

Statistical Methods for Particle Physics

Lecture 2: Discovery and Limits

<http://indico.ihep.ac.cn/event/5966/>



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Outline

Lecture 1: Introduction and review of fundamentals

Probability, random variables, pdfs

Parameter estimation, maximum likelihood

Statistical tests

→ Lecture 2: Discovery and Limits

Comments on multivariate methods (brief)

p -values

Testing the background-only hypothesis: discovery

Testing signal hypotheses: setting limits

Lecture 3: Systematic uncertainties and further topics

Nuisance parameters (Bayesian and frequentist)

Experimental sensitivity

The look-elsewhere effect

Comment on Multivariate Methods

Naively one might choose the input variables for a multivariate analysis to be those that, by themselves, give some discrimination between signal and background.

The following simple example shows that this is not always the best approach. Variables related to the quality of an event's reconstruction can help in a multivariate analysis, even if that variable by itself gives no discrimination between event types.

A simple example (2D)

Consider two variables, x_1 and x_2 , and suppose we have formulas for the joint pdfs for both signal (s) and background (b) events (in real problems the formulas are usually not available).

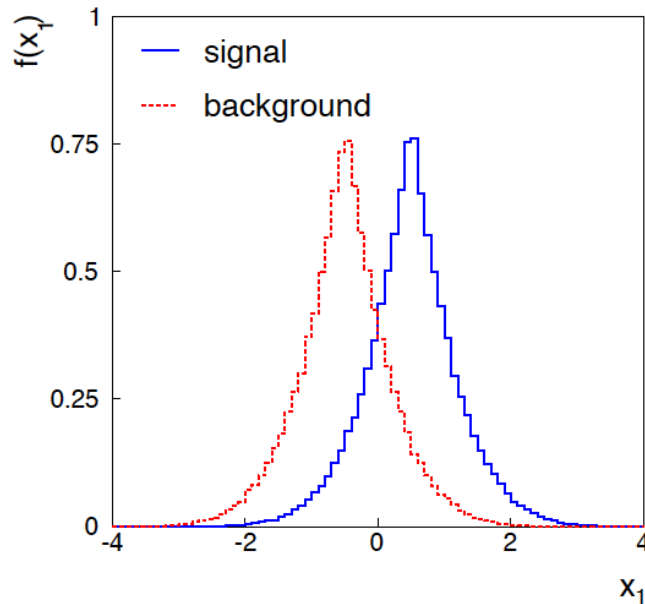
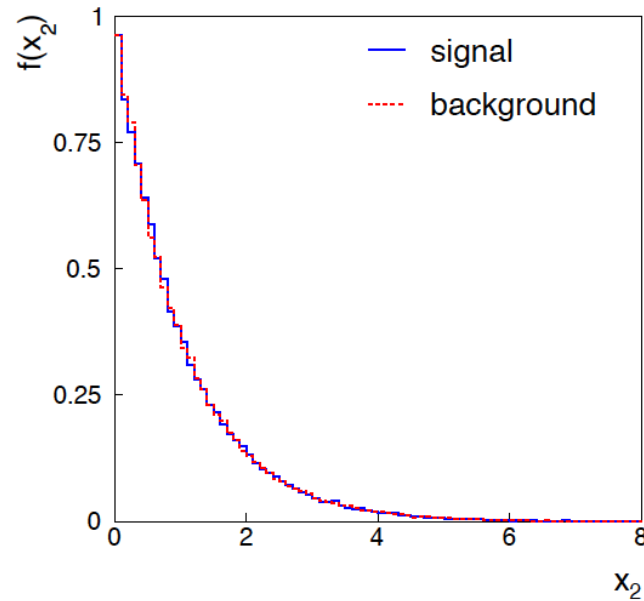
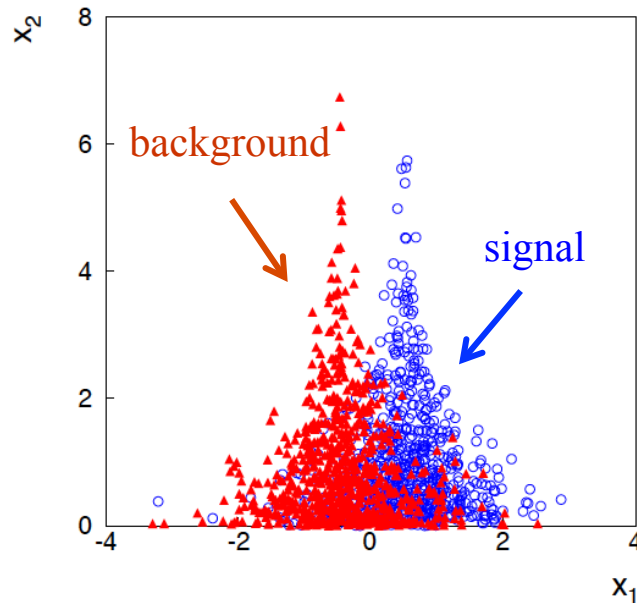
$f(x_1|x_2) \sim$ Gaussian, different means for s/b,
Gaussians have same σ , which depends on x_2 ,
 $f(x_2) \sim$ exponential, same for both s and b,
 $f(x_1, x_2) = f(x_1|x_2)f(x_2)$:

$$f(x_1, x_2|s) = \frac{1}{\sqrt{2\pi}\sigma(x_2)} e^{-(x_1 - \mu_s)^2 / 2\sigma^2(x_2)} \frac{1}{\lambda} e^{-x_2/\lambda}$$

$$f(x_1, x_2|b) = \frac{1}{\sqrt{2\pi}\sigma(x_2)} e^{-(x_1 - \mu_b)^2 / 2\sigma^2(x_2)} \frac{1}{\lambda} e^{-x_2/\lambda}$$

$$\sigma(x_2) = \sigma_0 e^{-x_2/\xi}$$

Joint and marginal distributions of x_1, x_2



Distribution $f(x_2)$ same for s, b.

So does x_2 help discriminate between the two event types?

Likelihood ratio for 2D example

Neyman-Pearson lemma says best critical region for classification is determined by the likelihood ratio:

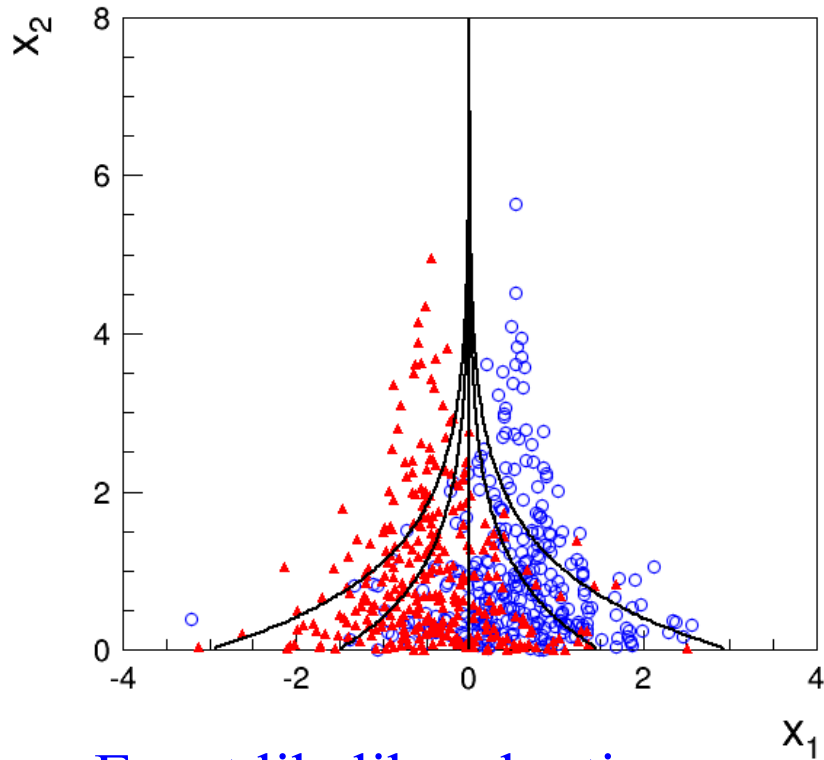
$$t(x_1, x_2) = \frac{f(x_1, x_2|s)}{f(x_1, x_2|b)}$$

Equivalently we can use any monotonic function of this as a test statistic, e.g.,

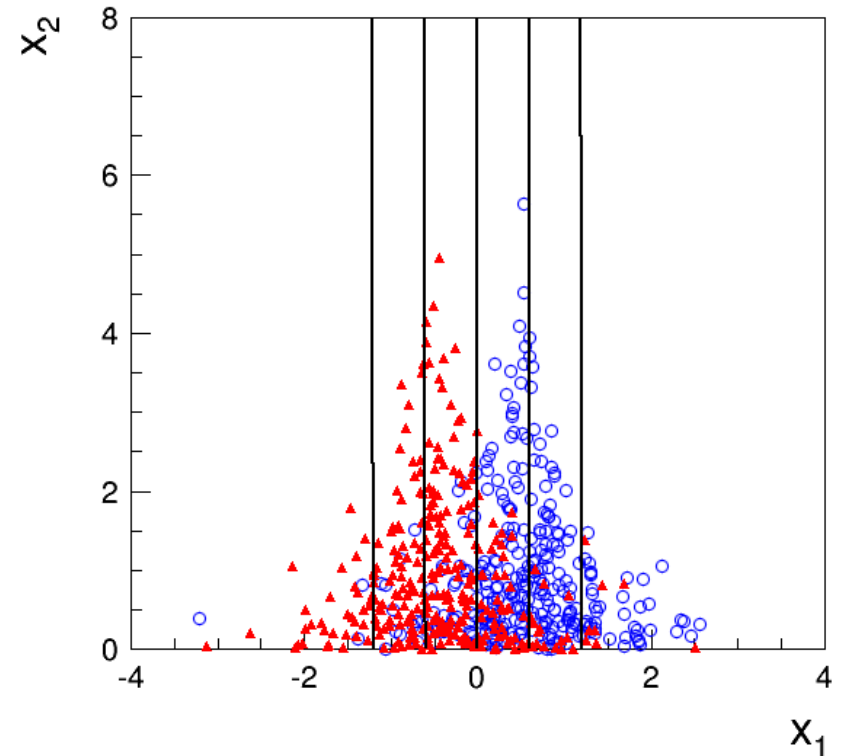
$$\ln t = \frac{\frac{1}{2}(\mu_b^2 - \mu_s^2) + (\mu_s - \mu_b)x_1}{\sigma_0^2 e^{-2x_2/\xi}}$$

Boundary of optimal critical region will be curve of constant $\ln t$, and this depends on x_2 !

