

PH4442 Advanced Particle Physics 2025/26

Lecture Week 5

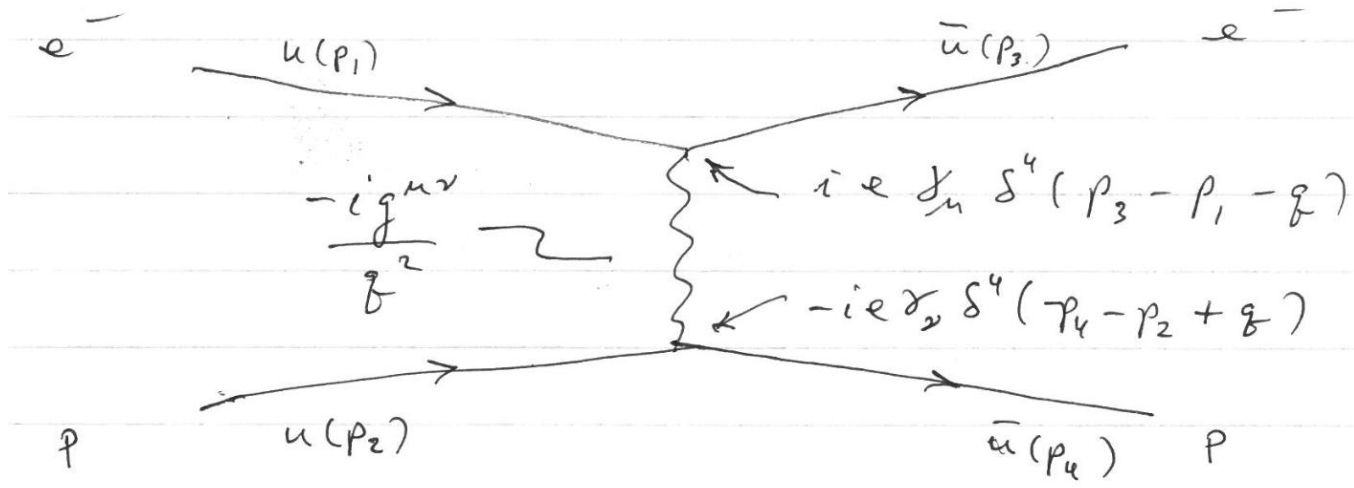


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- Summary of Feynman rules
- From amplitudes to observables
- $e^+e^- \rightarrow \mu^+\mu^-$
- The ALEPH detector at LEP

Recall S-matrix element for $ep \rightarrow ep$

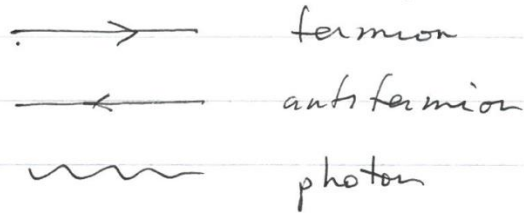
$$S_{fi}^{(1)} = (2\pi)^4 \delta^4(p_3 + p_4 - p_1 - p_2) [\bar{u}(p_3)(ie\gamma_\mu)u(p_1)] \frac{-i}{q^2} [\bar{u}(p_4)(-ie\gamma^\mu)u(p_2)]$$



Related to invariant amplitude \mathcal{M} by $S_{fi} = (2\pi)^4 \delta^4(P_f - P_i) \mathcal{M}$

$$\rightarrow \mathcal{M} = \frac{-ie^2}{(p_3 - p_1)^2} [\bar{u}(p_3)\gamma_\mu u(p_1)] [\bar{u}(p_4)\gamma^\mu u(p_2)]$$

Feynman rules



Incoming (outgoing) fermion $u(p)$ ($\bar{u}(p)$)

" " antifermion $v(p)$ ($\bar{v}(p)$)

" " photon $\epsilon_\mu(k)$ ($\epsilon_\mu^*(k)$)
 ↗ polarisation vector

virtual photon $\frac{-ig^{\mu\nu}}{q^2 + i\epsilon}$
 intermediate }
 virtual electron $\frac{i(\not{p} + m)}{p^2 - m^2 + i\epsilon}$ } For this course, won't need ϵ

