Statistics for Particle Physics
Lecture 2
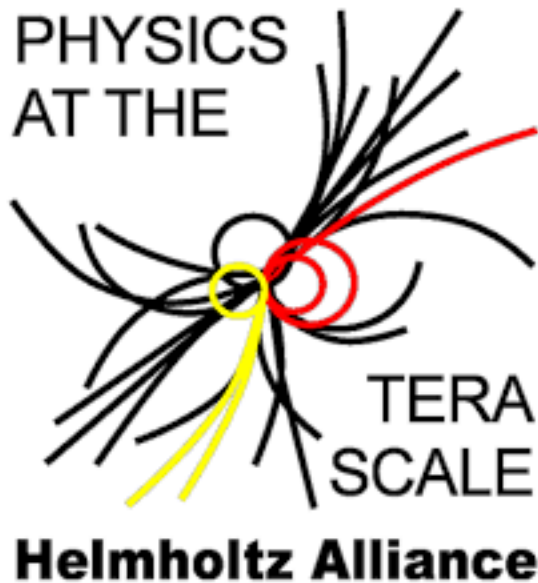
Terascale Statistics School
https://indico.desy.de/event/46667/

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Outline

Lectures/tutorials from me:

Monday 15:30   Hypothesis testing

Monday 16:30   Hypothesis testing tutorial

Tuesday 10:30   Parameter estimation tutorial

→ Tuesday 14:00   Setting limits

Wednesday 9:00   Bayesian parameter estimation

Wednesday 11:00 Errors on errors

Confidence intervals by inverting a test

In addition to a 'point estimate' of a parameter we should report an interval reflecting its statistical uncertainty.

Confidence intervals for a parameter θ can be found by defining a test of the hypothesized value θ (do this for all θ):

Specify values of the data that are 'disfavoured' by θ (critical region) such that $P($data in critical region$|\theta) \leq \alpha$ for a prespecified α, e.g., 0.05 or 0.1.

If data observed in the critical region, reject the value θ.

Now invert the test to define a confidence interval as:

set of θ values that are not rejected in a test of size α (confidence level CL is $1-\alpha$).

Relation between confidence interval and p-value

Equivalently we can consider a significance test for each hypothesized value of θ, resulting in a p-value, p_θ.

If $p_\theta \leq \alpha$, then we reject θ.

The confidence interval at CL $= 1 - \alpha$ consists of those values of θ that are not rejected.

E.g. an upper limit on θ is the greatest value for which $p_\theta > \alpha$.

In practice find by setting $p_\theta = \alpha$ and solve for θ.

For a multidimensional parameter space $\boldsymbol{\theta} = (\theta_1, \dots \theta_M)$ *use* same idea – result is a confidence "region" with boundary determined by $p_{\boldsymbol{\theta}} = \alpha$.

Coverage probability of confidence interval

If the true value of θ is rejected, then it's not in the confidence interval. The probability for this is by construction (equality for continuous data):

$$P(\text{reject } \theta | \theta) \le \alpha = \text{type-I error rate}$$

Therefore, the probability for the interval to contain or "cover" θ is

$$P(\text{conf. interval "covers" } \theta | \theta) \ge 1 \,\square\, \alpha$$

This assumes that the set of θ values considered includes the true value, i.e., it assumes the composite hypothesis $P(x|H,\theta)$.

Frequentist upper limit on Poisson parameter

Consider again the case of observing $n \sim \text{Poisson}(s + b)$.

Suppose $b = 4.5$, $n_{\text{obs}} = 5$. Find upper limit on s at 95% CL.

Relevant alternative is $s = 0$ (critical region at low n)

p-value of hypothesized s is $P(n \leq n_{\text{obs}}; s, b)$

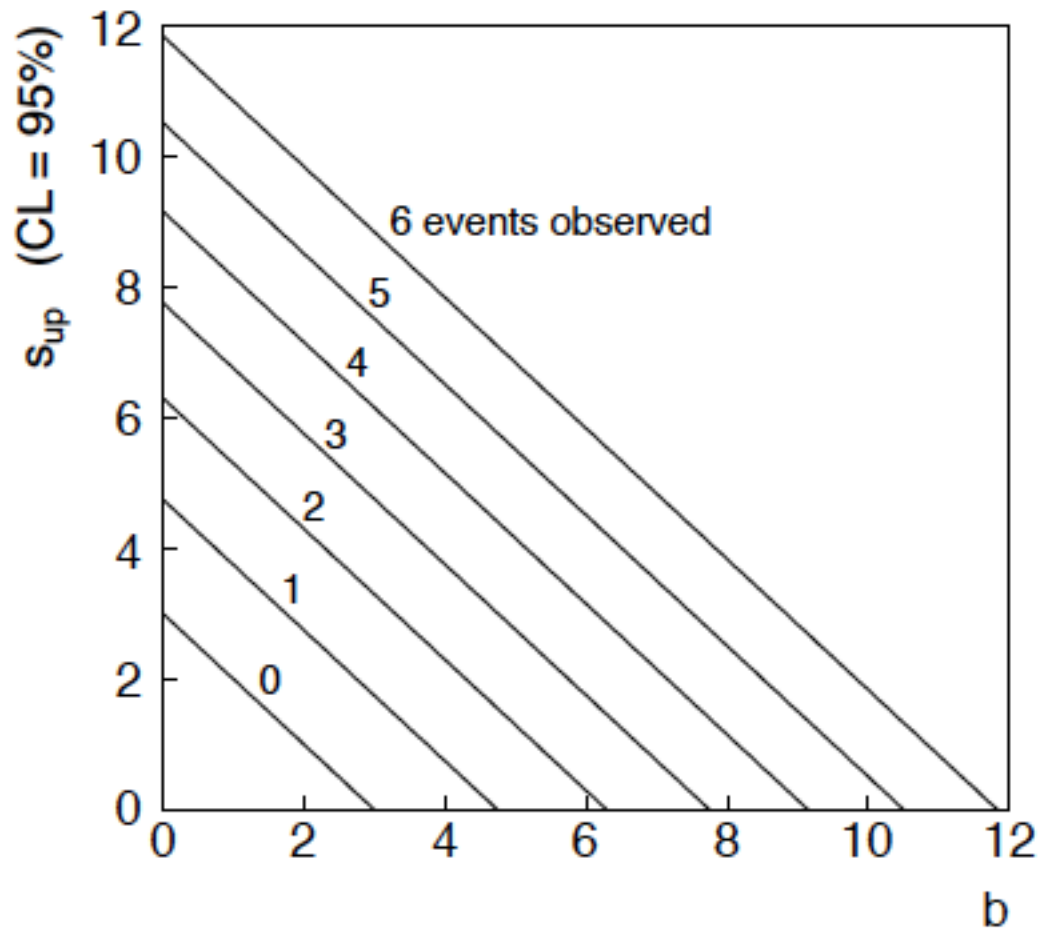Upper limit s_{up} at $\text{CL} = 1 - \alpha$ found from

$$\alpha = P(n \leq n_{\text{obs}}; s_{\text{up}}, b) = \sum_{n=0}^{n_{\text{obs}}} \frac{(s_{\text{up}} + b)^n}{n!} e^{-(s_{\text{up}}+b)}$$

$$s_{\text{up}} = \frac{1}{2} F_{\chi^2}^{-1}(1 - \alpha; 2(n_{\text{obs}} + 1)) - b$$

$$= \frac{1}{2} F_{\chi^2}^{-1}(0.95; 2(5 + 1)) - 4.5 = 6.0$$

$n \sim$ Poisson($s+b$): frequentist upper limit on s

For low fluctuation of n, formula can give negative result for s_{up}; i.e. confidence interval is empty; all values of $s \geq 0$ have $p_s \leq \alpha$.

Limits near a boundary of the parameter space

Suppose e.g. $b = 2.5$ and we observe $n = 0$.

If we choose CL $= 0.9$, we find from the formula for s_{up}

$$s_{\mathrm{up}} = -0.197 \quad (\mathrm{CL} = 0.90)$$

Physicist:

We already knew $s \geq 0$ before we started; can't use negative upper limit to report result of expensive experiment!

Statistician:

The interval is designed to cover the true value only 90% of the time — this was clearly not one of those times.