Contours of constant MVA output

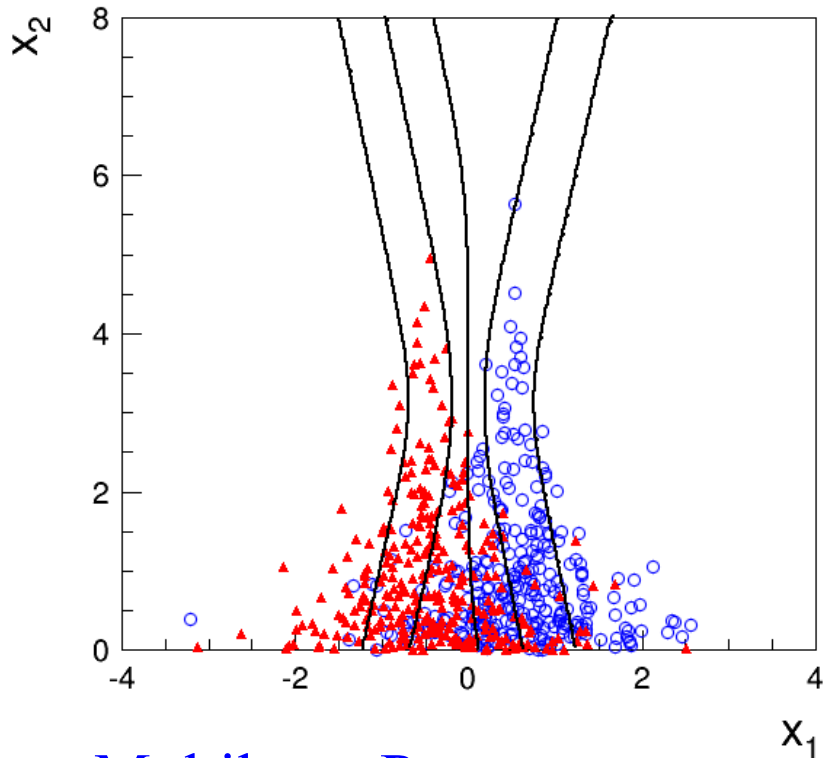


Exact likelihood ratio

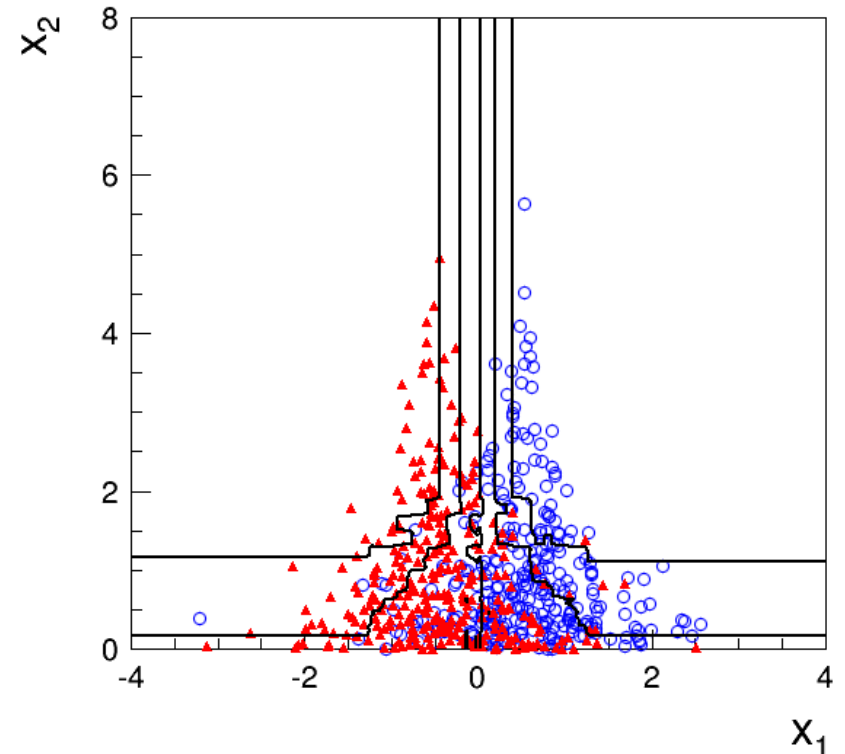


Fisher discriminant

Contours of constant MVA output



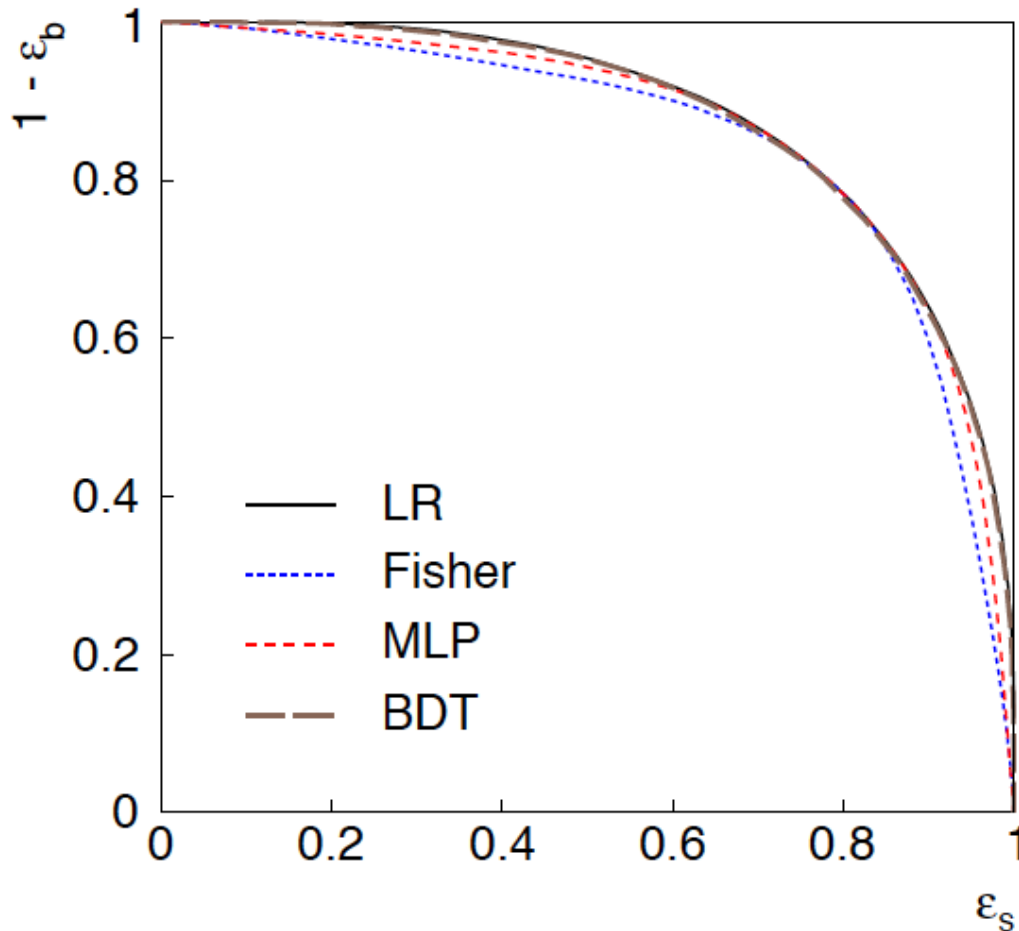
Multilayer Perceptron
1 hidden layer with 2 nodes



Boosted Decision Tree
200 iterations (AdaBoost)

Training samples: 10^5 signal and 10^5 background events

ROC curve



ROC = “receiver operating characteristic” (term from signal processing).

Shows (usually) background rejection ($1 - \epsilon_b$) versus signal efficiency ϵ_s .

Higher curve is better; usually analysis focused on a small part of the curve.

2D Example: discussion

Even though the distribution of x_2 is same for signal and background, x_1 and x_2 are not independent, so using x_2 as an input variable helps.

Here we can understand why: high values of x_2 correspond to a smaller σ for the Gaussian of x_1 . So high x_2 means that the value of x_1 was well measured.

If we don't consider x_2 , then all of the x_1 measurements are lumped together. Those with large σ (low x_2) “pollute” the well measured events with low σ (high x_2).

Often in HEP there may be variables that are characteristic of how well measured an event is (region of detector, number of pile-up vertices,...). Including these variables in a multivariate analysis preserves the information carried by the well-measured events, leading to improved performance.

Testing significance / goodness-of-fit

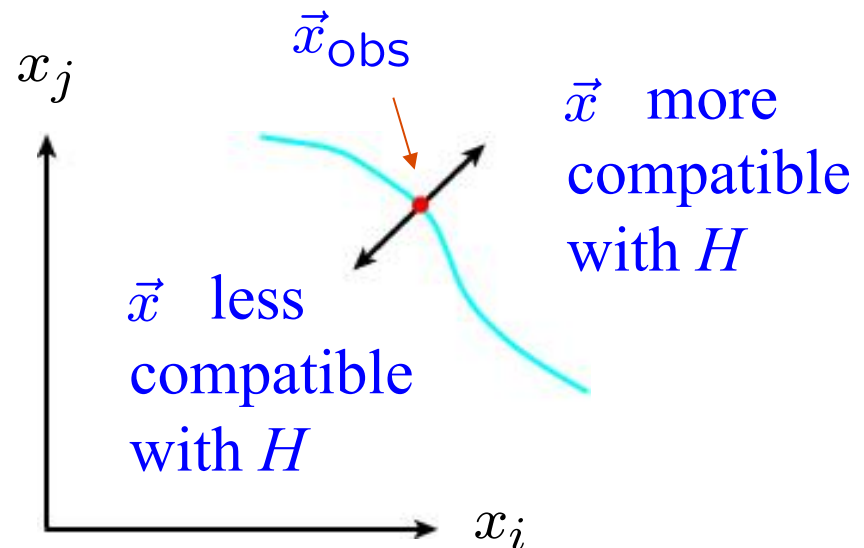
Suppose hypothesis H predicts pdf $f(\vec{x}|H)$ for a set of observations $\vec{x} = (x_1, \dots, x_n)$.

We observe a single point in this space: \vec{x}_{obs}

What can we say about the validity of H in light of the data?

Decide what part of the data space represents less compatibility with H than does the point \vec{x}_{obs} .

Note – “less compatible with H ” means “more compatible with some alternative H' ”.



p-values

Express ‘goodness-of-fit’ by giving the *p*-value for *H*:

p = probability, under assumption of *H*, to observe data with equal or lesser compatibility with *H* relative to the data we got.



This is not the probability that *H* is true!

In frequentist statistics we don’t talk about $P(H)$ (unless *H* represents a repeatable observation). In Bayesian statistics we do; use Bayes’ theorem to obtain

$$P(H|\vec{x}) = \frac{P(\vec{x}|H)\pi(H)}{\int P(\vec{x}|H)\pi(H) dH}$$

where $\pi(H)$ is the prior probability for *H*.

For now stick with the frequentist approach; result is *p*-value, regrettably easy to misinterpret as $P(H)$.

Distribution of the p -value

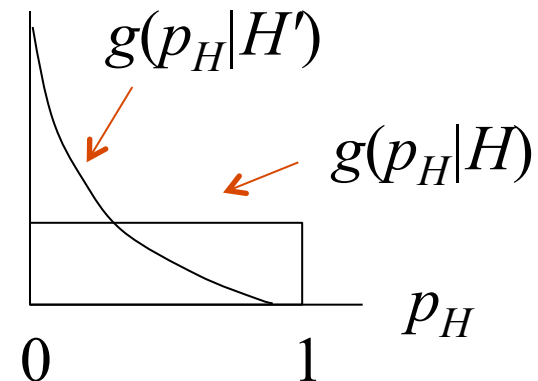
The p -value is a function of the data, and is thus itself a random variable with a given distribution. Suppose the p -value of H is found from a test statistic $t(\mathbf{x})$ as

$$p_H = \int_t^\infty f(t'|H) dt'$$

The pdf of p_H under assumption of H is

$$g(p_H|H) = \frac{f(t|H)}{|\partial p_H / \partial t|} = \frac{f(t|H)}{f(t|H)} = 1 \quad (0 \leq p_H \leq 1)$$

In general for continuous data, under assumption of H , $p_H \sim \text{Uniform}[0,1]$ and is concentrated toward zero for Some class of relevant alternatives.



Using a p -value to define test of H_0

One can show the distribution of the p -value of H , under assumption of H , is uniform in $[0,1]$.

So the probability to find the p -value of H_0 , p_0 , less than α is

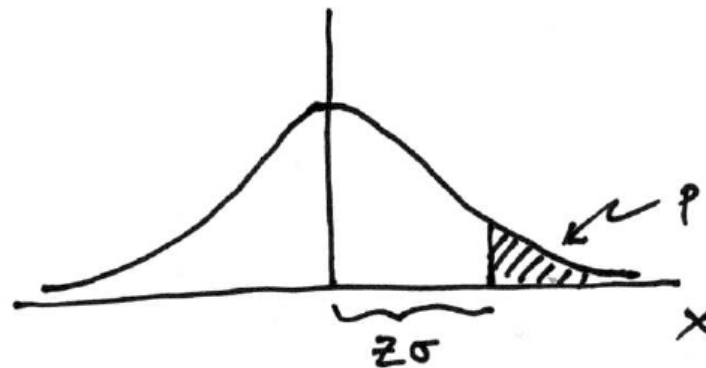
$$P(p_0 \leq \alpha | H_0) = \alpha$$

We can define the critical region of a test of H_0 with size α as the set of data space where $p_0 \leq \alpha$.

Formally the p -value relates only to H_0 , but the resulting test will have a given power with respect to a given alternative H_1 .

Significance from p -value

Often define significance Z as the number of standard deviations that a Gaussian variable would fluctuate in one direction to give the same p -value.



$$p = \int_Z^{\infty} \frac{1}{\sqrt{2\pi}} e^{-x^2/2} dx = 1 - \Phi(Z) \quad \mathbf{1 - TMath::Freq}$$

$$Z = \Phi^{-1}(1 - p) \quad \mathbf{TMath::NormQuantile}$$

E.g. $Z = 5$ (a “5 sigma effect”) corresponds to $p = 2.9 \times 10^{-7}$.

The Poisson counting experiment

Suppose we do a counting experiment and observe n events.

Events could be from *signal* process or from *background* – we only count the total number.

Poisson model:

$$P(n|s, b) = \frac{(s + b)^n}{n!} e^{-(s+b)}$$

s = mean (i.e., expected) # of signal events

b = mean # of background events

Goal is to make inference about s , e.g.,

test $s = 0$ (rejecting $H_0 \approx$ “discovery of signal process”)

test all non-zero s (values not rejected = confidence interval)

In both cases need to ask what is relevant alternative hypothesis.

