μ} find *p*-value: $p_{\mu} = \int_{q_{\mu,\text{obs}}}^{\infty} f(q_{\mu}|\mu) dq_{\mu}$

Large sample approximation:

$$p_{\mu} = 1 - \Phi\left(\sqrt{q_{\mu}}\right)$$

To find upper limit at CL = $1-\alpha$, set $p_{\mu} = \alpha$ and solve for μ .

G. Cowan / RHUL Physics

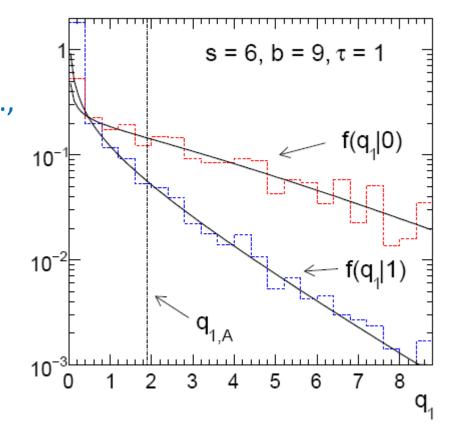
Monte Carlo test of asymptotic formulae

Consider again $n \sim \text{Poisson}(\mu s + b)$, $m \sim \text{Poisson}(\tau b)$ Use q_{μ} to find *p*-value of hypothesized μ values.

E.g. $f(q_1|1)$ for *p*-value of $\mu = 1$. Typically interested in 95% CL, i.e., *p*-value threshold = 0.05, i.e., $q_1 = 2.69$ or $Z_1 = \sqrt{q_1} = 1.64$. Median[$q_1|0$] gives "exclusion

sensitivity".

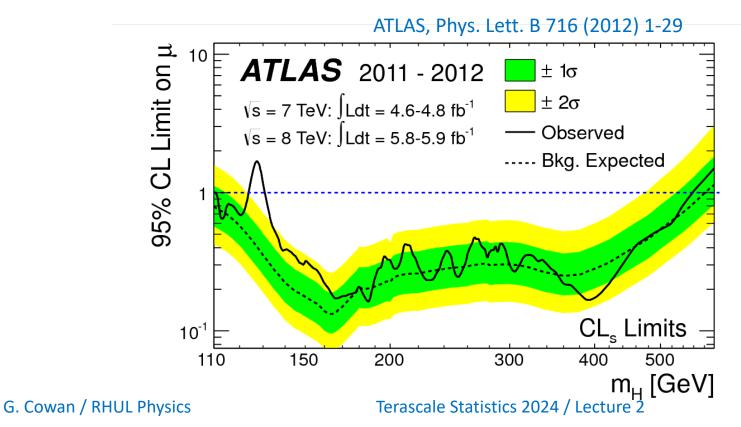
Here asymptotic formulae good for s = 6, b = 9.



How to read the green and yellow limit plots For every value of $m_{\rm H}$, find the upper limit on μ .

Also for each $m_{\rm H}$, determine the distribution of upper limits $\mu_{\rm up}$ one would obtain under the hypothesis of μ = 0.

The dashed curve is the median μ_{up} , and the green (yellow) bands give the $\pm 1\sigma$ (2σ) regions of this distribution.



Finally

Two lectures only enough for a brief introduction to:

Parameter estimation

Hypothesis tests (\rightarrow path to Machine Learning)

Limits (confidence intervals/regions)

Systematics (nuisance parameters)

No time for many other interesting topics:

Experimental sensitivity

Bayesian parameter estimation

Final thought: once the basic formalism is fixed, most of the work focuses on writing down the likelihood, e.g., $P(x|\theta)$, and including in it enough parameters to adequately describe the data (true for both Bayesian and frequentist approaches) so often best to invest most of your time with it.