Not uncommon dilemma when testing parameter values for which one has very little experimental sensitivity, e.g., very small s.
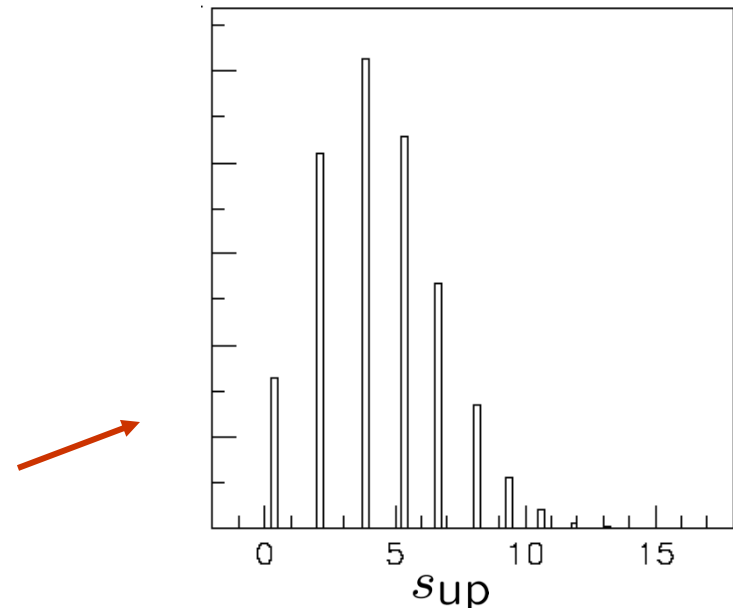
Expected limit for $s = 0$

Physicist: I should have used CL = 0.95 — then $s_{\mathrm{up}} = 0.496$

Even better: for CL = 0.917923 we get $s_{\mathrm{up}} = 10^{-4}$!

Reality check: with $b = 2.5$, typical Poisson fluctuation in n is at least $\sqrt{2.5} = 1.6$. How can the limit be so low?

Look at the mean limit for the no-signal hypothesis ($s = 0$) (sensitivity).

Distribution of 95% CL limits with $b = 2.5$, $s = 0$.
Mean upper limit = 4.44

The Bayesian approach to limits

In Bayesian statistics need to start with 'prior pdf' $\pi(\theta)$, this reflects degree of belief about θ before doing the experiment.

Bayes' theorem tells how our beliefs should be updated in light of the data x:

$$p(\theta|x) = \frac{p(x|\theta)\pi(\theta)}{\int p(x|\theta)\pi(\theta)\,d\theta} \propto p(x|\theta)\pi(\theta)$$

Integrate posterior pdf $p(\theta|x)$ to give interval with any desired probability content.

For e.g. $n \sim$ Poisson($s+b$), 95% CL upper limit on s from

$$0.95 = \int_{-\infty}^{s_{\mathrm{up}}} p(s|n)\,ds$$

Bayesian prior for Poisson parameter

Include knowledge that $s \geq 0$ by setting prior $\pi(s) = 0$ for $s < 0$.

Could try to reflect 'prior ignorance' with e.g.

$$\pi(s) = \begin{cases} 1 & s \geq 0 \\ 0 & \text{otherwise} \end{cases}$$

Not normalized; can be OK provided $p(n|s)$ dies off quickly for large s.

Not invariant under change of parameter — if we had used instead a flat prior for a nonlinear function of s, then this would imply a non-flat prior for s.

Doesn't really reflect a reasonable degree of belief, but often used as a point of reference; or viewed as a recipe for producing an interval whose frequentist properties can be studied (e.g., coverage probability, which will depend on true s).

Bayesian upper limit with flat prior for *s*

Put Poisson likelihood and flat prior into Bayes' theorem:

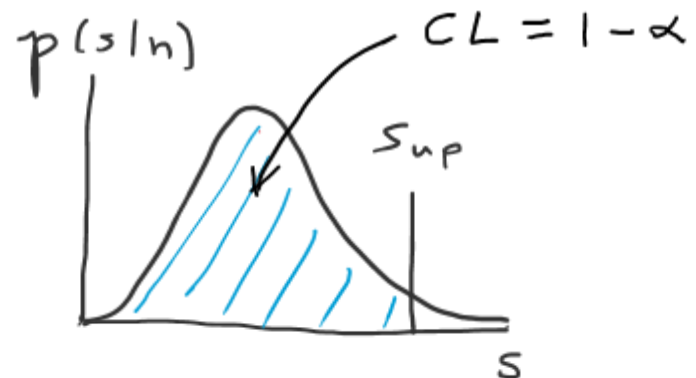$$p(s|n) \propto p(n|s)\pi(s) = \frac{(s+b)^n}{n!}e^{-(s+b)} \times 1, \quad s \geq 0$$

Normalize to unit area:

$$p(s|n) = \frac{(s+b)^n e^{-(s+b)}}{\Gamma(b, n+1)}$$

upper incomplete gamma function

Upper limit s_{up} determined by

$$1 - \alpha = \int_0^{s_{\mathrm{up}}} p(s|n)\, ds$$

Bayesian interval with flat prior for s

Solve to find limit s_{up}:

$$s_{\mathrm{up}} = \frac{1}{2} F_{\chi^2}^{-1}\left[p, 2(n+1)\right] - b$$

where

$$p = 1 - \alpha\left(1 - F_{\chi^2}\left[2b, 2(n+1)\right]\right)$$

For special case $b = 0$, Bayesian upper limit with flat prior numerically same as one-sided frequentist case ('coincidence').
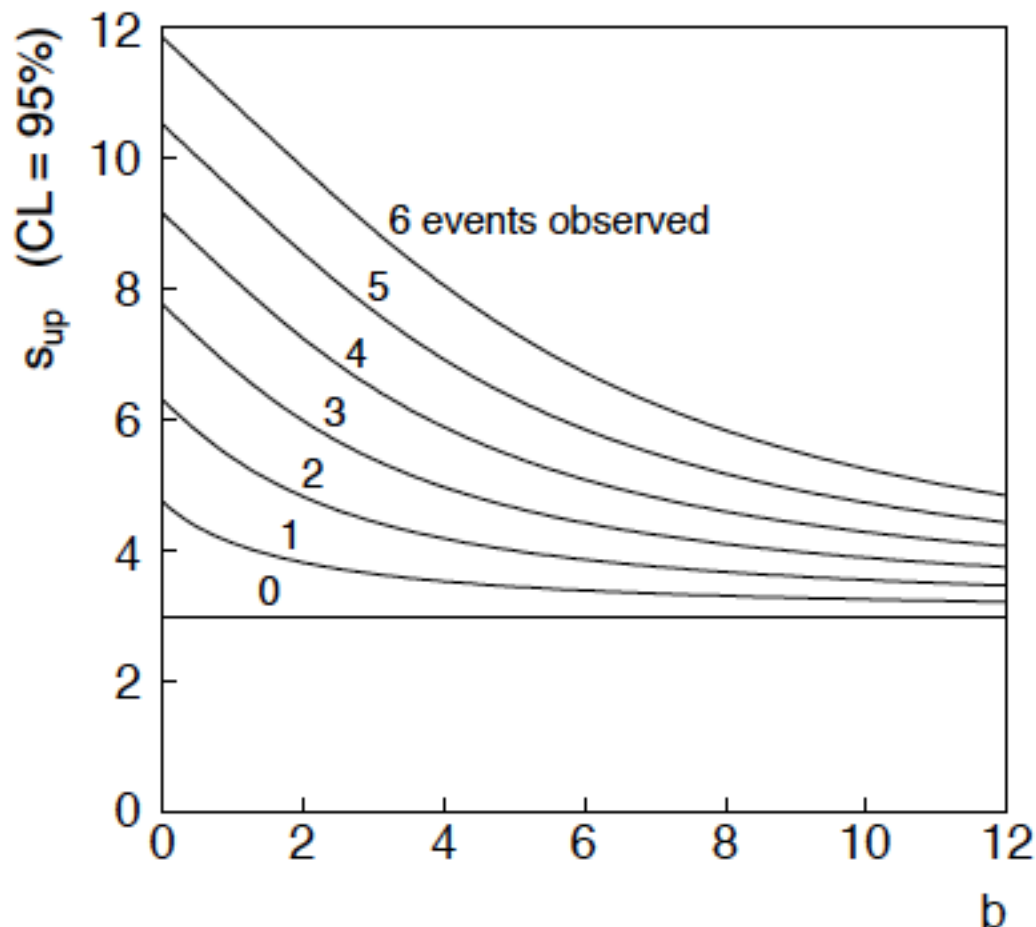
Bayesian interval with flat prior for *s*

For $b > 0$ Bayesian limit is everywhere greater than the (one sided) frequentist upper limit.