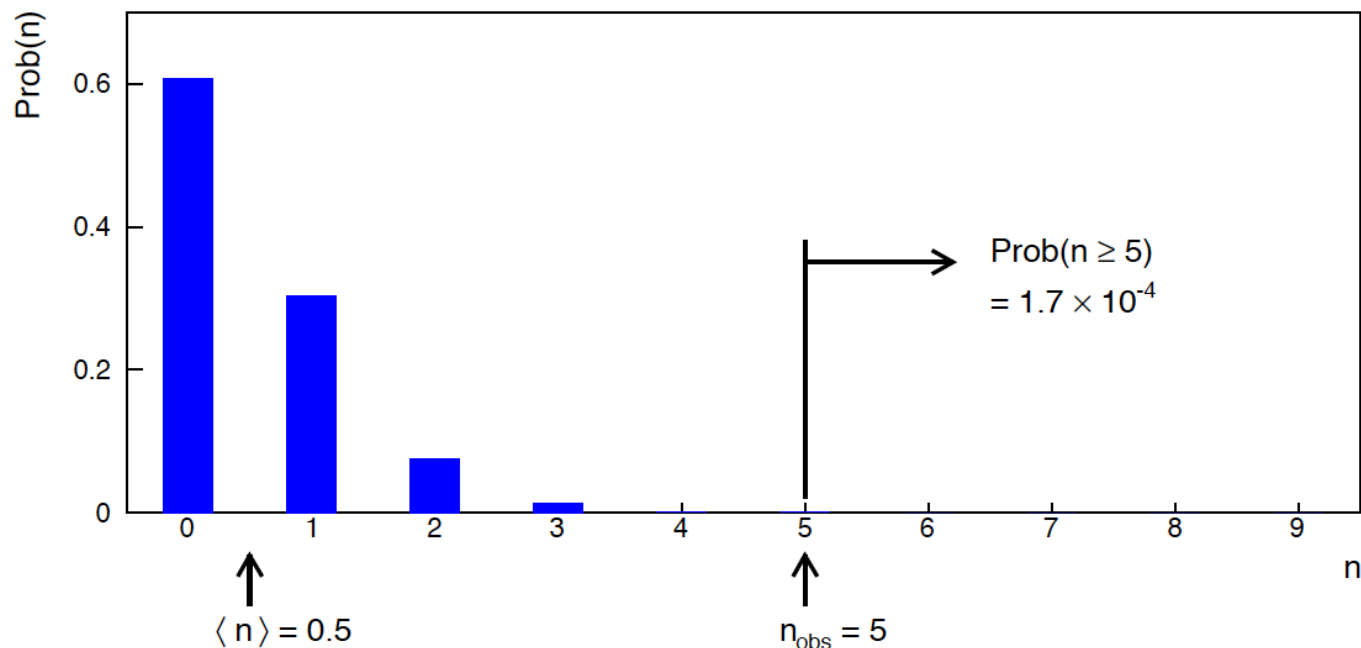
Poisson counting experiment: discovery p -value

Suppose $b = 0.5$ (known), and we observe $n_{\text{obs}} = 5$.

Should we claim evidence for a new discovery?

Give p -value for hypothesis $s = 0$:

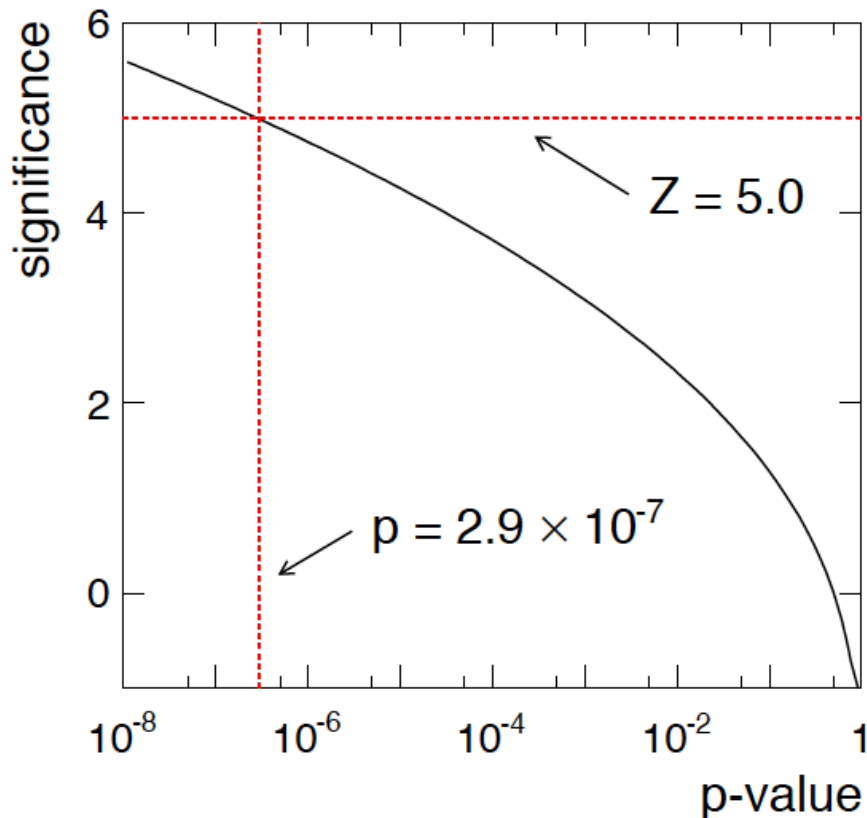
$$\begin{aligned} p\text{-value} &= P(n \geq 5; b = 0.5, s = 0) \\ &= 1.7 \times 10^{-4} \neq P(s = 0)! \end{aligned}$$



Poisson counting experiment: discovery significance

Equivalent significance for $p = 1.7 \times 10^{-4}$: $Z = \Phi^{-1}(1 - p) = 3.6$

Often claim discovery if $Z > 5$ ($p < 2.9 \times 10^{-7}$, i.e., a “5-sigma effect”)



In fact this tradition should be revisited: p -value intended to quantify probability of a signal-like fluctuation assuming background only; not intended to cover, e.g., hidden systematics, plausibility signal model, compatibility of data with signal, “look-elsewhere effect” (~multiple testing), etc.

Confidence intervals by inverting a test

Confidence intervals for a parameter θ can be found by defining a **test** of the hypothesized value θ (do this for all θ):

Specify values of the data that are ‘disfavoured’ by θ (critical region) such that $P(\text{data in critical region}) \leq \alpha$ for a prespecified α , e.g., 0.05 or 0.1.

If data observed in the critical region, reject the value θ .

Now **invert** the test to define a **confidence interval** as:

set of θ values that would **not** be rejected in a test of size α (confidence level is $1 - \alpha$).

The interval will cover the true value of θ with probability $\geq 1 - \alpha$.

Equivalently, the parameter values in the confidence interval have p -values of at least α .

To find edge of interval (the “limit”), set $p_\theta = \alpha$ and solve for θ .

Frequentist upper limit on Poisson parameter

Consider again the case of observing $n \sim \text{Poisson}(s + b)$.

Suppose $b = 4.5$, $n_{\text{obs}} = 5$. Find upper limit on s at 95% CL.

When testing s values to find upper limit, relevant alternative is $s = 0$ (or lower s), so critical region at low n and p -value of hypothesized s is $P(n \leq n_{\text{obs}}; s, b)$.

Upper limit s_{up} at $\text{CL} = 1 - \alpha$ from setting $\alpha = p_s$ and solving for s :

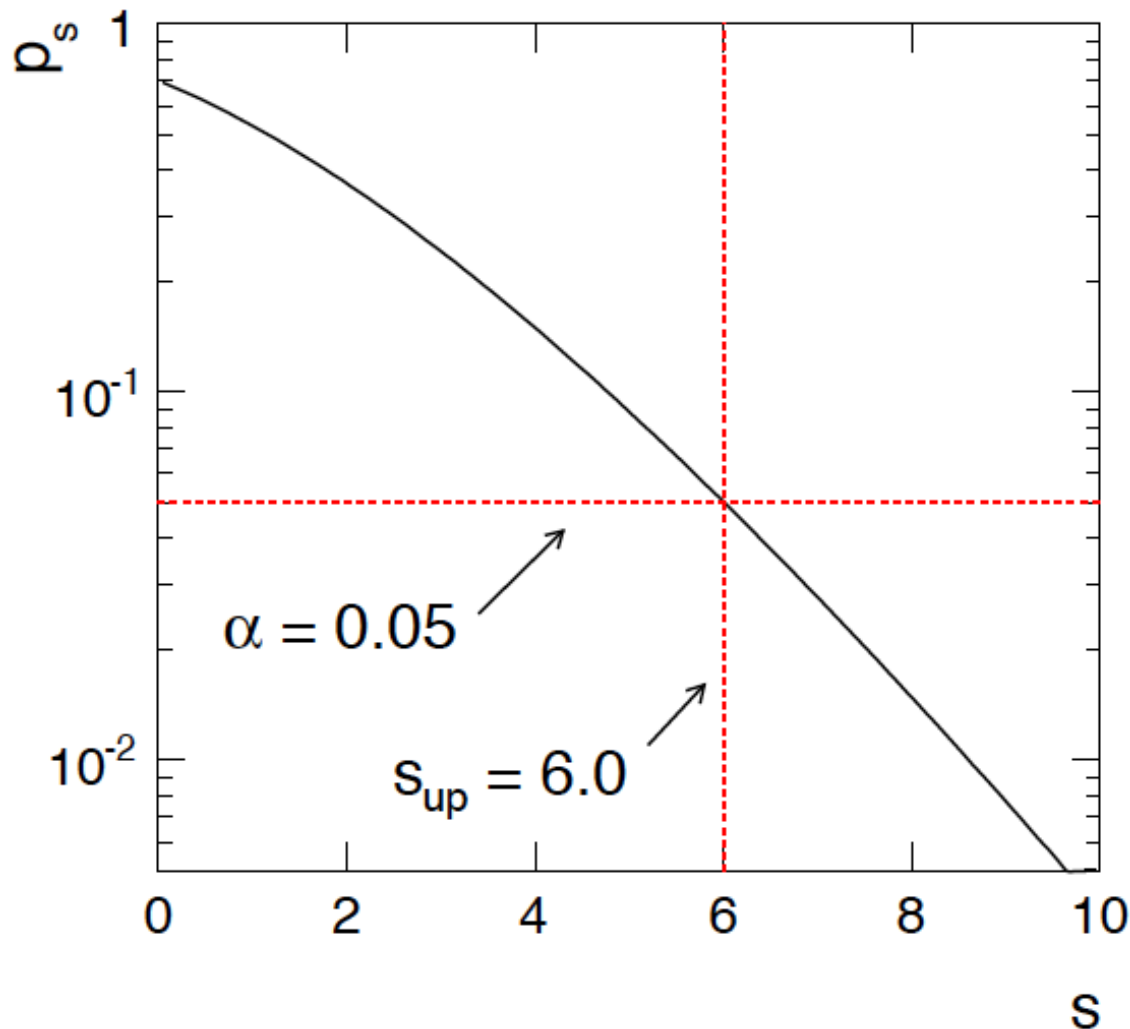
$$\alpha = P(n \leq n_{\text{obs}}; s_{\text{up}}, b) = \sum_{n=0}^{n_{\text{obs}}} \frac{(s_{\text{up}} + b)^n}{n!} e^{-(s_{\text{up}} + b)}$$

$$s_{\text{up}} = \frac{1}{2} F_{\chi^2}^{-1}(1 - \alpha; 2(n_{\text{obs}} + 1)) - b$$

$$= \frac{1}{2} F_{\chi^2}^{-1}(0.95; 2(5 + 1)) - 4.5 = 6.0$$

Frequentist upper limit on Poisson parameter

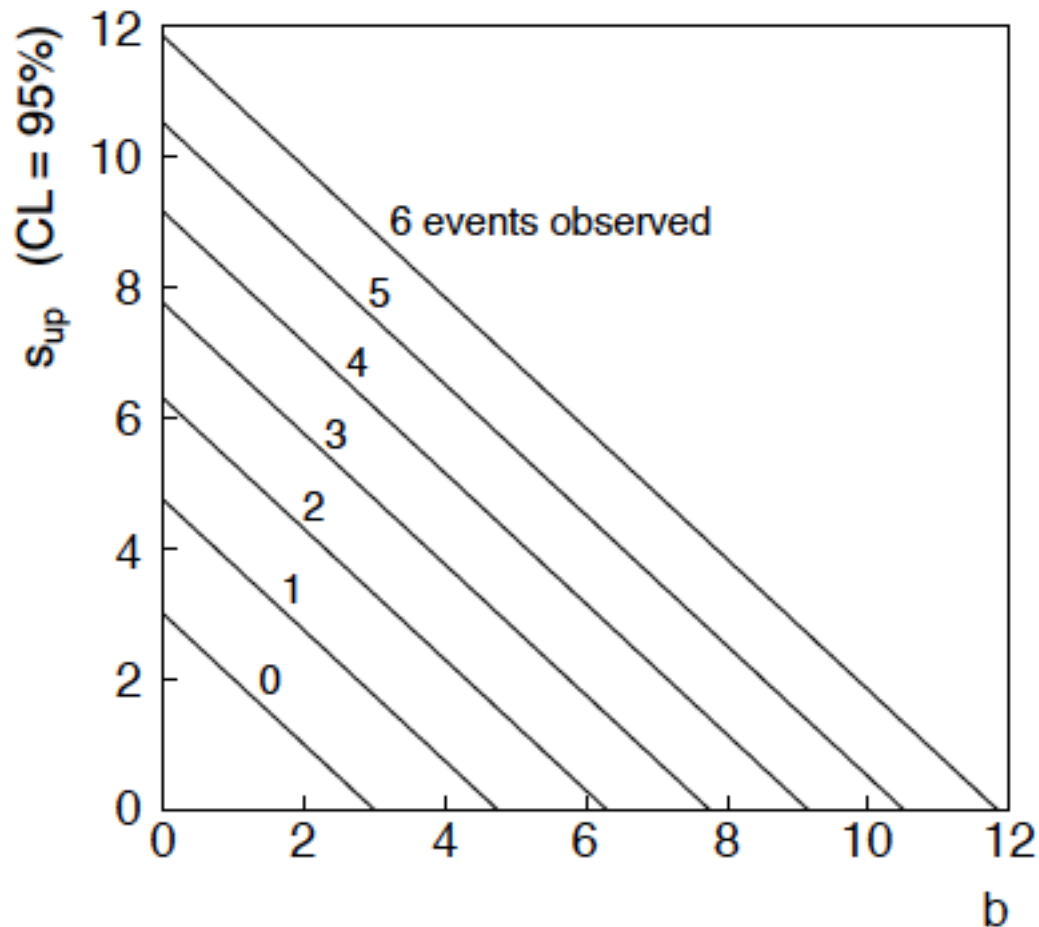
Upper limit s_{up} at $\text{CL} = 1 - \alpha$ found from $p_s = \alpha$.



$$n_{\text{obs}} = 5,$$
$$b = 4.5$$

$n \sim \text{Poisson}(s+b)$: frequentist upper limit on s

For low fluctuation of n formula can give negative result for s_{up} ; i.e. confidence interval is empty.