fermion- γ vertex: $-i(\pm e) \underbrace{\gamma_\mu}_{\epsilon}$

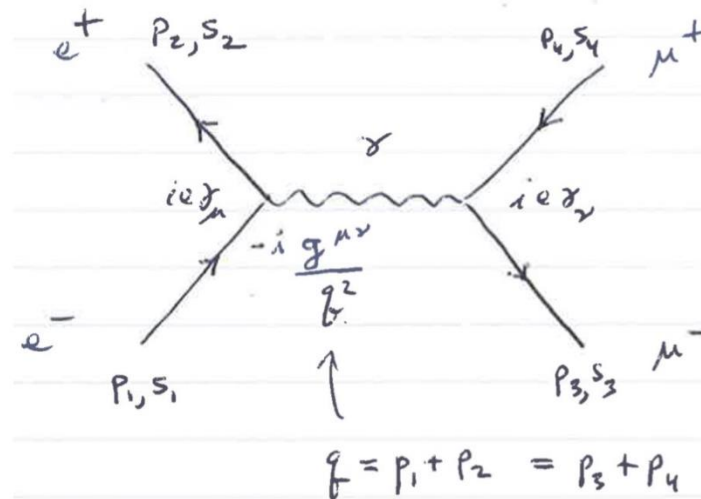
Feynman rules

- For a given reaction, write down all topologically distinct diagrams for that connect the initial and final state particles using the allowed vertices of the theory. Assign four-momenta to each particle so that it is conserved at each vertex.
- The conventions for lines were shown in Fig. 1.1. For fermions, the arrow points forward in time and for antifermions it points backwards. Photons are drawn with wavy lines.
- For each external fermion line, start at the tip of the arrow and write a Dirac spinor \bar{u} for a particle and \bar{v} for an antiparticle, evaluated with the particle's four-momentum.
- At a fermion-photon vertex, write down the factor $-iq\gamma_\mu$, where q is the charge of the particle, e.g., $ie\gamma_\mu$ for an electron.
- After the vertex factor, if the tail of the fermion line is in the initial or final state, write down a Dirac spinor u for a particle or v for antiparticle.
- For a photon in an initial or final state, include a factor of the polarisation vector ε^μ or $\varepsilon^{\mu*}$, respectively. Its Lorentz index is contracted with that of the gamma matrix at the vertex to which the photon is coupled.
- For an internal photon, include the propagator $-ig^{\mu\nu}/q^2$. The Lorentz indices are contracted with those of the gamma matrices of the vertices at each end of the photon.
- For an internal electron, include the propagator $i(\not{p} + m)/(p^2 - m^2)$.
- For a diagram that differs from another by exchange of an identical fermion, include a minus sign.

Amplitude for $e^+e^- \rightarrow \mu^+\mu^-$

In QED only one leading-order diagram (SM also has annihilation into Z).

Assign four-momenta, enforce conservation at every vertex.



Follow the fermion lines backwards, bar at head, no bar at tail, u for particle, v for antiparticle:

$$\mathcal{M} = [\bar{v}(p_2, s_2)(ie\gamma_\mu)u(p_1, s_1)] \frac{-ig^{\mu\nu}}{q^2} [\bar{u}(p_3, s_3)(ie\gamma_\nu)v(p_4, s_4)]$$

Contract Lorentz index ν , use $q^2 = E_{\text{cm}}^2 \equiv s$:

$$\mathcal{M} = i\frac{e^2}{s} [\bar{v}(p_2, s_2)\gamma_\mu u(p_1, s_1)] [\bar{u}(p_3, s_3)\gamma^\mu v(p_4, s_4)]$$

Trace theorems

$$\text{Tr}(\gamma^\mu) = \text{Tr}(\text{any product of odd number of } \gamma \text{ matrices}) = 0 ,$$

$$\text{Tr}(\gamma^\mu \gamma^\nu) = 4g^{\mu\nu} ,$$

$$\text{Tr}(\gamma^\mu \gamma^\nu \gamma^\rho \gamma^\sigma) = 4 [g^{\mu\nu} g^{\rho\sigma} + g^{\mu\sigma} g^{\nu\rho} - g^{\mu\rho} g^{\nu\sigma}] ,$$

$$\text{Tr}(\gamma^5) = 0 ,$$

$$\text{Tr}(\gamma^\mu \gamma^\nu \gamma^\rho \gamma^\sigma \gamma^5) = -4i\epsilon^{\mu\nu\rho\sigma} ,$$

where $\epsilon^{\mu\nu\rho\sigma}$ is the totally antisymmetric tensor with $\epsilon^{0123} = 1$.