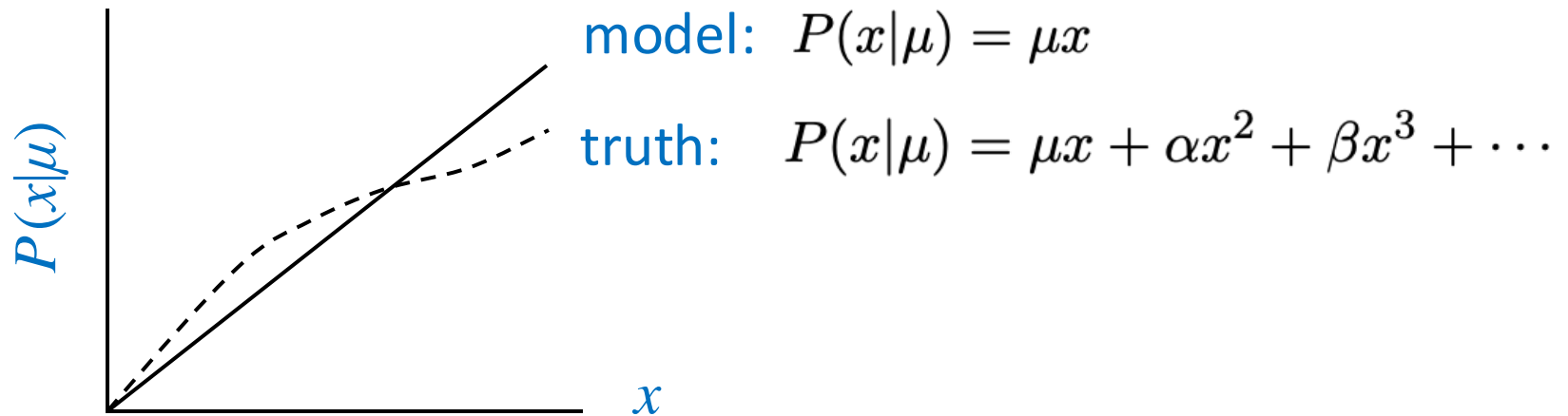For $b = 0$, Bayesian and frequentist upper limits come out equal.

Never goes negative.

Doesn't depend on b if $n = 0$.

Systematic uncertainties and nuisance parameters

In general, our model of the data is not perfect:

model: $P(x|\mu) = \mu x$

truth: $P(x|\mu) = \mu x + \alpha x^2 + \beta x^3 + \cdots$

(plot: $P(x|\mu)$ vs x)

Can improve model by including additional adjustable parameters.

$$P(x|\mu) \to P(x|\mu, \boldsymbol{\theta})$$

Nuisance parameter $\leftrightarrow$ systematic uncertainty. Some point in the parameter space of the enlarged model should be "true".

Presence of nuisance parameter decreases sensitivity of analysis to the parameter of interest (e.g., increases variance of estimate).

Profile Likelihood

Suppose we have a likelihood $L(\boldsymbol{\mu}, \boldsymbol{\theta}) = P(\boldsymbol{x}|\boldsymbol{\mu}, \boldsymbol{\theta})$ with N parameters of interest $\boldsymbol{\mu} = (\mu_1, ..., \mu_N)$ and M nuisance parameters $\boldsymbol{\theta} = (\theta_1, ..., \theta_M)$. The "profiled" (or "constrained") values of $\boldsymbol{\theta}$ are:

$$\hat{\hat{\boldsymbol{\theta}}}(\boldsymbol{\mu}) = \operatorname*{argmax}_{\boldsymbol{\theta}} L(\boldsymbol{\mu}, \boldsymbol{\theta})$$

and the profile likelihood is: $\quad L_{\mathrm{p}}(\boldsymbol{\mu}) = L(\boldsymbol{\mu}, \hat{\hat{\boldsymbol{\theta}}})$

The profile likelihood depends only on the parameters of interest; the nuisance parameters are replaced by their profiled values.

The profile likelihood can be used to obtain confidence intervals/regions for the parameters of interest in the same way as one would for all of the parameters from the full likelihood.

Profile Likelihood Ratio – Wilks theorem

Goal is to test/reject regions of $\boldsymbol{\mu}$ space (param. of interest).

Rejecting a point $\boldsymbol{\mu}$ should mean $p_{\boldsymbol{\mu}} \leq \alpha$ for all possible values of the nuisance parameters $\boldsymbol{\theta}$.

Test $\boldsymbol{\mu}$ using the "profile likelihood ratio":
$$\lambda(\boldsymbol{\mu}) = \frac{L(\boldsymbol{\mu}, \hat{\hat{\boldsymbol{\theta}}})}{L(\hat{\boldsymbol{\mu}}, \hat{\boldsymbol{\theta}})}$$

Let $t_{\boldsymbol{\mu}} = -2\ln\lambda(\boldsymbol{\mu})$. Wilks' theorem says in large-sample limit:

$$t_{\boldsymbol{\mu}} \sim \text{chi-square}(N)$$

where the number of degrees of freedom is the number of parameters of interest (components of $\boldsymbol{\mu}$). So p-value for $\boldsymbol{\mu}$ is

$$p_{\boldsymbol{\mu}} = \int_{t_{\mu,\text{obs}}}^{\infty} f(t_{\boldsymbol{\mu}}|\boldsymbol{\mu}, \boldsymbol{\theta})\, dt_{\boldsymbol{\mu}} = 1 - F_{\chi_N^2}(t_{\boldsymbol{\mu},\text{obs}})$$

Profile Likelihood Ratio – Wilks theorem (2)

Set $p_{\boldsymbol{\mu}} = \alpha$ and solve for $t_{\boldsymbol{\mu}}$ (here means observed value)

$$t_{\boldsymbol{\mu}} = F^{-1}_{\chi^2_N}(1 - \alpha) \qquad \text{which is also} \qquad t_{\boldsymbol{\mu}} = -2\ln \frac{L_{\mathrm{p}}(\boldsymbol{\mu})}{L(\hat{\boldsymbol{\mu}}, \hat{\boldsymbol{\theta}})}$$

Equate and solve for $\ln L_{\mathrm{p}}(\boldsymbol{\mu})$ gives the recipe to get confidence regions/intervals for the parameters of interest at CL $= 1 - \alpha$:

$$\ln L_{\mathrm{p}}(\boldsymbol{\mu}) = \ln L_{\mathrm{max}} - \frac{1}{2} F^{-1}_{\chi^2_N}(1 - \alpha)$$

where the number of degrees of freedom N for the chi-square quantile is equal to the number of parameters of interest.

If the large-sample limit is not justified, then use e.g. Monte Carlo to get distribution of $t_{\boldsymbol{\mu}}$.

Prototype search analysis

Search for signal in a region of phase space; result is histogram of some variable x giving numbers:

$$\mathbf{n} = (n_1, \ldots, n_N)$$

Assume the n_i are Poisson distributed with expectation values

$$E[n_i] = \mu s_i + b_i$$

strength parameter

where

$$s_i = s_{\text{tot}} \int_{\text{bin } i} f_s(x; \boldsymbol{\theta}_s) \, dx \,, \quad b_i = b_{\text{tot}} \int_{\text{bin } i} f_b(x; \boldsymbol{\theta}_b) \, dx \,.$$

signal

background

Prototype analysis (II)

Often also have a subsidiary measurement that constrains some of the background and/or shape parameters:

$$\mathbf{m} = (m_1, \ldots, m_M)$$

Assume the m_i are Poisson distributed with expectation values

$$E[m_i] = u_i(\boldsymbol{\theta})$$

nuisance parameters ($\boldsymbol{\theta}_\mathrm{s}$, $\boldsymbol{\theta}_\mathrm{b}$, b_tot)

Likelihood function is

$$L(\mu, \boldsymbol{\theta}) = \prod_{j=1}^{N} \frac{(\mu s_j + b_j)^{n_j}}{n_j!} e^{-(\mu s_j + b_j)} \prod_{k=1}^{M} \frac{u_k^{m_k}}{m_k!} e^{-u_k}$$
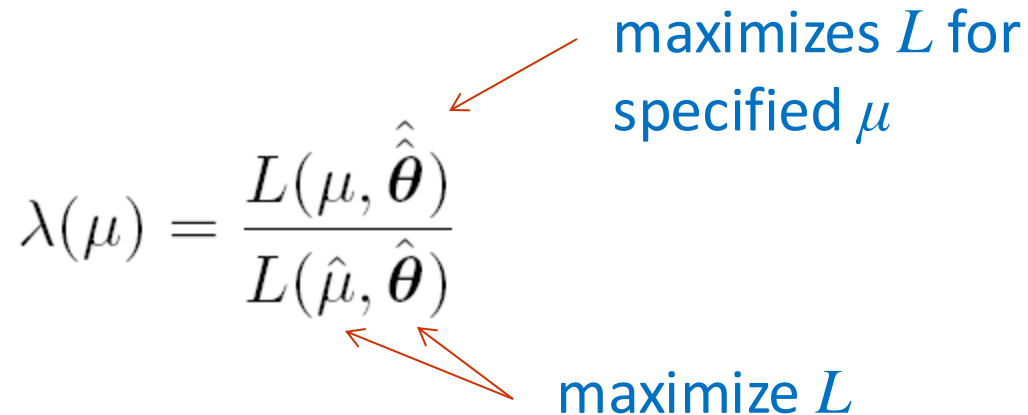
The profile likelihood ratio

Base significance test on the profile likelihood ratio:

$$\lambda(\mu) = \frac{L(\mu, \hat{\hat{\boldsymbol{\theta}}})}{L(\hat{\mu}, \hat{\boldsymbol{\theta}})}$$

maximizes L for specified μ

maximize L

Define critical region of test of μ by the region of data space that gives the lowest values of $\lambda(\mu)$.