Limits near a physical boundary

Suppose e.g. $b = 2.5$ and we observe $n = 0$.

If we choose $CL = 0.9$, we find from the formula for s_{up}

$$s_{\text{up}} = -0.197 \quad (CL = 0.90)$$

Physicist:

We already knew $s \geq 0$ before we started; can't use negative upper limit to report result of expensive experiment!

Statistician:

The interval is designed to cover the true value only 90% of the time — this was clearly not one of those times.

Not uncommon dilemma when testing parameter values for which one has very little experimental sensitivity, e.g., very small s .

Expected limit for $s = 0$

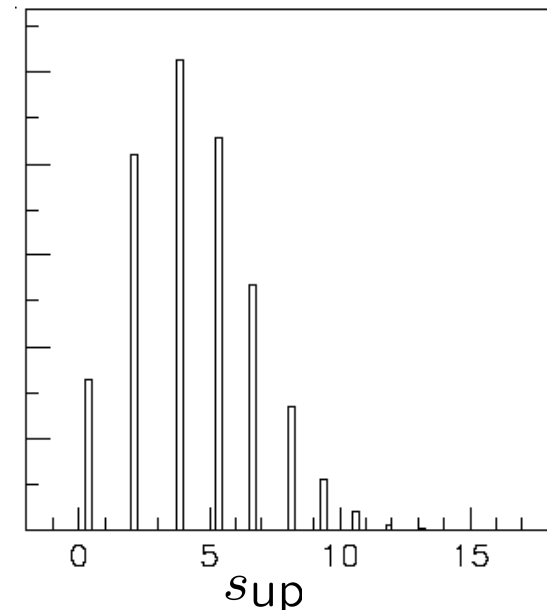
Physicist: I should have used $CL = 0.95$ — then $s_{\text{up}} = 0.496$

Even better: for $CL = 0.917923$ we get $s_{\text{up}} = 10^{-4}$!

Reality check: with $b = 2.5$, typical Poisson fluctuation in n is at least $\sqrt{2.5} = 1.6$. How can the limit be so low?

Look at the mean limit for the no-signal hypothesis ($s = 0$) (sensitivity).

Distribution of 95% CL limits with $b = 2.5$, $s = 0$.
Mean upper limit = 4.44



The Bayesian approach to limits

In Bayesian statistics need to start with ‘prior pdf’ $\pi(\theta)$, this reflects degree of belief about θ before doing the experiment.

Bayes’ theorem tells how our beliefs should be updated in light of the data x :

$$p(\theta|x) = \frac{L(x|\theta)\pi(\theta)}{\int L(x|\theta')\pi(\theta') d\theta'} \propto L(x|\theta)\pi(\theta)$$

Integrate posterior pdf $p(\theta|x)$ to give interval with any desired probability content.

For e.g. $n \sim \text{Poisson}(s+b)$, 95% CL upper limit on s from

$$0.95 = \int_{-\infty}^{s_{\text{sup}}} p(s|n) ds$$

Bayesian prior for Poisson parameter

Include knowledge that $s \geq 0$ by setting prior $\pi(s) = 0$ for $s < 0$.

Could try to reflect ‘prior ignorance’ with e.g.

$$\pi(s) = \begin{cases} 1 & s \geq 0 \\ 0 & \text{otherwise} \end{cases}$$

Not normalized but this is OK as long as $L(s)$ dies off for large s .

Not invariant under change of parameter — if we had used instead a flat prior for, say, the mass of the Higgs boson, this would imply a non-flat prior for the expected number of Higgs events.

Doesn’t really reflect a reasonable degree of belief, but often used as a point of reference;

or viewed as a recipe for producing an interval whose frequentist properties can be studied (coverage will depend on true s).

Bayesian interval with flat prior for s

Solve to find limit s_{up} :

$$s_{\text{up}} = \frac{1}{2} F_{\chi^2}^{-1} [p, 2(n+1)] - b$$

where

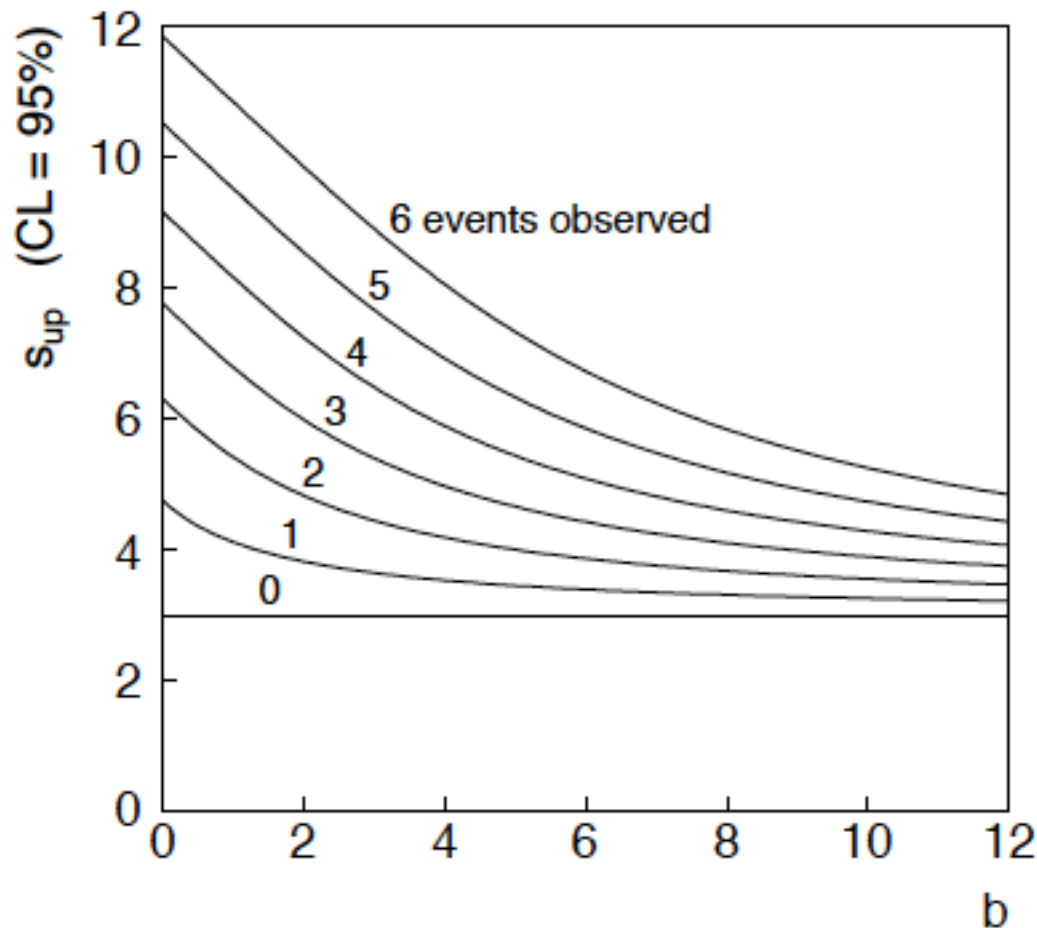
$$p = 1 - \alpha \left(1 - F_{\chi^2} [2b, 2(n+1)] \right)$$

For special case $b = 0$, Bayesian upper limit with flat prior numerically same as one-sided frequentist case ('coincidence').

Bayesian interval with flat prior for s

For $b > 0$ Bayesian limit is everywhere greater than the (one sided) frequentist upper limit.

Never goes negative. Doesn't depend on b if $n = 0$.



Priors from formal rules

Because of difficulties in encoding a vague degree of belief in a prior, one often attempts to derive the prior from formal rules, e.g., to satisfy certain invariance principles or to provide maximum information gain for a certain set of measurements.

Often called “objective priors”

Form basis of Objective Bayesian Statistics

The priors do not reflect a degree of belief (but might represent possible extreme cases).

In Objective Bayesian analysis, can use the intervals in a frequentist way, i.e., regard Bayes’ theorem as a recipe to produce an interval with certain coverage properties.

Priors from formal rules (cont.)

For a review of priors obtained by formal rules see, e.g.,

Robert E. Kass and Larry Wasserman, *The Selection of Prior Distributions by Formal Rules*, J. Am. Stat. Assoc., Vol. 91, No. 435, pp. 1343-1370 (1996).

Formal priors have not been widely used in HEP, but there is recent interest in this direction, especially the reference priors of Bernardo and Berger; see e.g.

L. Demortier, S. Jain and H. Prosper, *Reference priors for high energy physics*, Phys. Rev. D 82 (2010) 034002, arXiv:1002.1111.

D. Casadei, *Reference analysis of the signal + background model in counting experiments*, JINST 7 (2012) 01012; arXiv:1108.4270.

Approximate confidence intervals/regions from the likelihood function

Suppose we test parameter value(s) $\theta = (\theta_1, \dots, \theta_n)$ using the ratio


$$\lambda(\theta) = \frac{L(\theta)}{L(\hat{\theta})} \quad 0 \leq \lambda(\theta) \leq 1$$

Lower $\lambda(\theta)$ means worse agreement between data and hypothesized θ . Equivalently, usually define

$$t_\theta = -2 \ln \lambda(\theta)$$

so higher t_θ means worse agreement between θ and the data.

p -value of θ therefore $p_\theta = \int_{t_{\theta, \text{obs}}}^{\infty} f(t_\theta | \theta) dt_\theta$

 need pdf

Confidence region from Wilks' theorem

Wilks' theorem says (in large-sample limit and providing certain conditions hold...)

$$f(t_{\theta}|\theta) \sim \chi_n^2$$

chi-square dist. with # d.o.f. =
of components in $\theta = (\theta_1, \dots, \theta_n)$.