G. Cowan / RHUL Physics

Extra slides

Information inequality for N parameters Suppose we have estimated N parameters $\theta = (\theta_1, ..., \theta_N)$ The Fisher information matrix is

$$I_{ij} = -E\left[\frac{\partial^2 \ln L}{\partial \theta_i \partial \theta_j}\right] = -\int \frac{\partial^2 \ln P(\mathbf{x}|\boldsymbol{\theta})}{\partial \theta_i \partial \theta_j} P(\mathbf{x}|\boldsymbol{\theta}) \, d\mathbf{x}$$

and the covariance matrix of estimators $\hat{\theta}$ is $V_{ij} = ext{cov}[\hat{ heta}_i, \hat{ heta}_j]$

The information inequality states that the matrix

$$M_{ij} = V_{ij} - \sum_{k,l} \left(\delta_{ik} + \frac{\partial b_i}{\partial \theta_k} \right) I_{kl}^{-1} \left(\delta_{lj} + \frac{\partial b_l}{\partial \theta_j} \right)$$

is positive semi-definite:

 $z^{\mathrm{T}}Mz \ge 0$ for all $z \ne 0$, diagonal elements ≥ 0

Information inequality for N parameters (2)

In practice the inequality is ~always used in the large-sample limit: bias $\rightarrow 0$ inequality \rightarrow equality, i.e, M = 0, and therefore $V^{-1} = I$

That is,
$$V_{ij}^{-1} = -E\left[\frac{\partial^2 \ln L}{\partial \theta_i \partial \theta_j}\right]$$

This can be estimated from data using $\hat{V}_{ij}^{-1} = -\frac{\partial^2 \ln L}{\partial \theta_i \partial \theta_j}\Big|_{\hat{\theta}}$

Find the matrix V^{-1} numerically (or with automatic differentiation), then invert to get the covariance matrix of the estimators

$$\widehat{V}_{ij} = \widehat{\text{cov}}[\widehat{\theta}_i, \widehat{\theta}_j]$$

Multiparameter graphical method for variances

Expand $\ln L(\theta)$ to 2nd order about MLE:

$$\ln L(\boldsymbol{\theta}) \approx \ln L(\hat{\boldsymbol{\theta}}) + \sum_{i} \frac{\partial \ln L}{\partial \theta_{i}} \Big|_{\hat{\boldsymbol{\theta}}} (\theta_{i} - \hat{\theta}_{i}) + \frac{1}{2!} \sum_{i,j} \frac{\partial^{2} \ln L}{\partial \theta_{i} \partial \theta_{j}} \Big|_{\hat{\boldsymbol{\theta}}} (\theta_{i} - \hat{\theta}_{i})(\theta_{j} - \hat{\theta}_{j})$$

$$\int_{\ln L_{\max}} zero relate to covariance matrix of$$

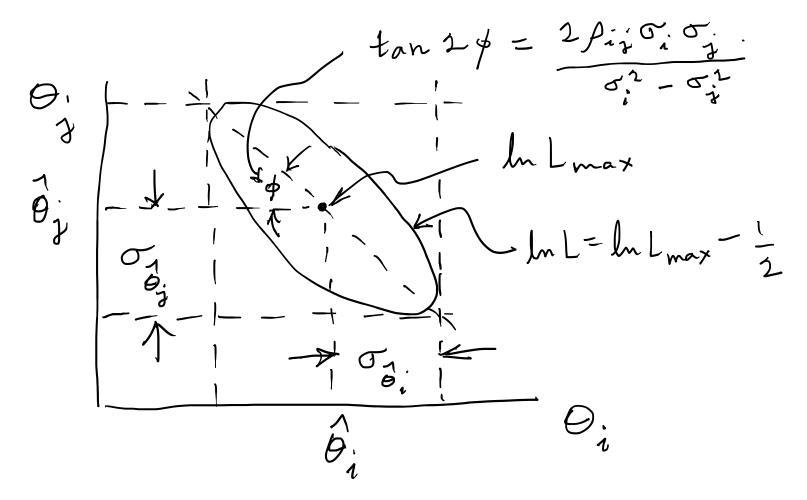
relate to covariance matrix of MLEs using information (in)equality.

Result:
$$\ln L(\boldsymbol{\theta}) = \ln L_{\max} - \frac{1}{2} \sum_{i,j} (\theta_i - \hat{\theta}_i) V_{ij}^{-1} (\theta_j - \hat{\theta}_j)$$

So the surface $\ln L(\theta) = \ln L_{\max} - \frac{1}{2}$ corresponds to

 $(\theta - \hat{\theta})^T V^{-1} (\theta - \hat{\theta}) = 1$, which is the equation of a (hyper-) ellipse.

Multiparameter graphical method (2)



Distance from MLE to tangent planes gives standard deviations.

Sensitivity for Poisson counting experiment

Count a number of events $n \sim \text{Poisson}(s+b)$, where

- s = expected number of events from signal,
- b = expected number of background events.