Important advantage of profile LR is that its distribution becomes independent of nuisance parameters in large sample limit.

Test statistic for discovery

Suppose relevant alternative to background-only ($\mu = 0$) is $\mu \geq 0$.

So take critical region for test of $\mu = 0$ corresponding to high q_0 and $\hat{m} > 0$ (data characteristic for $\mu \geq 0$).

That is, to test background-only hypothesis define statistic

$$q_0 = \begin{cases} -2\ln\lambda(0) & \hat{\mu} \geq 0 \\ 0 & \hat{\mu} < 0 \end{cases}$$

i.e. here only large (positive) observed signal strength is evidence against the background-only hypothesis.

Note that even though here physically $\mu \geq 0$, we allow $\hat{m}$ to be negative. In large sample limit its distribution becomes Gaussian, and this will allow us to write down simple expressions for distributions of our test statistics.

Distribution of q_0 in large-sample limit

Assuming approximations valid in the large sample (asymptotic) limit, we can write down the full distribution of q_0 as

$$f(q_0|\mu') = \left(1 - \Phi\left(\frac{\mu'}{\sigma}\right)\right)\delta(q_0) + \frac{1}{2}\frac{1}{\sqrt{2\pi}}\frac{1}{\sqrt{q_0}}\exp\left[-\frac{1}{2}\left(\sqrt{q_0} - \frac{\mu'}{\sigma}\right)^2\right]$$

The special case $\mu' = 0$ is a "half chi-square" distribution:

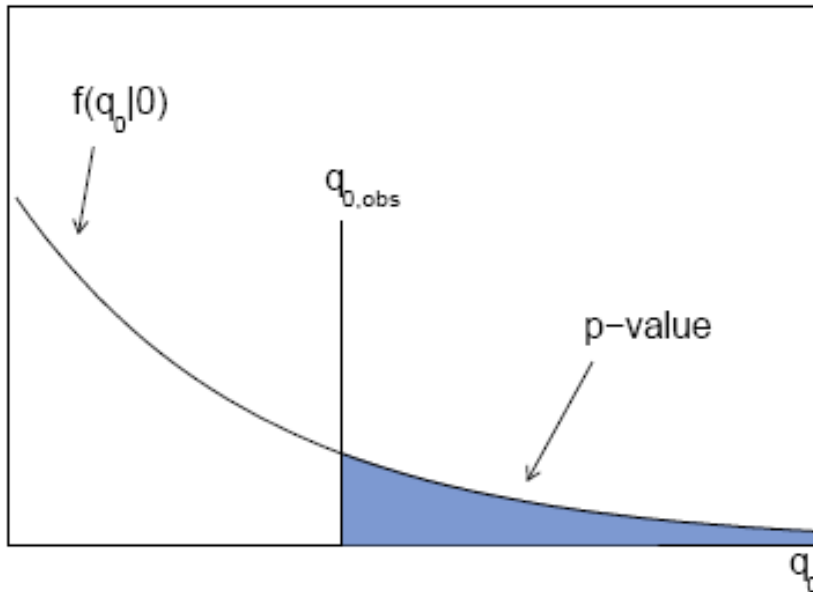$$f(q_0|0) = \frac{1}{2}\delta(q_0) + \frac{1}{2}\frac{1}{\sqrt{2\pi}}\frac{1}{\sqrt{q_0}}e^{-q_0/2}$$

In large sample limit, $f(q_0|0)$ independent of nuisance parameters; $f(q_0|\mu')$ depends on nuisance parameters through σ.

p-value for discovery

Large q_0 means increasing incompatibility between the data and hypothesis, therefore p-value for an observed $q_{0,\mathrm{obs}}$ is

$$p_0 = \int_{q_{0,\mathrm{obs}}}^{\infty} f(q_0|0)\, dq_0$$

use e.g. asymptotic formula

From p-value get equivalent significance,

$$Z = \Phi^{-1}(1 - p)$$

Cumulative distribution of q_0, significance

From the pdf, the cumulative distribution of q_0 is found to be

$$F(q_0|\mu') = \Phi\left(\sqrt{q_0} - \frac{\mu'}{\sigma}\right)$$

The special case $\mu' = 0$ is

$$F(q_0|0) = \Phi\left(\sqrt{q_0}\right)$$

The *p*-value of the $\mu = 0$ hypothesis is

$$p_0 = 1 - F(q_0|0)$$

Therefore the discovery significance Z is simply

$$Z = \Phi^{-1}(1 - p_0) = \sqrt{q_0}$$

Monte Carlo test of asymptotic formula
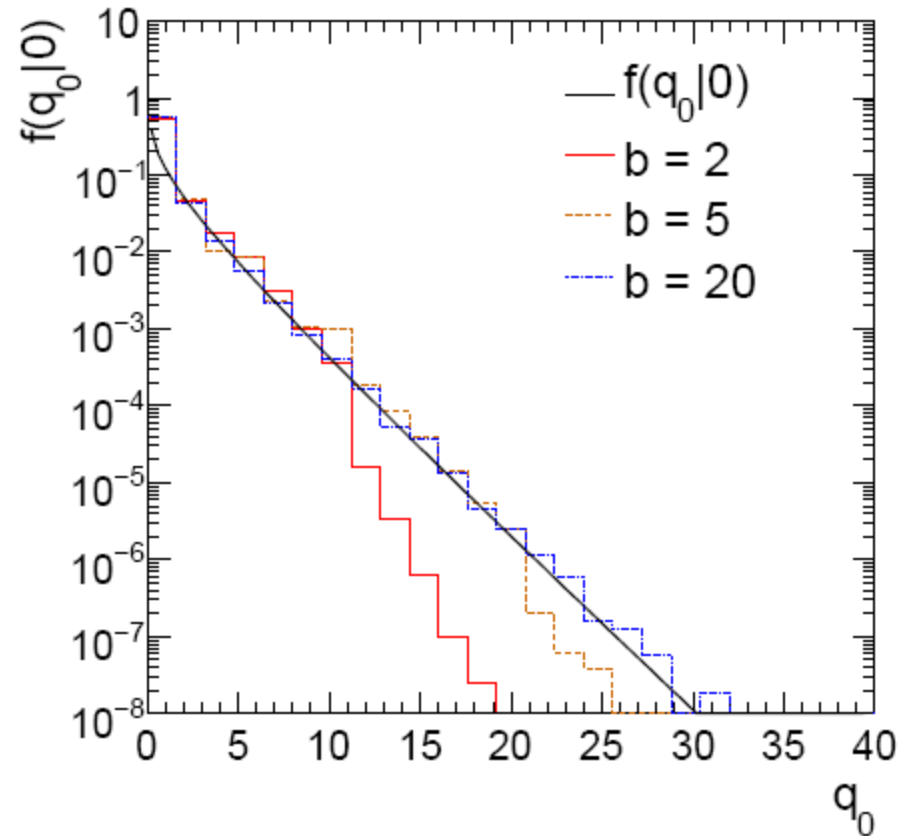
$$n \sim \text{Poisson}(\mu s + b)$$

$$m \sim \text{Poisson}(\tau b)$$

μ = param. of interest

b = nuisance parameter

Here take s known, $\tau = 1$.

Asymptotic formula is good approximation to 5σ level ($q_0 = 25$) already for $b \sim 20$.

How to read the p_0 plot

The "local" p_0 means the p-value of the background-only hypothesis obtained from the test of $\mu = 0$ at each individual m_H, without any correct for the Look-Elsewhere Effect.

The "Expected" (dashed) curve gives the median p_0 under assumption of the SM Higgs ($\mu = 1$) at each m_H.