Assuming this holds, the p -value is

$$p_{\theta} = 1 - F_{\chi_n^2}(t_{\theta}) \quad \text{where} \quad F_{\chi_n^2}(t_{\theta}) \equiv \int_0^{t_{\theta}} f_{\chi_n^2}(t'_{\theta}) t'_{\theta}$$

To find boundary of confidence region set $p_{\theta} = \alpha$ and solve for t_{θ} :

$$t_{\theta} = -2 \ln \frac{L(\theta)}{L(\hat{\theta})} = F_{\chi_n^2}^{-1}(1 - \alpha)$$

Confidence region from Wilks' theorem (cont.)

i.e., boundary of confidence region in θ space is where

$$\ln L(\theta) = \ln L(\hat{\theta}) - \frac{1}{2} F_{\chi_n^2}^{-1}(1 - \alpha)$$

For example, for $1 - \alpha = 68.3\%$ and $n = 1$ parameter,

$$F_{\chi_1^2}^{-1}(0.683) = 1$$

and so the 68.3% confidence level interval is determined by

$$\ln L(\theta) = \ln L(\hat{\theta}) - \frac{1}{2}$$

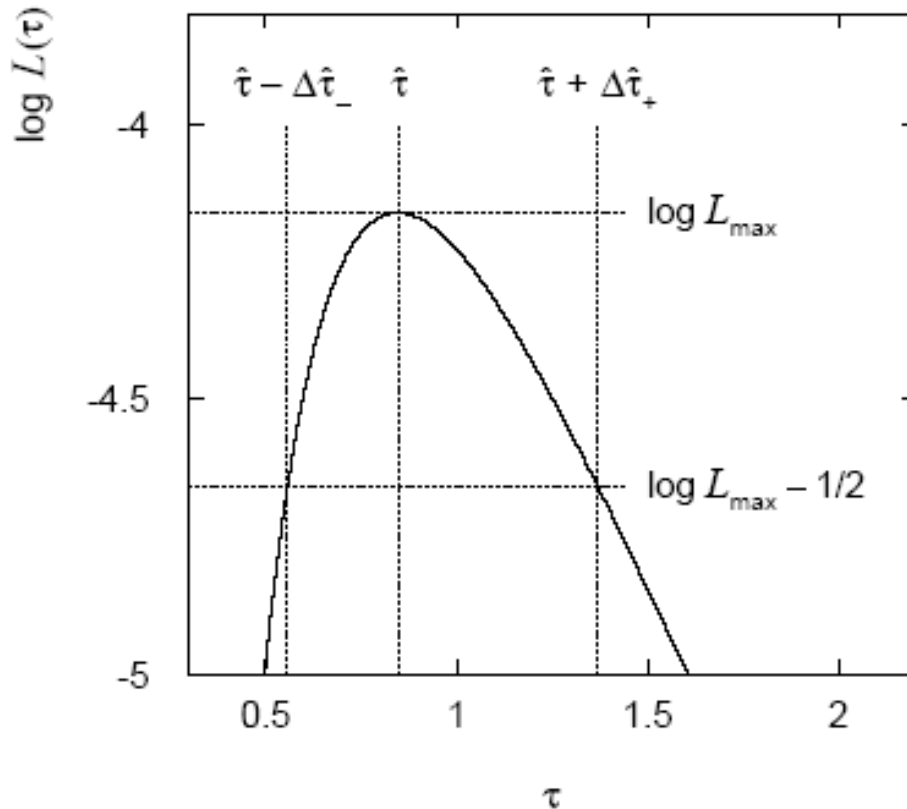
Same as recipe for finding the estimator's standard deviation, i.e.,

$[\hat{\theta} - \sigma_{\hat{\theta}}, \hat{\theta} + \sigma_{\hat{\theta}}]$ is a 68.3% CL confidence interval.

Example of interval from $\ln L$

For $n = 1$ parameter, $CL = 0.683$, $Q_\alpha = 1$.

Exponential example, now with only 5 events:



Parameter estimate and
approximate 68.3% CL
confidence interval:

$$\hat{\tau} = 0.85^{+0.52}_{-0.30}$$

Multiparameter case

For increasing number of parameters, $CL = 1 - \alpha$ decreases for confidence region determined by a given

$$Q_\alpha = F_{\chi_n^2}^{-1}(1 - \alpha)$$

Q_α	$1 - \alpha$				
	$n = 1$	$n = 2$	$n = 3$	$n = 4$	$n = 5$
1.0	0.683	0.393	0.199	0.090	0.037
2.0	0.843	0.632	0.428	0.264	0.151
4.0	0.954	0.865	0.739	0.594	0.451
9.0	0.997	0.989	0.971	0.939	0.891

Multiparameter case (cont.)

Equivalently, Q_α increases with n for a given $\text{CL} = 1 - \alpha$.

$1 - \alpha$	\bar{Q}_α				
	$n = 1$	$n = 2$	$n = 3$	$n = 4$	$n = 5$
0.683	1.00	2.30	3.53	4.72	5.89
0.90	2.71	4.61	6.25	7.78	9.24
0.95	3.84	5.99	7.82	9.49	11.1
0.99	6.63	9.21	11.3	13.3	15.1

Prototype search analysis

Search for signal in a region of phase space; result is histogram of some variable x giving numbers:

$$\mathbf{n} = (n_1, \dots, n_N)$$

Assume the n_i are Poisson distributed with expectation values

$$E[n_i] = \mu s_i + b_i$$

strength parameter

where

$$s_i = s_{\text{tot}} \int_{\text{bin } i} f_s(x; \boldsymbol{\theta}_s) dx, \quad b_i = b_{\text{tot}} \int_{\text{bin } i} f_b(x; \boldsymbol{\theta}_b) dx.$$

signal

background

Prototype analysis (II)

Often also have a subsidiary measurement that constrains some of the background and/or shape parameters:

$$\mathbf{m} = (m_1, \dots, m_M)$$

Assume the m_i are Poisson distributed with expectation values

$$E[m_i] = u_i(\boldsymbol{\theta})$$

nuisance parameters ($\boldsymbol{\theta}_s, \boldsymbol{\theta}_b, b_{\text{tot}}$)

Likelihood function is

$$L(\mu, \boldsymbol{\theta}) = \prod_{j=1}^N \frac{(\mu s_j + b_j)^{n_j}}{n_j!} e^{-(\mu s_j + b_j)} \prod_{k=1}^M \frac{u_k^{m_k}}{m_k!} e^{-u_k}$$

The profile likelihood ratio

Base significance test on the profile likelihood ratio:

$$\lambda(\mu) = \frac{L(\mu, \hat{\boldsymbol{\theta}})}{L(\hat{\mu}, \hat{\boldsymbol{\theta}})}$$

maximizes L for specified μ

maximize L

The likelihood ratio of point hypotheses gives optimum test (Neyman-Pearson lemma). In practice the profile LR is near-optimal.

Important advantage of profile LR is that its distribution becomes **independent of nuisance parameters** in large sample limit.

Test statistic for discovery

Try to reject background-only ($\mu=0$) hypothesis using

$$q_0 = \begin{cases} -2 \ln \lambda(0) & \hat{\mu} \geq 0 \\ 0 & \hat{\mu} < 0 \end{cases}$$

i.e. here only regard upward fluctuation of data as evidence against the background-only hypothesis.

Note that even though here physically $\mu \geq 0$, we allow $\hat{\mu}$ to be negative. In large sample limit its distribution becomes Gaussian, and this will allow us to write down simple expressions for distributions of our test statistics.

Distribution of q_0 in large-sample limit

Assuming approximations valid in the large sample (asymptotic) limit, we can write down the full distribution of q_0 as

$$f(q_0|\mu') = \left(1 - \Phi\left(\frac{\mu'}{\sigma}\right)\right) \delta(q_0) + \frac{1}{2} \frac{1}{\sqrt{2\pi}} \frac{1}{\sqrt{q_0}} \exp\left[-\frac{1}{2} \left(\sqrt{q_0} - \frac{\mu'}{\sigma}\right)^2\right]$$

The special case $\mu' = 0$ is a “half chi-square” distribution:

$$f(q_0|0) = \frac{1}{2} \delta(q_0) + \frac{1}{2} \frac{1}{\sqrt{2\pi}} \frac{1}{\sqrt{q_0}} e^{-q_0/2}$$

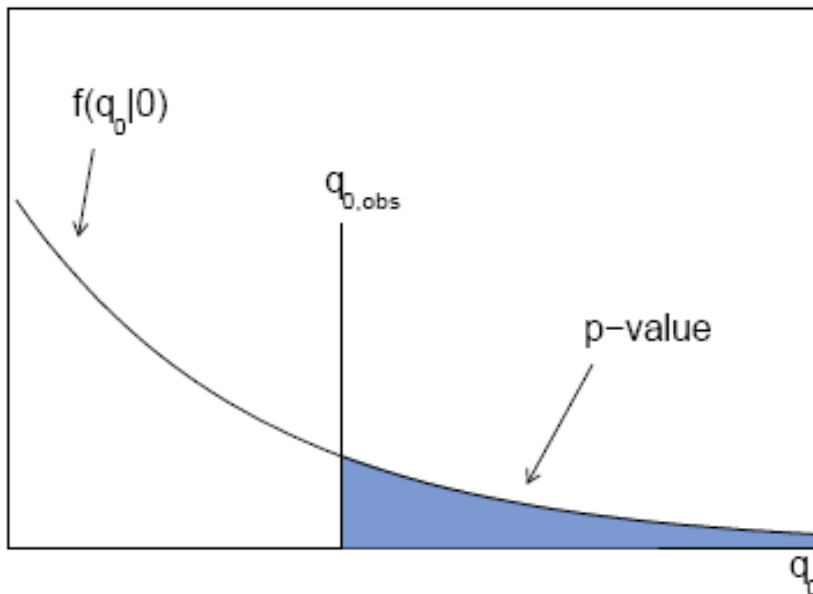
In large sample limit, $f(q_0|0)$ independent of nuisance parameters; $f(q_0|\mu')$ depends on nuisance parameters through σ .

p -value for discovery

Large q_0 means increasing incompatibility between the data and hypothesis, therefore p -value for an observed $q_{0,\text{obs}}$ is

$$p_0 = \int_{q_{0,\text{obs}}}^{\infty} f(q_0|0) dq_0$$

use e.g. asymptotic formula



From p -value get equivalent significance,

$$Z = \Phi^{-1}(1 - p)$$

Cumulative distribution of q_0 , significance

From the pdf, the cumulative distribution of q_0 is found to be

$$F(q_0|\mu') = \Phi\left(\sqrt{q_0} - \frac{\mu'}{\sigma}\right)$$

The special case $\mu' = 0$ is

$$F(q_0|0) = \Phi(\sqrt{q_0})$$

The p -value of the $\mu = 0$ hypothesis is

$$p_0 = 1 - F(q_0|0)$$

Therefore the discovery significance Z is simply

$$Z = \Phi^{-1}(1 - p_0) = \sqrt{q_0}$$

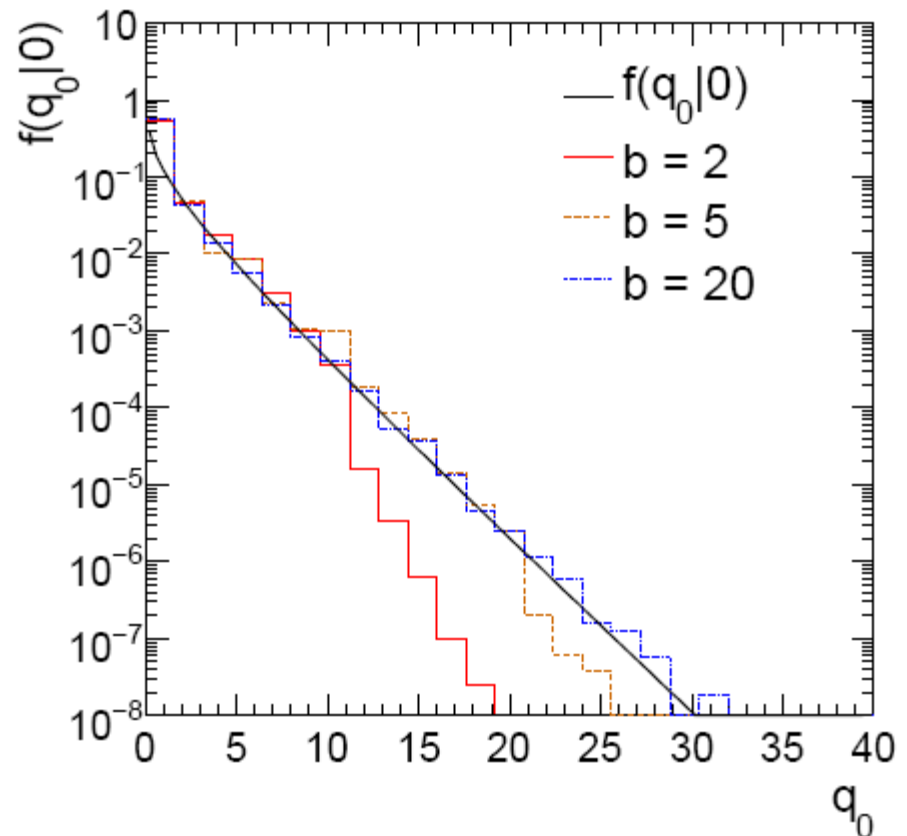
Monte Carlo test of asymptotic formula

$$n \sim \text{Poisson}(\mu s + b)$$

$$m \sim \text{Poisson}(\tau b)$$

Here take $\tau = 1$.