To test for discovery of signal compute p-value of s = 0 hypothesis,

$$p = P(n \ge n_{\text{obs}}|b) = \sum_{n=n_{\text{obs}}}^{\infty} \frac{b^n}{n!} e^{-b} = 1 - F_{\chi^2}(2b; 2n_{\text{obs}})$$

Usually convert to equivalent significance: $Z = \Phi^{-1}(1-p)$ where Φ is the standard Gaussian cumulative distribution, e.g., Z > 5 (a 5 sigma effect) means $p < 2.9 \times 10^{-7}$.

To characterize sensitivity to discovery, give expected (mean or median) Z under assumption of a given s.

G. Cowan / RHUL Physics

 s/\sqrt{b} for expected discovery significance For large s + b, $n \to x \sim \text{Gaussian}(\mu, \sigma)$, $\mu = s + b$, $\sigma = \sqrt{(s + b)}$. For observed value x_{obs} , p-value of s = 0 is $\text{Prob}(x > x_{\text{obs}} | s = 0)$,:

$$p_0 = 1 - \Phi\left(\frac{x_{\rm obs} - b}{\sqrt{b}}\right)$$

Significance for rejecting s = 0 is therefore

$$Z_0 = \Phi^{-1}(1 - p_0) = \frac{x_{\text{obs}} - b}{\sqrt{b}}$$

Expected (median) significance assuming signal rate s is

$$\mathrm{median}[Z_0|s+b] = \frac{s}{\sqrt{b}}$$

G. Cowan / RHUL Physics

Better approximation for significance

Poisson likelihood for parameter s is

 $L(s) = \frac{(s+b)^n}{n!} e^{-(s+b)}$ For now no nuisance params.

To test for discovery use profile likelihood ratio:

$$q_0 = \begin{cases} -2\ln\lambda(0) & \hat{s} \ge 0 \ , \\ 0 & \hat{s} < 0 \ . \end{cases} \qquad \lambda(s) = \frac{L(s, \hat{\hat{\theta}}(s))}{L(\hat{s}, \hat{\theta})}$$

So the likelihood ratio statistic for testing s = 0 is

$$q_0 = -2\ln\frac{L(0)}{L(\hat{s})} = 2\left(n\ln\frac{n}{b} + b - n\right) \quad \text{for } n > b, \ 0 \text{ otherwise}$$

G. Cowan / RHUL Physics

Approximate Poisson significance (continued)

For sufficiently large s + b, (use Wilks' theorem),

$$Z = \sqrt{2\left(n\ln\frac{n}{b} + b - n\right)} \quad \text{for } n > b \text{ and } Z = 0 \text{ otherwise.}$$

To find median[Z|s], let $n \rightarrow s + b$ (i.e., the Asimov data set):

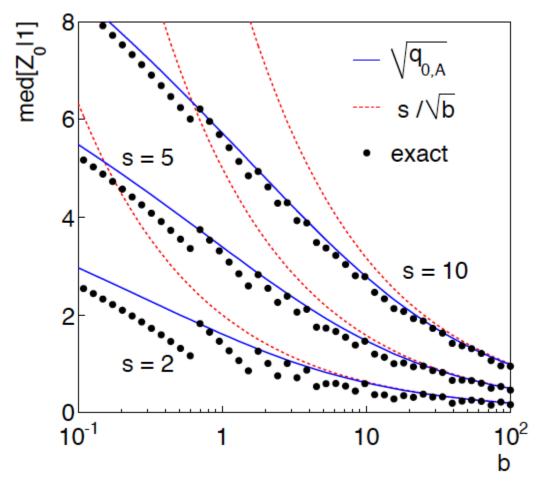
$$Z_{\rm A} = \sqrt{2\left(\left(s+b\right)\ln\left(1+\frac{s}{b}\right) - s\right)}$$

This reduces to s/\sqrt{b} for $s \ll b$.

G. Cowan / RHUL Physics

 $n \sim \text{Poisson}(s+b)$, median significance, assuming *s*, of the hypothesis s = 0

CCGV, EPJC 71 (2011) 1554, arXiv:1007.1727



"Exact" values from MC, jumps due to discrete data.

Asimov $\sqrt{q_{0,A}}$ good approx. for broad range of *s*, *b*.

 s/\sqrt{b} only good for $s \ll b$.

G. Cowan / RHUL Physics