ATLAS, Phys. Lett. B 716 (2012) 1-29
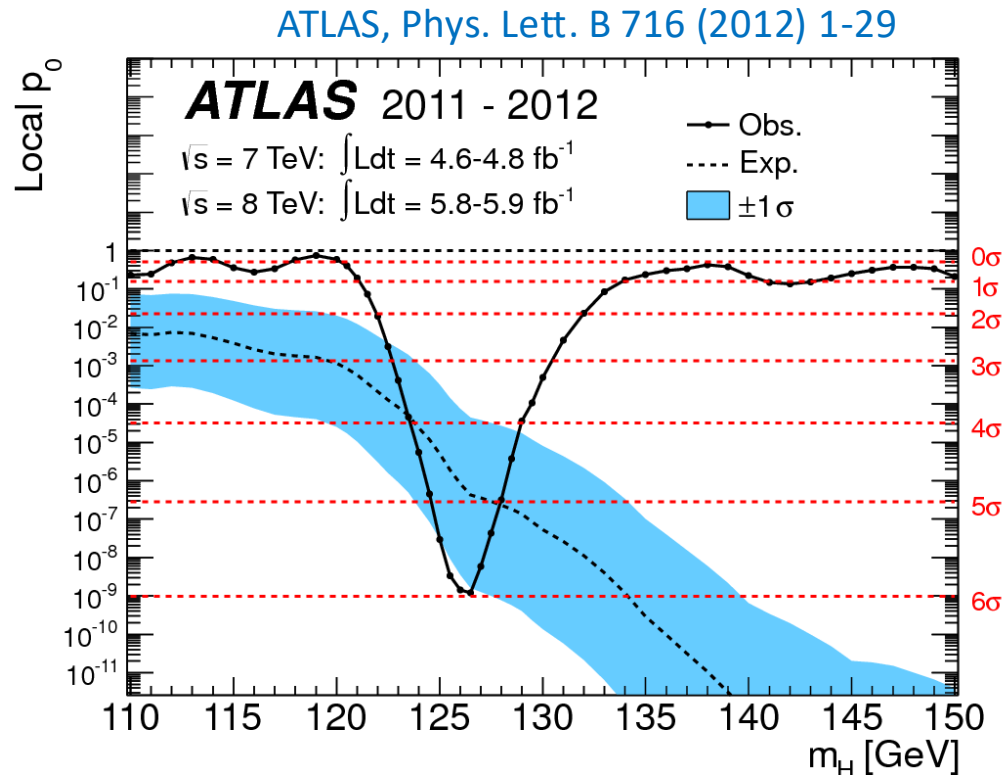
The blue band gives the width of the distribution ($\pm 1\sigma$) of significances under assumption of the SM Higgs.

Test statistic for upper limits

For purposes of setting an upper limit on μ use

$$q_\mu = \begin{cases} -2\ln\lambda(\mu) & \hat\mu \leq \mu \\ 0 & \hat\mu > \mu \end{cases} \qquad \text{where} \qquad \lambda(\mu) = \frac{L(\mu, \hat{\hat{\boldsymbol\theta}})}{L(\hat\mu, \hat{\boldsymbol\theta})}$$

I.e. when setting an upper limit, an upwards fluctuation of the data is not taken to mean incompatibility with the hypothesized μ :

From observed q_μ find p-value: $\qquad p_\mu = \int_{q_{\mu,\text{obs}}}^{\infty} f(q_\mu|\mu)\, dq_\mu$

Large sample approximation: $\qquad \boxed{p_\mu = 1 - \Phi\left(\sqrt{q_\mu}\right)}$

To find upper limit at CL = $1-\alpha$, set $p_\mu = \alpha$ and solve for μ.

Monte Carlo test of asymptotic formulae

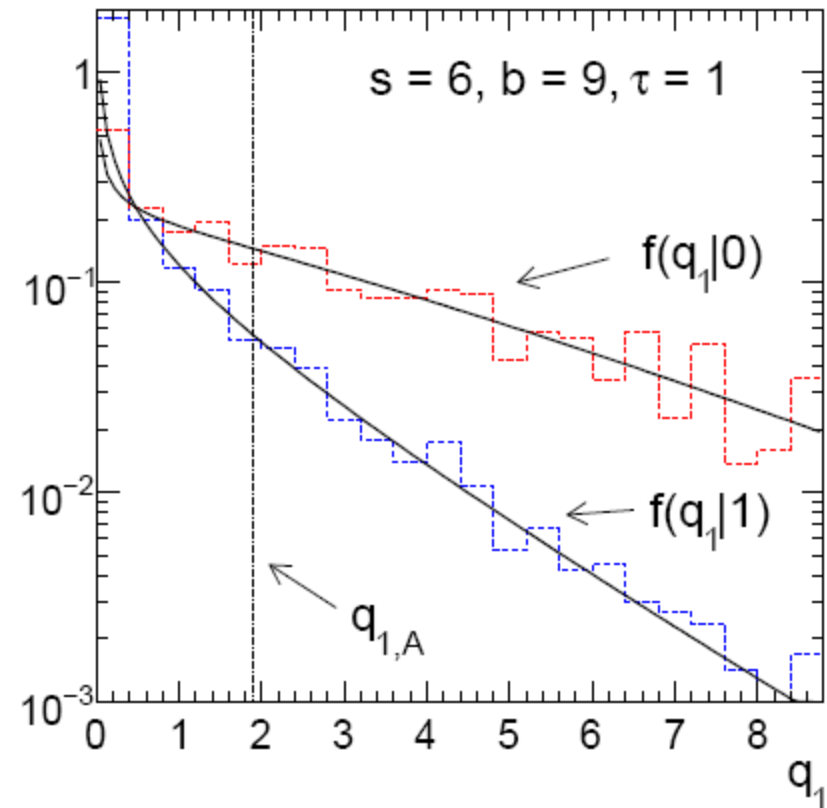Consider again $n \sim \text{Poisson}(\mu s + b)$, $m \sim \text{Poisson}(\tau b)$
Use q_μ to find p-value of hypothesized μ values.

E.g. $f(q_1|1)$ for p-value of $\mu = 1$.

Typically interested in 95% CL, i.e., p-value threshold = 0.05, i.e., $q_1 = 2.69$ or $Z_1 = \sqrt{q_1} = 1.64$.

Median$[q_1|0]$ gives "exclusion sensitivity".

Here asymptotic formulae good for $s = 6$, $b = 9$.

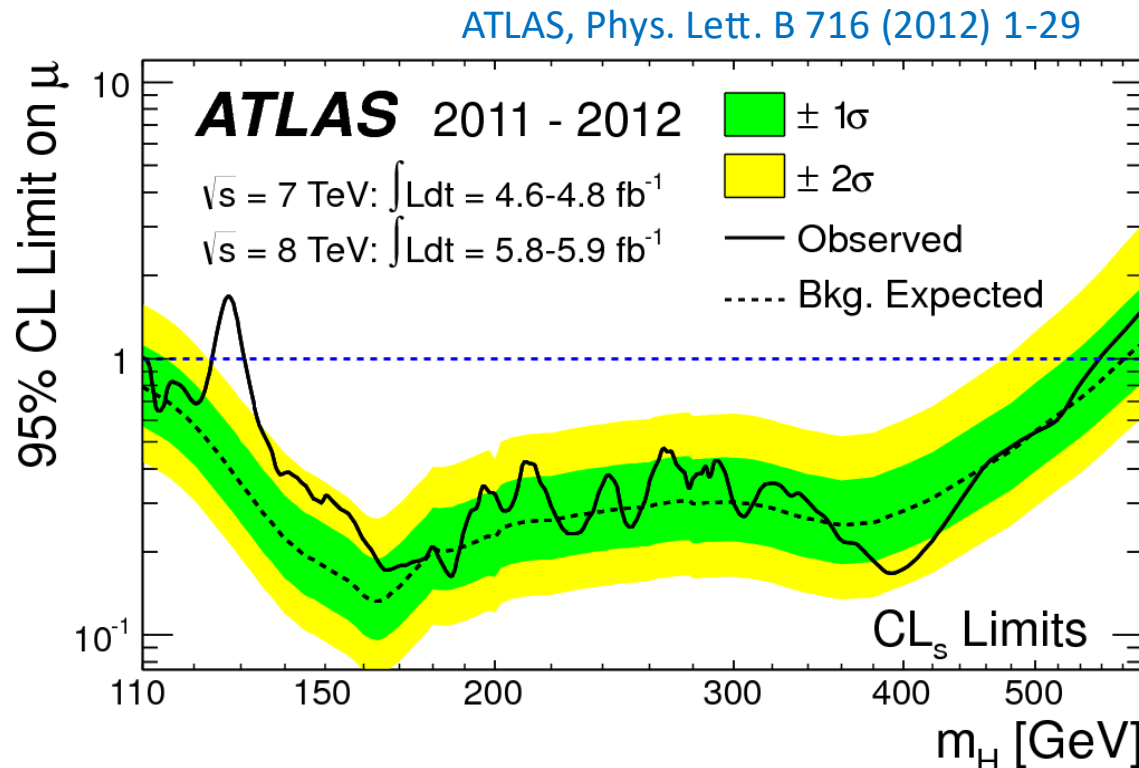$s = 6$, $b = 9$, $\tau = 1$

$f(q_1|0)$

$f(q_1|1)$

$q_{1,A}$

How to read the green and yellow limit plots

For every value of m_H, find the upper limit on μ.