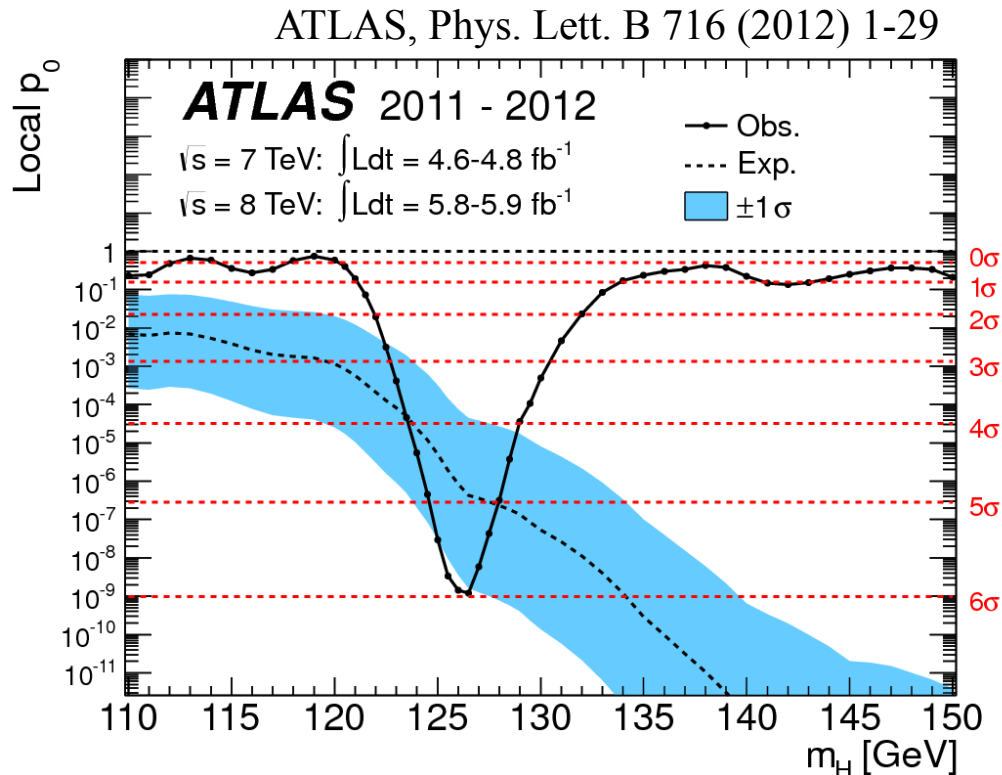
Asymptotic formula is
good approximation to 5σ
level ($q_0 = 25$) already for
 $b \sim 20$.



Example of discovery: the p_0 plot

The “local” p_0 means the p -value of the background-only hypothesis obtained from the test of $\mu = 0$ at each individual m_H , without any correction for the Look-Elsewhere Effect.

The “Expected” (dashed) curve gives the median p_0 under assumption of the SM Higgs ($\mu = 1$) at each m_H .



The blue band gives the width of the distribution ($\pm 1\sigma$) of significances under assumption of the SM Higgs.

Return to interval estimation

Suppose a model contains a parameter μ ; we want to know which values are consistent with the data and which are disfavoured.

Carry out a test of size α for all values of μ .

The values that are not rejected constitute a *confidence interval* for μ at confidence level $CL = 1 - \alpha$.

The probability that the true value of μ will be rejected is not greater than α , so by construction the confidence interval will contain the true value of μ with probability $\geq 1 - \alpha$.

The interval depends on the choice of the test (critical region).

If the test is formulated in terms of a p -value, p_μ , then the confidence interval represents those values of μ for which $p_\mu > \alpha$.

To find the end points of the interval, set $p_\mu = \alpha$ and solve for μ .

Test statistic for upper limits

cf. Cowan, Cranmer, Gross, Vitells, arXiv:1007.1727, EPJC 71 (2011) 1554.

For purposes of setting an upper limit on μ one can use

$$q_\mu = \begin{cases} -2 \ln \lambda(\mu) & \hat{\mu} \leq \mu \\ 0 & \hat{\mu} > \mu \end{cases} \quad \text{where} \quad \lambda(\mu) = \frac{L(\mu, \hat{\boldsymbol{\theta}})}{L(\hat{\mu}, \hat{\boldsymbol{\theta}})}$$

I.e. when setting an upper limit, an upwards fluctuation of the data is not taken to mean incompatibility with the hypothesized μ :

From observed q_μ find p -value:
$$p_\mu = \int_{q_{\mu, \text{obs}}}^{\infty} f(q_\mu | \mu) dq_\mu$$

Large sample approximation:

$$p_\mu = 1 - \Phi(\sqrt{q_\mu})$$

95% CL upper limit on μ is highest value for which p -value is not less than 0.05.

Monte Carlo test of asymptotic formulae

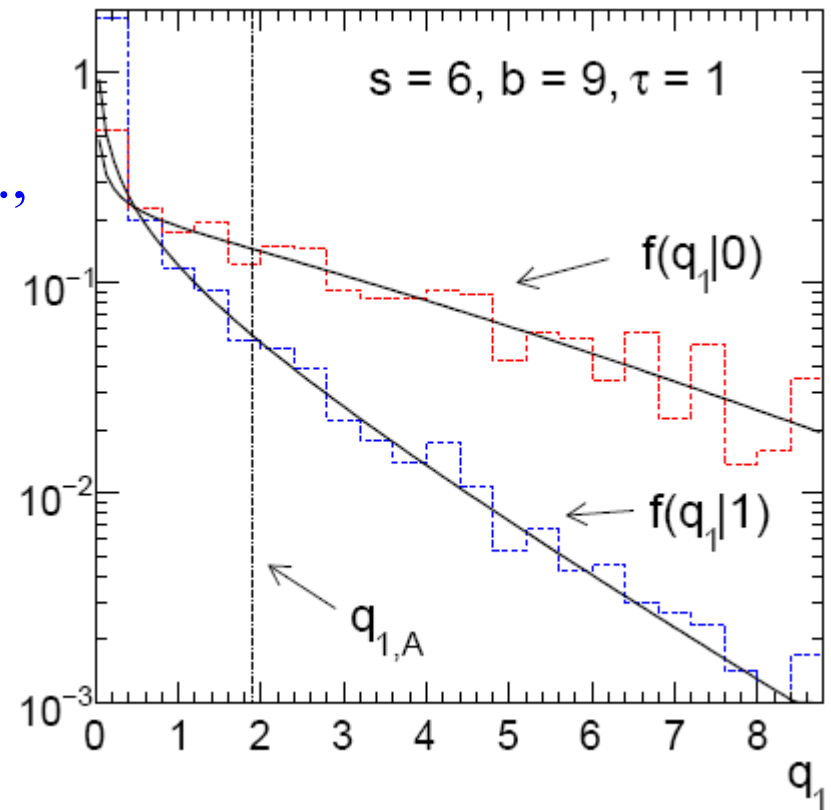
Consider again $n \sim \text{Poisson}(\mu s + b)$, $m \sim \text{Poisson}(\tau b)$
 Use q_μ to find p -value of hypothesized μ values.

E.g. $f(q_1|1)$ for p -value of $\mu=1$.

Typically interested in 95% CL, i.e.,
 p -value threshold = 0.05, i.e.,
 $q_1 = 2.69$ or $Z_1 = \sqrt{q_1} = 1.64$.

Median[$q_1|0$] gives “exclusion sensitivity”.

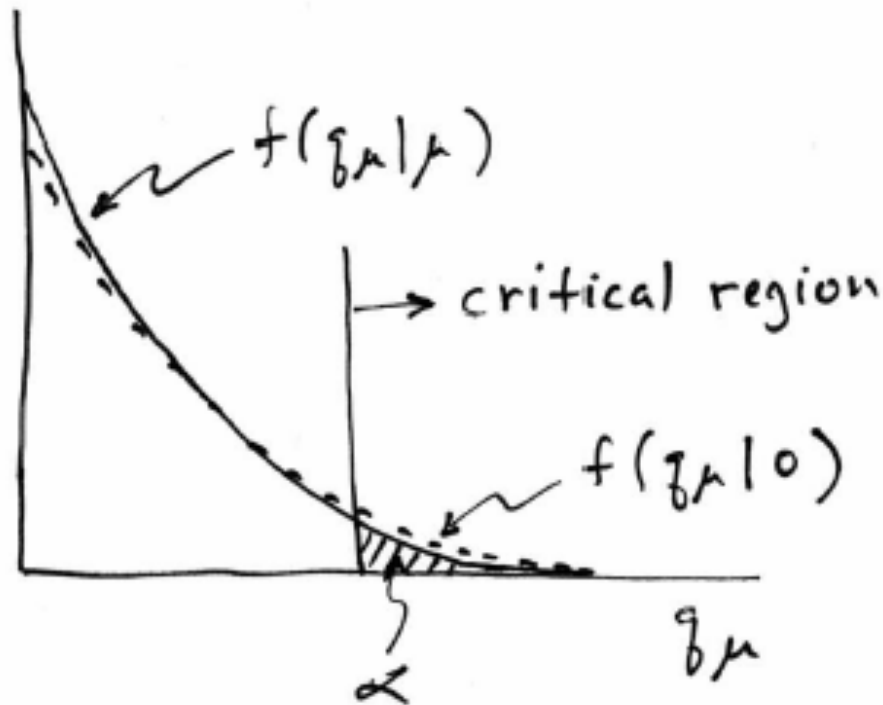
Here asymptotic formulae good
 for $s = 6$, $b = 9$.



Low sensitivity to μ

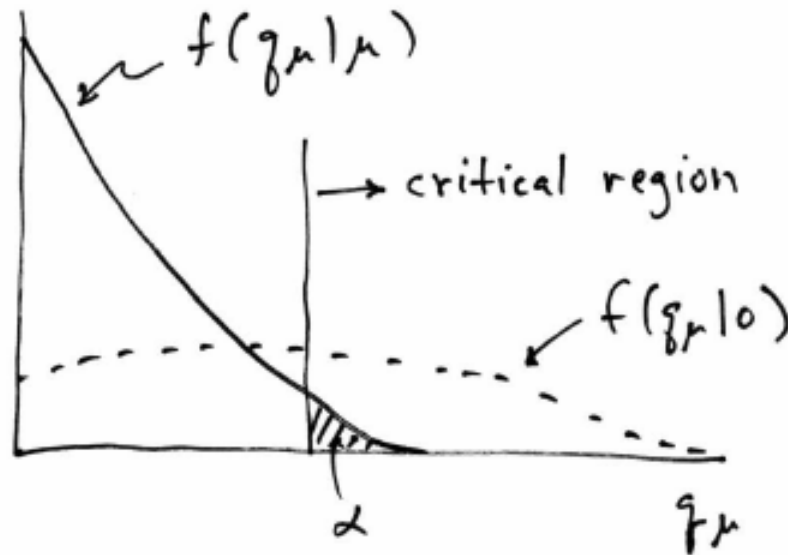
It can be that the effect of a given hypothesized μ is very small relative to the background-only ($\mu = 0$) prediction.

This means that the distributions $f(q_\mu|\mu)$ and $f(q_\mu|0)$ will be almost the same:



Having sufficient sensitivity

In contrast, having sensitivity to μ means that the distributions $f(q_\mu|\mu)$ and $f(q_\mu|0)$ are more separated:

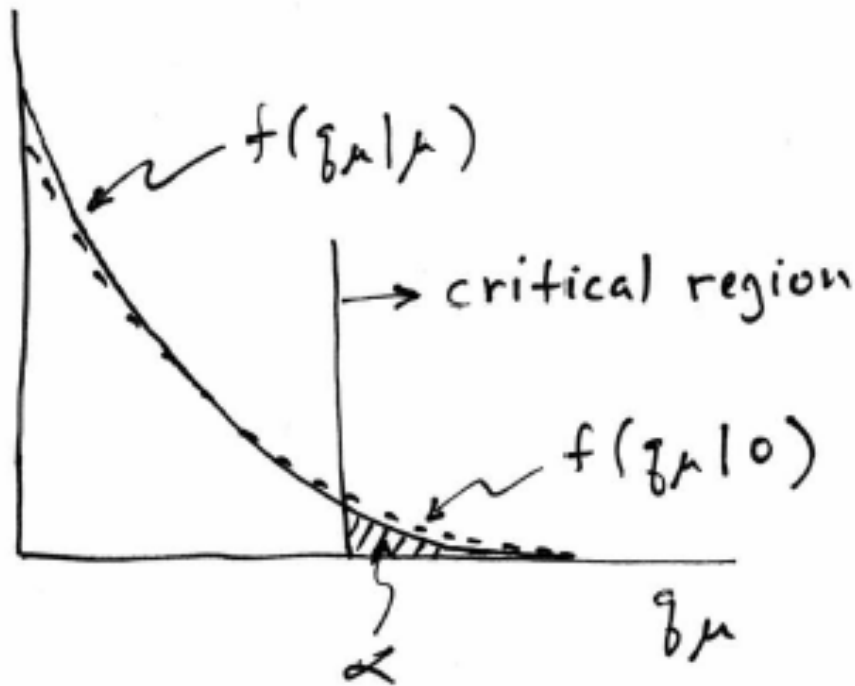


That is, the power (probability to reject μ if $\mu = 0$) is substantially higher than α . Use this power as a measure of the sensitivity.

Spurious exclusion

Consider again the case of low sensitivity. By construction the probability to reject μ if μ is true is α (e.g., 5%).

And the probability to reject μ if $\mu = 0$ (the power) is only slightly greater than α .



This means that with probability of around $\alpha = 5\%$ (slightly higher), one excludes hypotheses to which one has essentially no sensitivity (e.g., $m_H = 1000$ TeV).

“Spurious exclusion”

Ways of addressing spurious exclusion

The problem of excluding parameter values to which one has no sensitivity known for a long time; see e.g.,

Virgil L. Highland, *Estimation of Upper Limits from Experimental Data*, July 1986, Revised February 1987, Temple University Report C00-3539-38.

In the 1990s this was re-examined for the LEP Higgs search by Alex Read and others

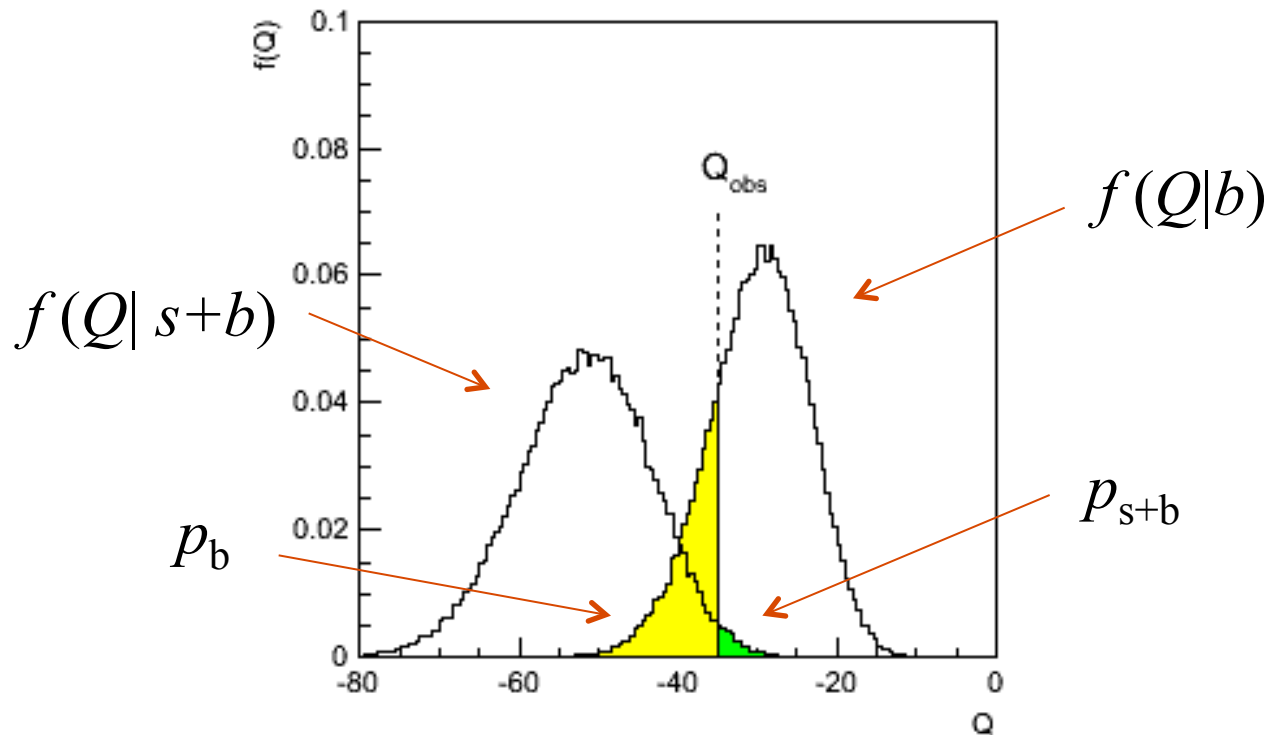
T. Junk, Nucl. Instrum. Methods Phys. Res., Sec. A **434**, 435 (1999); A.L. Read, J. Phys. G **28**, 2693 (2002).

and led to the “ CL_s ” procedure for upper limits.

Unified intervals also effectively reduce spurious exclusion by the particular choice of critical region.

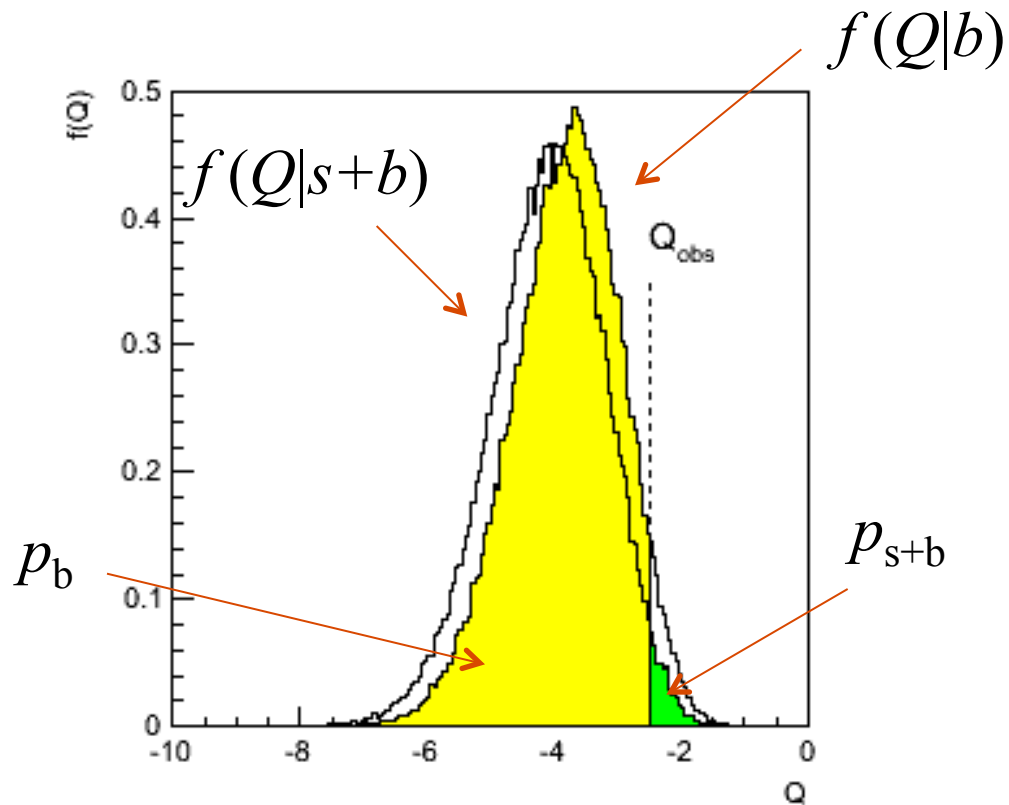
The CL_s procedure

In the usual formulation of CL_s , one tests both the $\mu = 0$ (b) and $\mu > 0$ ($\mu s+b$) hypotheses with the same statistic $Q = -2\ln L_{s+b}/L_b$:



The CL_s procedure (2)

As before, “low sensitivity” means the distributions of Q under b and $s+b$ are very close:



The CL_s procedure (3)

The CL_s solution (A. Read et al.) is to base the test not on the usual p -value (CL_{s+b}), but rather to divide this by CL_b (\sim one minus the p -value of the b -only hypothesis), i.e.,

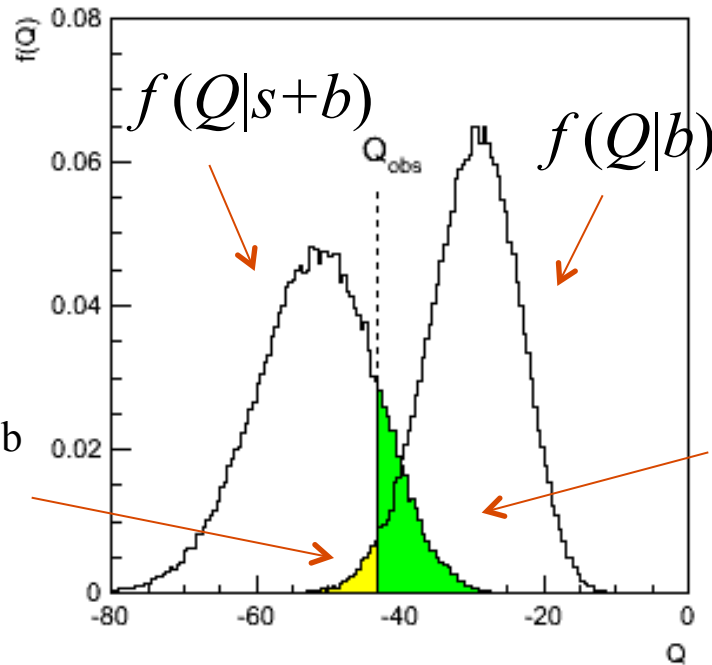
Define:

$$CL_s = \frac{CL_{s+b}}{CL_b} = \frac{p_{s+b}}{1 - p_b}$$

Reject $s+b$ hypothesis if:

$$CL_s \leq \alpha$$

$$1 - CL_b = p_b$$



$$CL_{s+b} = p_{s+b}$$

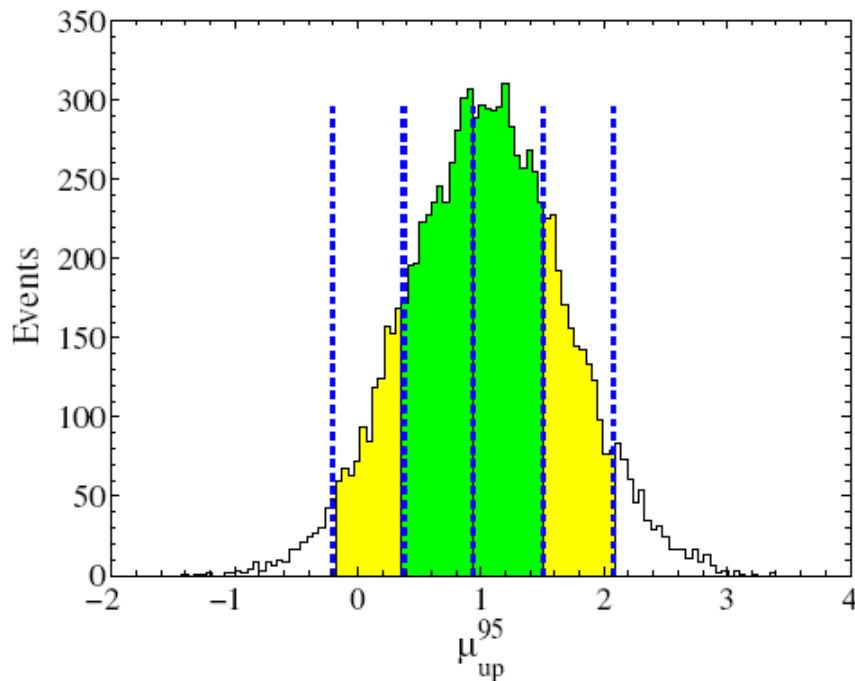
Increases “effective” p -value when the two distributions become close (prevents exclusion if sensitivity is low).

Setting upper limits on $\mu = \sigma/\sigma_{\text{SM}}$

Carry out the CLs procedure for the parameter $\mu = \sigma/\sigma_{\text{SM}}$, resulting in an upper limit μ_{up} .

In, e.g., a Higgs search, this is done for each value of m_{H} .

At a given value of m_{H} , we have an observed value of μ_{up} , and we can also find the distribution $f(\mu_{\text{up}}|0)$:



$\pm 1\sigma$ (green) and $\pm 2\sigma$ (yellow) bands from toy MC;

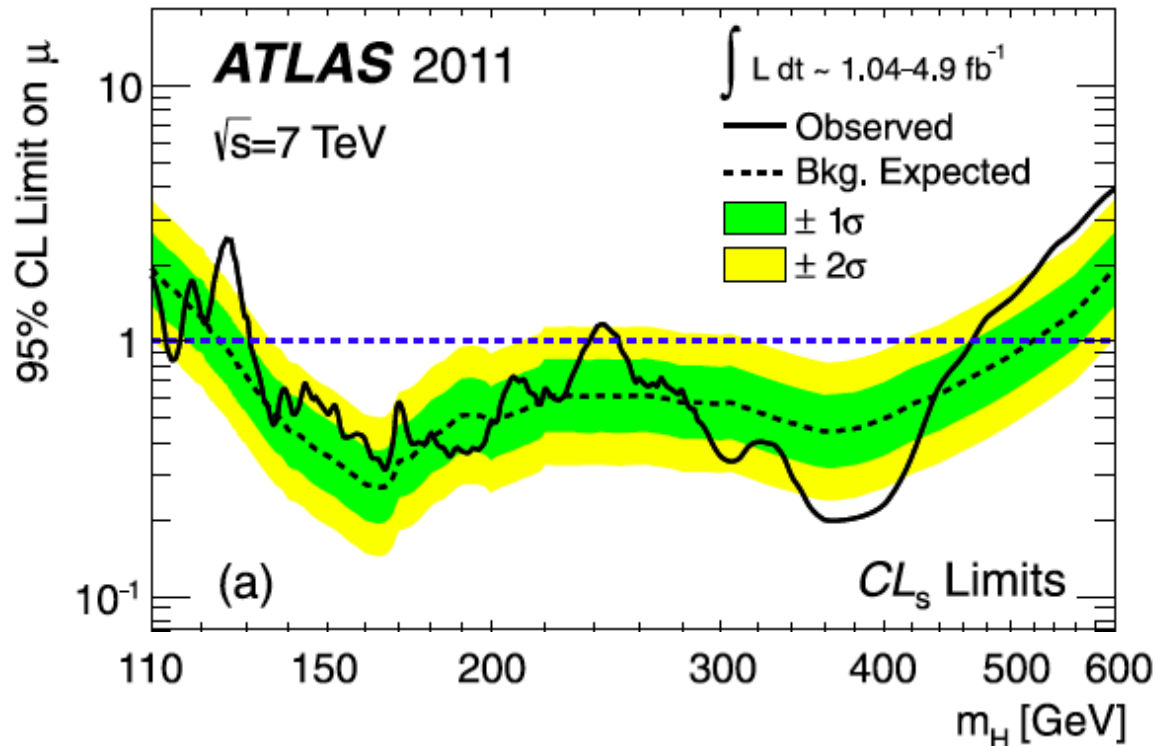
Vertical lines from asymptotic formulae.