Also for each m_H, determine the distribution of upper limits μ_{up} one would obtain under the hypothesis of $\mu = 0$.

The dashed curve is the median μ_{up}, and the green (yellow) bands give the $\pm\, 1\sigma$ (2σ) regions of this distribution.

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Extra slides

Approximate confidence intervals/regions from the likelihood function

Suppose we test parameter value(s) $\boldsymbol{\theta} = (\theta_1, ..., \theta_n)$ using the ratio

$$\lambda(\boldsymbol{\theta}) = \frac{L(\boldsymbol{\theta})}{L(\hat{\boldsymbol{\theta}})} \qquad\qquad 0 \leq \lambda(\boldsymbol{\theta}) \leq 1$$

Lower $\lambda(\boldsymbol{\theta})$ means worse agreement between data and hypothesized $\boldsymbol{\theta}$. Equivalently, usually define

$$t_{\boldsymbol{\theta}} = -2\ln\lambda(\boldsymbol{\theta})$$

so higher $t_{\boldsymbol{\theta}}$ means worse agreement between $\boldsymbol{\theta}$ and the data.

p-value of $\boldsymbol{\theta}$ therefore

$$p_{\boldsymbol{\theta}} = \int_{t_{\boldsymbol{\theta},\text{obs}}}^{\infty} f(t_{\boldsymbol{\theta}}|\boldsymbol{\theta})\, dt_{\boldsymbol{\theta}}$$

need pdf

Confidence region from Wilks' theorem

Wilks' theorem says (in large-sample limit and provided certain conditions hold…)

$$f(t_\theta | \boldsymbol{\theta}) \sim \chi_n^2$$

chi-square dist. with # d.o.f. = # of components in $\boldsymbol{\theta} = (\theta_1, ..., \theta_n)$.

Assuming this holds, the p-value is

$$p_\theta = 1 - F_{\chi_n^2}(t_\theta) \qquad \leftarrow \text{set equal to } \alpha$$

To find boundary of confidence region set $p_\theta = \alpha$ and solve for t_θ:

$$t_\theta = F_{\chi_n^2}^{-1}(1 - \alpha)$$

Recall also

$$t_\theta = -2 \ln \frac{L(\theta)}{L(\hat{\theta})}$$

Confidence region from Wilks' theorem (cont.)

i.e., boundary of confidence region in $\boldsymbol{\theta}$ space is where

$$\ln L(\boldsymbol{\theta}) = \ln L(\hat{\boldsymbol{\theta}}) - \tfrac{1}{2}F^{-1}_{\chi^2_n}(1-\alpha)$$

For example, for $1-\alpha = 68.3\%$ and $n=1$ parameter,

$$F^{-1}_{\chi^2_1}(0.683) = 1$$

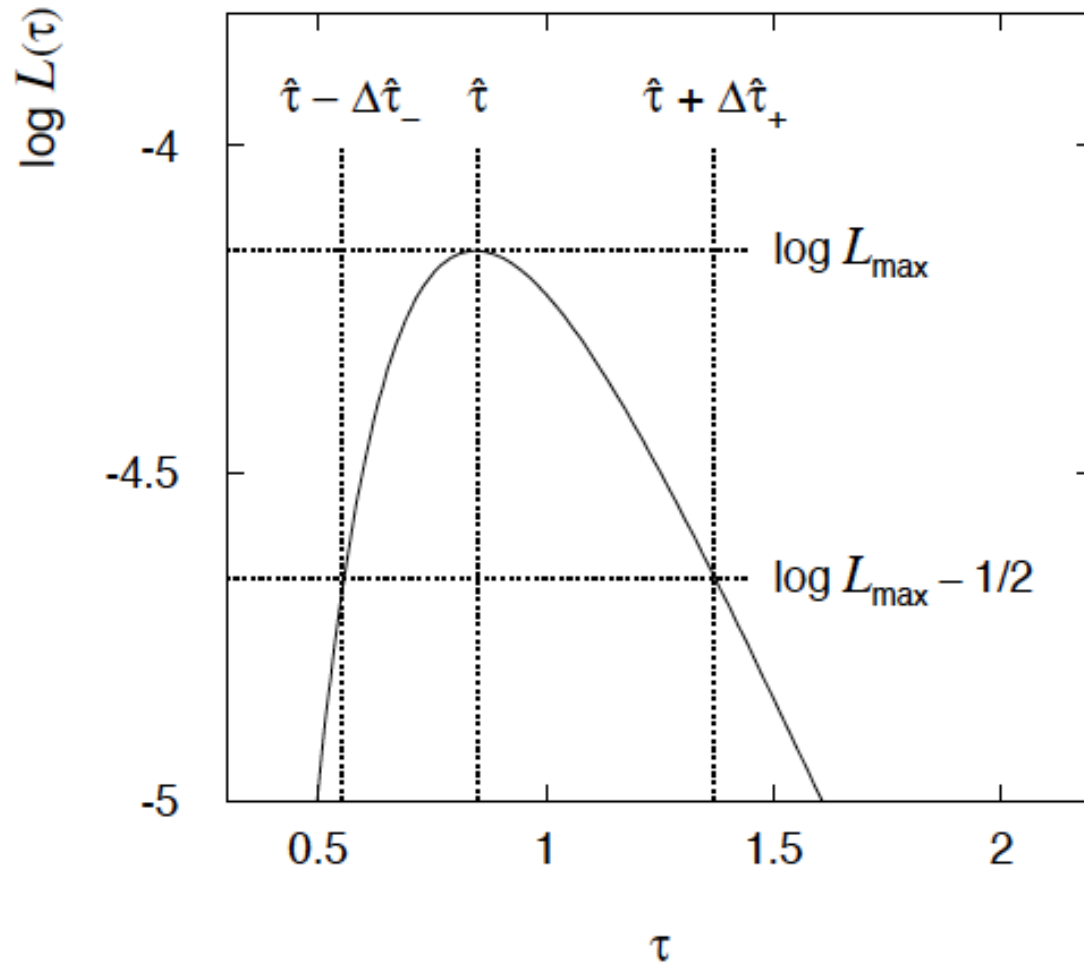and so the 68.3% confidence level interval is determined by

$$\ln L(\theta) = \ln L(\hat{\theta}) - \frac{1}{2}$$

Same as recipe for finding the estimator's standard deviation, i.e.,

$[\hat{\theta} - \sigma_{\hat{\theta}}, \hat{\theta} + \sigma_{\hat{\theta}}]$ is a 68.3% CL confidence interval.

Example of interval from $\ln L(\theta)$

For $n=1$ parameter, CL = 0.683, $Q_\alpha = 1$.

Our exponential example, now with only $n = 5$ events.

Can report ML estimate with approx. confidence interval from $\ln L_{max} - 1/2$ as "asymmetric error bar":

$$\hat{\tau} = 0.85^{+0.52}_{-0.30}$$

Multiparameter case

For increasing number of parameters, CL $= 1 - \alpha$ decreases for confidence region determined by a given

$$Q_\alpha = F^{-1}_{\chi^2_n}(1 - \alpha)$$

Q_α	$1 - \alpha$				
	$n = 1$	$n = 2$	$n = 3$	$n = 4$	$n = 5$
1.0	0.683	0.393	0.199	0.090	0.037
2.0	0.843	0.632	0.428	0.264	0.151
4.0	0.954	0.865	0.739	0.594	0.451
9.0	0.997	0.989	0.971	0.939	0.891

Multiparameter case (cont.)

Equivalently, Q_α increases with n for a given CL $= 1 - \alpha$.

$1 - \alpha$	$\hat{Q}_\alpha$				
	$n = 1$	$n = 2$	$n = 3$	$n = 4$	$n = 5$
0.683	1.00	2.30	3.53	4.72	5.89
0.90	2.71	4.61	6.25	7.78	9.24
0.95	3.84	5.99	7.82	9.49	11.1
0.99	6.63	9.21	11.3	13.3	15.1

Expected significance (sensitivity)

To characterize sensitivity to discovery, give expected (mean or median) Z under assumption of a given s.

For large $s + b$, $n \rightarrow x \sim$ Gaussian(μ, σ) , $\mu = s + b$, $\sigma = \sqrt{(s + b)}$.