How to read the green and yellow limit plots

For every value of m_H , find the CLs upper limit on μ .

Also for each m_H , determine the distribution of upper limits μ_{up} one would obtain under the hypothesis of $\mu = 0$.

The dashed curve is the median μ_{up} , and the green (yellow) bands give the $\pm 1\sigma$ (2σ) regions of this distribution.

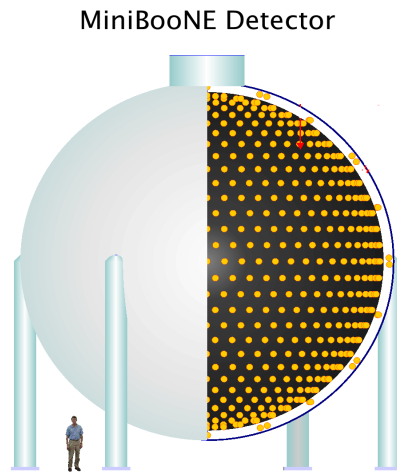


ATLAS, Phys. Lett.
B 710 (2012) 49-66

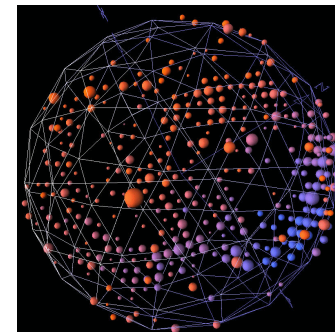
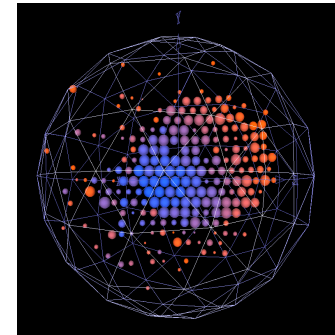
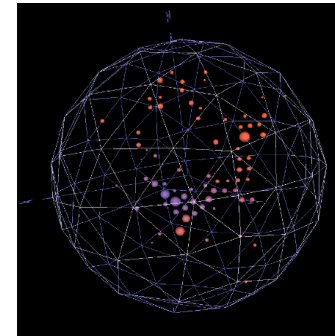
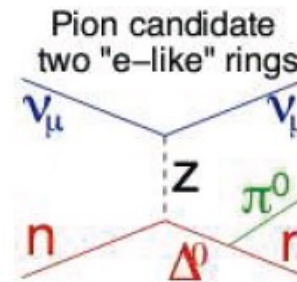
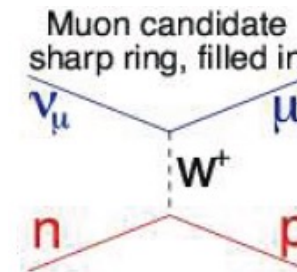
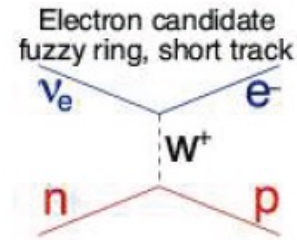
Extra slides

Particle i.d. in MiniBooNE

Detector is a 12-m diameter tank of mineral oil exposed to a beam of neutrinos and viewed by 1520 photomultiplier tubes:



Search for ν_μ to ν_e oscillations required particle i.d. using information from the PMTs.



H.J. Yang, MiniBooNE PID, DNP06

Decision trees

Out of all the input variables, find the one for which with a single cut gives best improvement in signal purity:

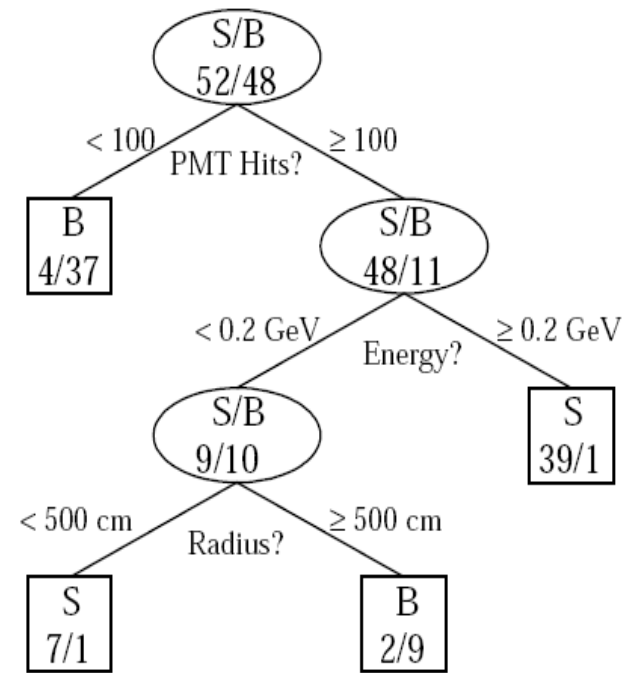
$$P = \frac{\sum_{\text{signal}} w_i}{\sum_{\text{signal}} w_i + \sum_{\text{background}} w_i}$$

where w_i is the weight of the i th event.

Resulting nodes classified as either signal/background.

Iterate until stop criterion reached based on e.g. purity or minimum number of events in a node.

The set of cuts defines the decision boundary.



Example by MiniBooNE experiment, B. Roe et al., NIM 543 (2005) 577

Finding the best single cut

The level of separation within a node can, e.g., be quantified by the *Gini coefficient*, calculated from the (s or b) purity as:

$$G = p(1 - p)$$

For a cut that splits a set of events a into subsets b and c , one can quantify the improvement in separation by the change in weighted Gini coefficients:

$$\Delta = W_a G_a - W_b G_b - W_c G_c \quad \text{where, e.g.,} \quad W_a = \sum_{i \in a} w_i$$

Choose e.g. the cut to the maximize Δ ; a variant of this scheme can use instead of Gini e.g. the misclassification rate:

$$\varepsilon = 1 - \max(p, 1 - p)$$

Decision tree classifier

The terminal nodes (**leaves**) are classified a signal or background depending on majority vote (or e.g. signal fraction greater than a specified threshold).

This classifies every point in input-variable space as either signal or background, a **decision tree classifier**, with discriminant function

$$f(\mathbf{x}) = 1 \text{ if } \mathbf{x} \text{ in signal region, } -1 \text{ otherwise}$$

Decision trees tend to be very sensitive to statistical fluctuations in the training sample.

Methods such as **boosting** can be used to stabilize the tree.

Boosting

Boosting is a general method of creating a set of classifiers which can be combined to achieve a new classifier that is more stable and has a smaller error than any individual one.

Often applied to decision trees but, can be applied to any classifier.

Suppose we have a training sample T consisting of N events with

$\mathbf{x}_1, \dots, \mathbf{x}_N$ event data vectors (each \mathbf{x} multivariate)

y_1, \dots, y_N true class labels, +1 for signal, -1 for background

w_1, \dots, w_N event weights

Now define a rule to create from this an ensemble of training samples T_1, T_2, \dots , derive a classifier from each and average them.

Trick is to create modifications in the training sample that give classifiers with smaller error rates than those of the preceding ones.

A successful example is **AdaBoost** (Freund and Schapire, 1997).

AdaBoost

First initialize the training sample T_1 using the original

$\mathbf{x}_1, \dots, \mathbf{x}_N$ event data vectors

$\mathbf{y}_1, \dots, \mathbf{y}_N$ true class labels (+1 or -1)

$\mathbf{w}_1^{(1)}, \dots, \mathbf{w}_N^{(1)}$ event weights

with the weights equal and normalized such that

$$\sum_{i=1}^N w_i^{(1)} = 1$$

Then train the classifier $f_1(\mathbf{x})$ (e.g., a decision tree) with a method that uses the event weights. Recall for an event at point \mathbf{x} ,

$f_1(\mathbf{x}) = +1$ for \mathbf{x} in signal region, -1 in background region

We will define an iterative procedure that gives a series of classifiers $f_1(\mathbf{x}), f_2(\mathbf{x}), \dots$

Error rate of the k th classifier

At the k th iteration the classifier $f_k(\mathbf{x})$ has an error rate

$$\varepsilon_k = \sum_{i=1}^N w_i^{(k)} I(y_i f_k(\mathbf{x}_i) \leq 0)$$

where $I(X) = 1$ if X is true and is zero otherwise.

Next assign a score to the k th classifier based on its error rate,

$$\alpha_k = \frac{1}{2} \ln \frac{1 - \varepsilon_k}{\varepsilon_k}$$

Updating the event weights

The classifier at each iterative step is found from an updated training sample, in which the weight of event i is modified from step k to step $k+1$ according to

$$w_i^{(k+1)} = w_i^{(k)} \frac{e^{-\alpha_k f_k(\mathbf{x}_i) y_i}}{Z_k}$$

Here Z_k is a normalization factor defined such that the sum of the weights over all events is equal to one.

That is, the weight for event i is increased in the $k+1$ training sample if it was classified incorrectly in step k .

Idea is that next time around the classifier should pay more attention to this event and try to get it right.

Defining the classifier

After K boosting iterations, the final classifier is defined as a weighted linear combination of the $f_k(\mathbf{x})$,

$$t(\mathbf{x}) = \sum_{k=1}^K \alpha_k f_k(\mathbf{x})$$

One can show that the error rate on the training data of the final classifier satisfies the bound

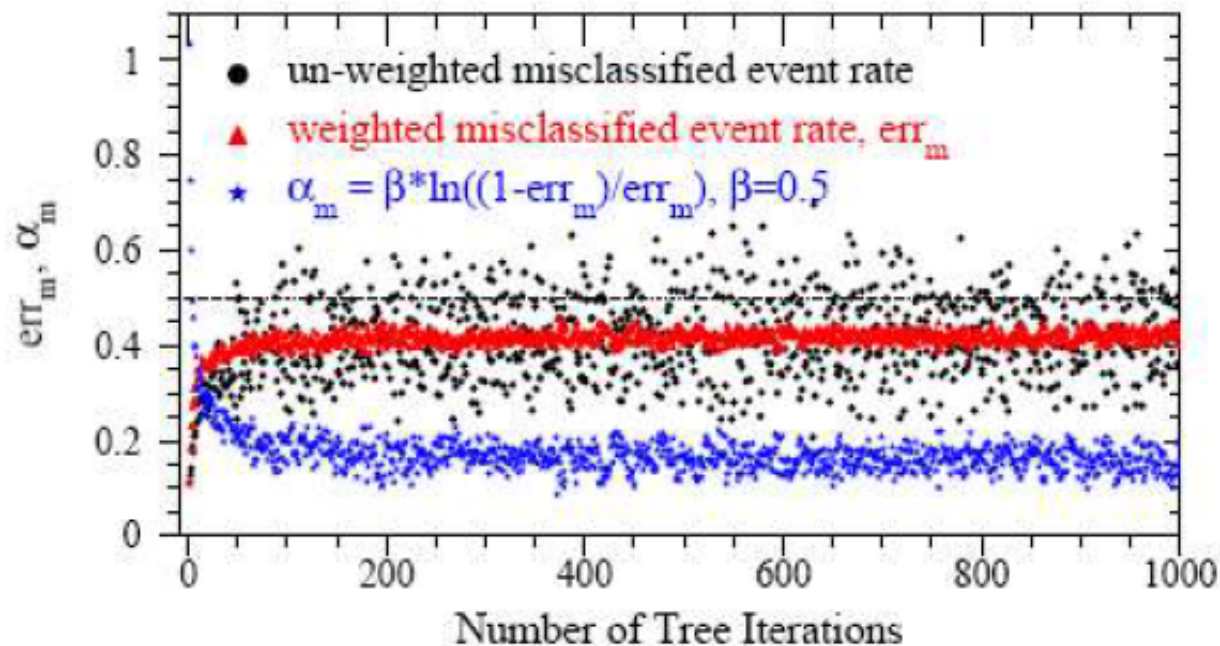
$$\varepsilon \leq \prod_{k=1}^K 2\sqrt{\varepsilon_k(1 - \varepsilon_k)}$$

i.e. as long as the $\varepsilon_k < \frac{1}{2}$ (better than random guessing), with enough boosting iterations every event in the training sample will be classified correctly.

BDT example from MiniBooNE

~200 input variables for each event (ν interaction producing e , μ or π).

Each individual tree is relatively weak, with a misclassification error rate $\sim 0.4 - 0.45$



B. Roe et al., NIM 543 (2005) 577

Monitoring overtraining

From MiniBooNE
example:

Performance stable
after a few hundred
trees.

Training MC Samples .VS. Testing MC Samples