For observed value x_{obs}, p-value of $s = 0$ is Prob$(x > x_{\mathrm{obs}} \mid s = 0)$,:

$$p_0 = 1 - \Phi \left(\frac{x_{\mathrm{obs}} - b}{\sqrt{b}} \right)$$

Significance for rejecting $s = 0$ is therefore

$$Z_0 = \Phi^{-1}(1 - p_0) = \frac{x_{\mathrm{obs}} - b}{\sqrt{b}}$$

Expected (median) significance assuming signal rate s is

$$\mathrm{median}[Z_0 | s + b] = \frac{s}{\sqrt{b}}$$

Better approximation for significance

Poisson likelihood for parameter *s* is

$$L(s) = \frac{(s+b)^n}{n!} e^{-(s+b)}$$

For now no nuisance params.

To test for discovery use profile likelihood ratio:

$$q_0 = \begin{cases} -2\ln\lambda(0) & \hat{s} \geq 0 \,, \\ 0 & \hat{s} < 0 \,. \end{cases}$$

$$\lambda(s) = \frac{L(s, \hat{\hat{\boldsymbol{\theta}}}(s))}{L(\hat{s}, \hat{\boldsymbol{\theta}})}$$

So the likelihood ratio statistic for testing $s = 0$ is

$$q_0 = -2\ln\frac{L(0)}{L(\hat{s})} = 2\left(n\ln\frac{n}{b} + b - n\right) \quad \text{for } n > b, \ 0 \text{ otherwise}$$

Approximate Poisson significance (continued)

For sufficiently large $s + b$, (use Wilks' theorem),

$$Z = \sqrt{2\left(n \ln \frac{n}{b} + b - n\right)} \quad \text{for } n > b \text{ and } Z = 0 \text{ otherwise.}$$

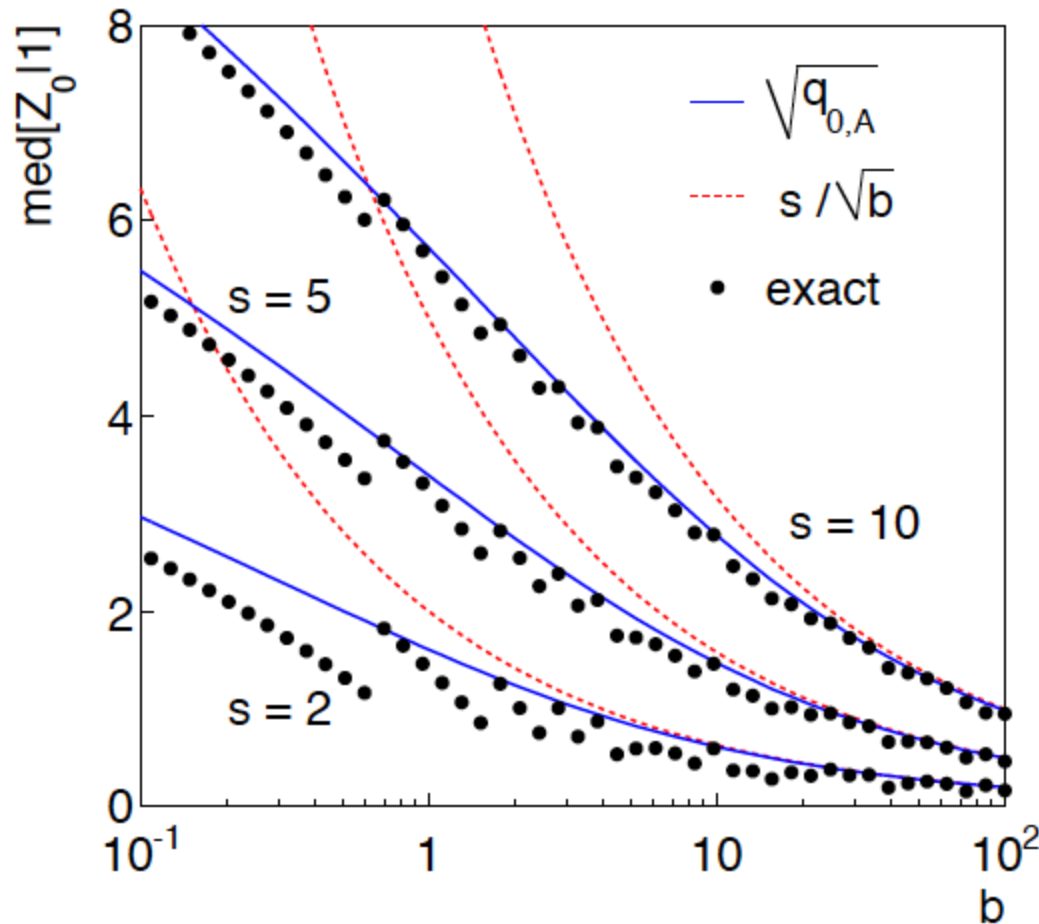To find median$[Z|s]$, let $n \to s + b$ (i.e., the Asimov data set):

$$Z_A = \sqrt{2\left((s+b) \ln\left(1 + \frac{s}{b}\right) - s\right)}$$

This reduces to $s/\sqrt{b}$ for $s \ll b$.

$n \sim$ Poisson($s+b$), median significance, assuming s, of the hypothesis $s = 0$

CCGV, EPJC 71 (2011) 1554, arXiv:1007.1727

"Exact" values from MC, jumps due to discrete data.

Asimov $\sqrt{q_{0,A}}$ good approx. for broad range of s, b.

$s/\sqrt{b}$ only good for $s \ll b$.

Extending $s/\sqrt{b}$ to case where b uncertain

The intuitive explanation of $s/\sqrt{b}$ is that it compares the signal, s, to the standard deviation of n assuming no signal, $\sqrt{b}$.

Now suppose the value of b is uncertain, characterized by a standard deviation σ_b.

A reasonable guess is to replace $\sqrt{b}$ by the quadratic sum of $\sqrt{b}$ and σ_b, i.e.,

$$\mathrm{med}[Z|s] = \frac{s}{\sqrt{b + \sigma_b^2}}$$

This has been used to optimize some analyses e.g. where σ_b cannot be neglected.

Profile likelihood with b uncertain

This is the well studied "on/off" problem: Cranmer 2005; Cousins, Linnemann, and Tucker 2008; Li and Ma 1983,...

Measure two Poisson distributed values:

$n \sim$ Poisson($s+b$)   (primary or "search" measurement)

$m \sim$ Poisson(τb)   (control measurement, τ known)

The likelihood function is

$$L(s,b) = \frac{(s+b)^n}{n!} e^{-(s+b)} \frac{(\tau b)^m}{m!} e^{-\tau b}$$

Use this to construct profile likelihood ratio (b is nuisance parameter):

$$\lambda(0) = \frac{L(0, \hat{\hat{b}}(0))}{L(\hat{s}, \hat{b})}$$

Ingredients for profile likelihood ratio

To construct profile likelihood ratio from this need estimators:

$$\hat{s} = n - m/\tau \,,$$

$$\hat{b} = m/\tau \,,$$

$$\hat{\hat{b}}(s) = \frac{n + m - (1 + \tau)s + \sqrt{(n + m - (1 + \tau)s)^2 + 4(1 + \tau)sm}}{2(1 + \tau)} \,.$$

and in particular to test for discovery ($s = 0$),

$$\hat{\hat{b}}(0) = \frac{n + m}{1 + \tau}$$

Asymptotic significance

Use profile likelihood ratio for q_0, and then from this get discovery significance using asymptotic approximation (Wilks' theorem):

$$Z = \sqrt{q_0}$$

$$= \left[-2 \left(n \ln \left[\frac{n+m}{(1+\tau)n} \right] + m \ln \left[\frac{\tau(n+m)}{(1+\tau)m} \right] \right) \right]^{1/2}$$

for $n > \hat{b}$ and $Z = 0$ otherwise.

Essentially same as in:

Robert D. Cousins, James T. Linnemann and Jordan Tucker, NIM A 595 (2008) 480–501; arXiv:physics/0702156.

Tipei Li and Yuqian Ma, Astrophysical Journal 272 (1983) 317–324.

Asimov approximation for median significance

To get median discovery significance, replace n, m by their expectation values assuming background-plus-signal model:

$$n \to s + b$$

$$m \to \tau b$$

$$Z_A = \left[-2 \left((s+b) \ln \left[\frac{s + (1+\tau)b}{(1+\tau)(s+b)} \right] + \tau b \ln \left[1 + \frac{s}{(1+\tau)b} \right] \right) \right]^{1/2}$$

Or use the variance of $\hat{b} = m/\tau$, $V[\hat{b}] \equiv \sigma_b^2 = \dfrac{b}{\tau}$, to eliminate τ:

$$Z_A = \left[2 \left((s+b) \ln \left[\frac{(s+b)(b+\sigma_b^2)}{b^2 + (s+b)\sigma_b^2} \right] - \frac{b^2}{\sigma_b^2} \ln \left[1 + \frac{\sigma_b^2 s}{b(b+\sigma_b^2)} \right] \right) \right]^{1/2}$$
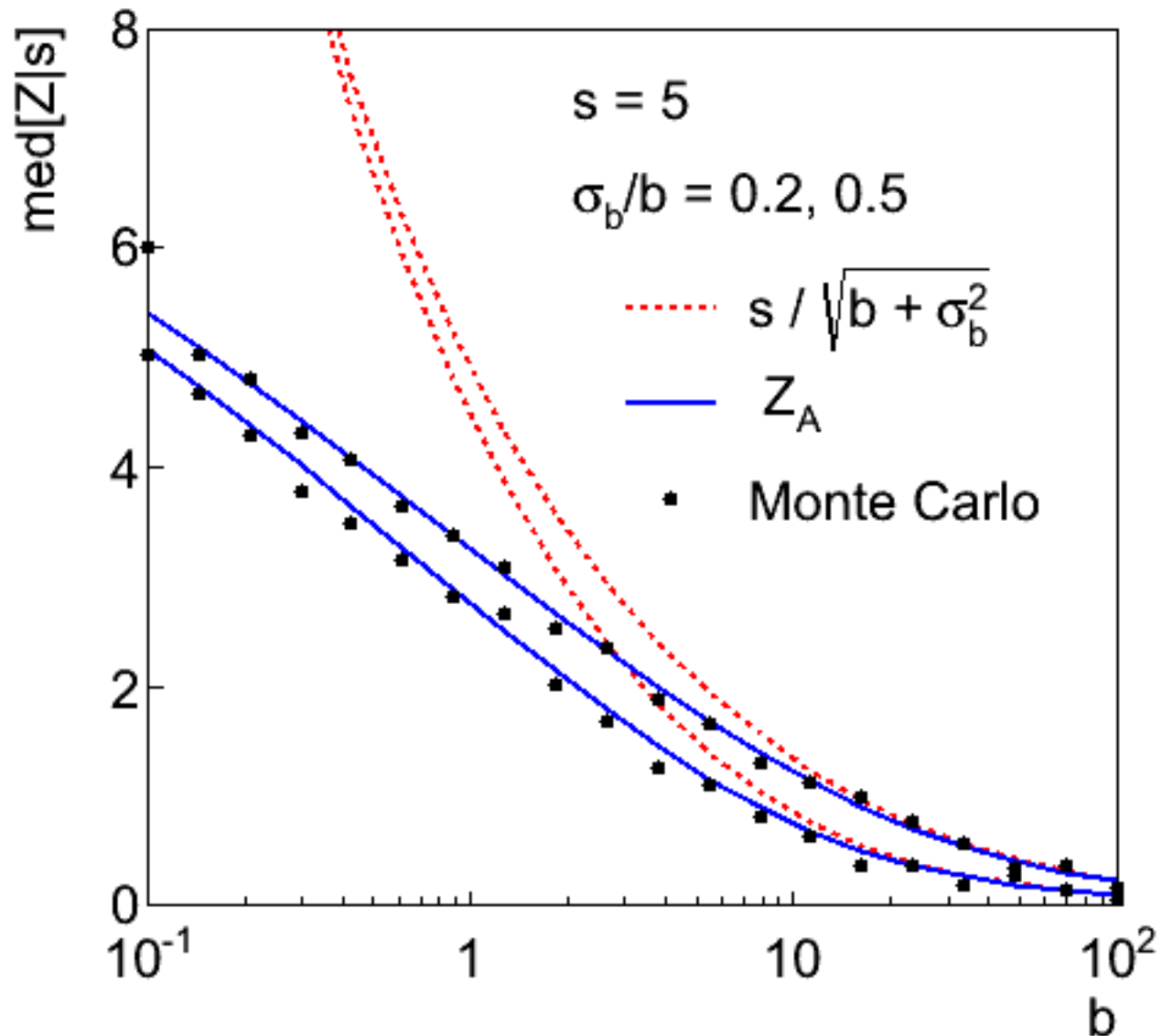
Limiting cases

Expanding the Asimov formula in powers of s/b and $\sigma_b{}^2/b$ $(= 1/\tau)$ gives

$$Z_{\mathrm{A}} = \frac{s}{\sqrt{b + \sigma_b^2}} \left(1 + \mathcal{O}(s/b) + \mathcal{O}(\sigma_b^2/b)\right)$$

So the "intuitive" formula can be justified as a limiting case of the significance from the profile likelihood ratio test evaluated with the Asimov data set.

Testing the formulae: $s = 5$

$s = 5$

$\sigma_b/b = 0.2, 0.5$

$------ \quad s / \sqrt{b + \sigma_b^2}$

$\quad\quad Z_A$

$\bullet \quad$ Monte Carlo

Using sensitivity to optimize a cut
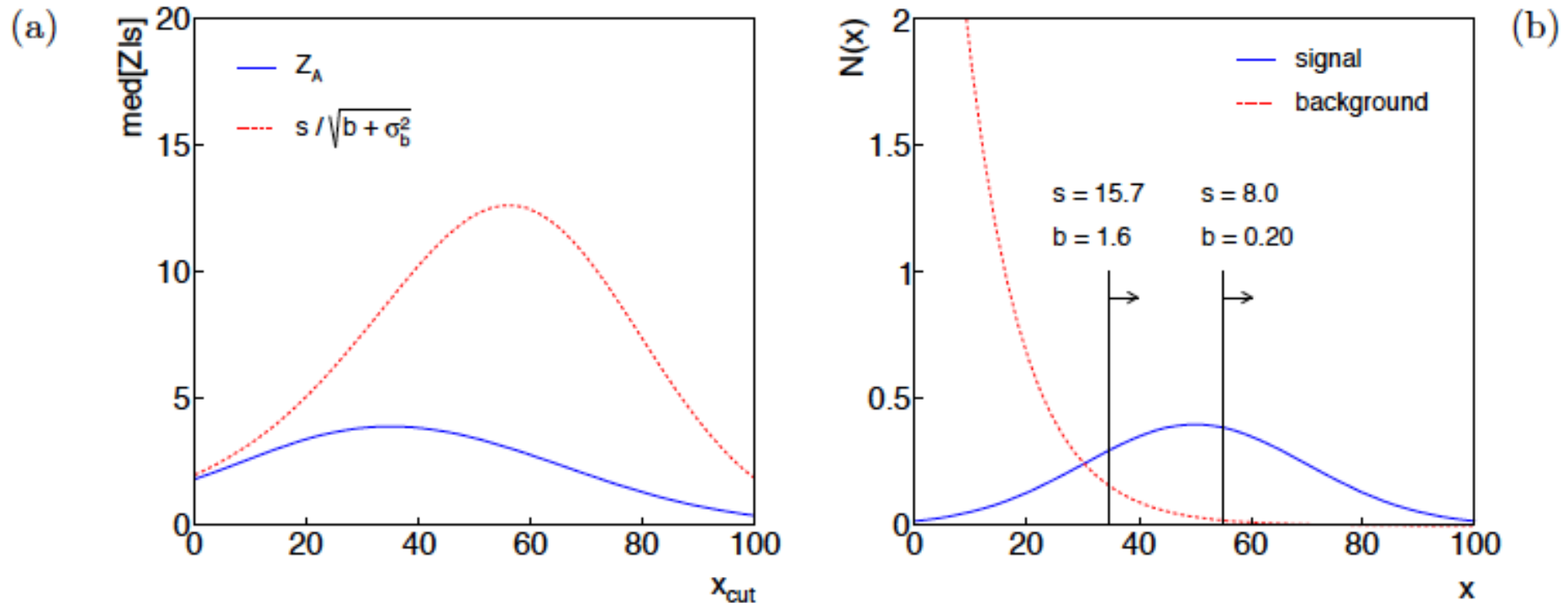

Figure 1: (a) The expected significance as a function of the cut value x_{cut}; (b) the distributions of signal and background with the optimal cut value indicated.

Summary on discovery sensitivity

Simple formula for expected discovery significance based on profile likelihood ratio test and Asimov approximation:

$$Z_A = \left[2 \left((s+b) \ln \left[\frac{(s+b)(b+\sigma_b^2)}{b^2 + (s+b)\sigma_b^2} \right] - \frac{b^2}{\sigma_b^2} \ln \left[1 + \frac{\sigma_b^2 s}{b(b+\sigma_b^2)} \right] \right) \right]^{1/2}$$

For large b, all formulae OK.

For small b, $s/\sqrt{b}$ and $s/\sqrt{(b+\sigma_b^2)}$ overestimate the significance.

Could be important in optimization of searches with low background.

Formula maybe also OK if model is not simple on/off experiment, e.g., several background control measurements (check this).