Useful variants:

$$\text{Tr}(\not{p} \gamma^\mu \not{k} \gamma^\nu) = 4 [p^\mu k^\nu + p^\nu k^\mu - (p \cdot k) g^{\mu\nu}] ,$$

$$\text{Tr}(\not{a} \not{b} \not{c} \not{d}) = 4 [(a \cdot b)(c \cdot d) + (a \cdot d)(b \cdot c) - (a \cdot c)(b \cdot d)]$$

Spin averaged $|\mathcal{M}|^2$ for $e^+e^- \rightarrow \mu^+\mu^-$

Square the amplitude:

$$|\mathcal{M}|^2 = \frac{e^4}{s^2} [\bar{v}_2 \gamma_\mu u_1] [\bar{v}_2 \gamma_\nu u_1]^* [\bar{u}_3 \gamma^\mu v_4] [\bar{u}_3 \gamma^\nu v_4]^*$$

Use Casimir's trick (neglect rest masses)

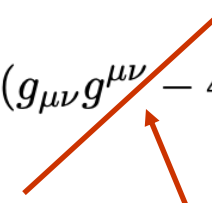
$$\sum_{s_1, s_2} [\bar{u}_1 \Gamma_A u_2] [\bar{u}_1 \Gamma_B u_2]^* \approx \text{Tr}[\not{p}_1 \Gamma_A \not{p}_2 \bar{\Gamma}_B]$$

to average over initial-state, sum over final-state spins:

$$\langle |\mathcal{M}|^2 \rangle = \frac{1}{4} \sum_{s_1, s_2, s_3, s_4} |\mathcal{M}|^2 = \frac{e^4}{4s^2} \text{Tr}[\not{p}_2 \gamma_\mu \not{p}_1 \gamma_\nu] \text{Tr}[\not{p}_3 \gamma^\mu \not{p}_4 \gamma^\nu]$$

Spin averaged $|\mathcal{M}|^2$

Use trace theorems to evaluate:

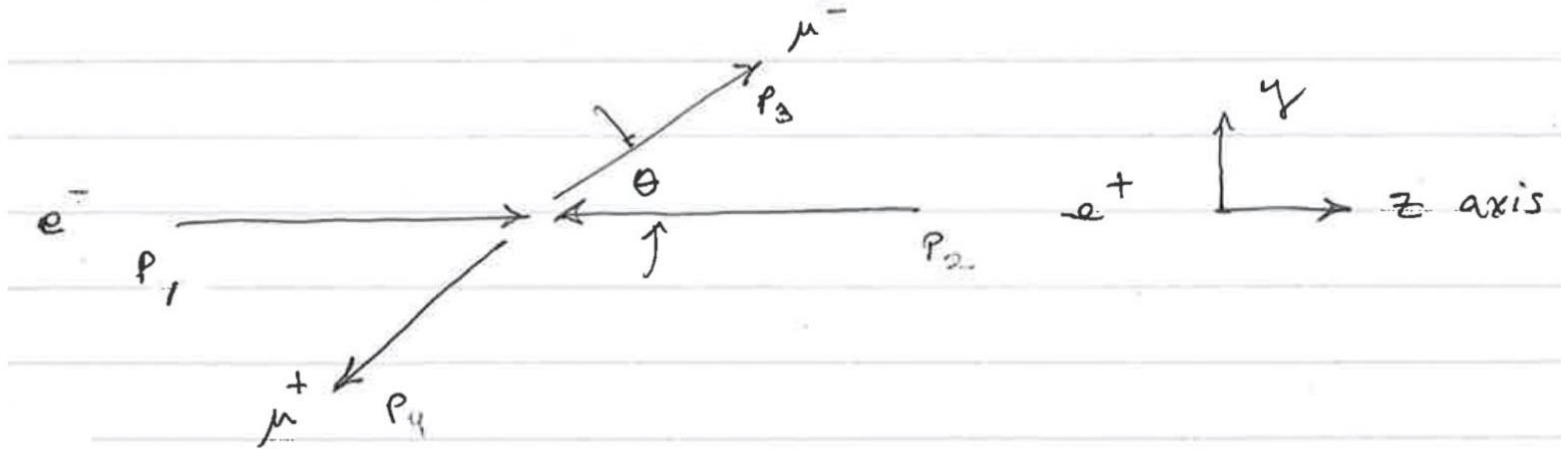
$$\begin{aligned}\langle |\mathcal{M}|^2 \rangle &= \frac{4e^4}{s^2} [p_{1\mu} p_{2\nu} + p_{2\mu} p_{1\nu} - g_{\mu\nu} (p_1 \cdot p_2)] [p_3^\mu p_4^\nu + p_4^\mu p_3^\nu - g^{\mu\nu} (p_3 \cdot p_4)] \\ &= \frac{4e^4}{s^2} [2(p_1 \cdot p_3)(p_2 \cdot p_4) + 2(p_1 \cdot p_4)(p_2 \cdot p_3) + (g_{\mu\nu} g^{\mu\nu} - 4)(p_1 \cdot p_2)(p_3 \cdot p_4)]\end{aligned}$$

$$g_{\mu\nu} g^{\mu\nu} - 4 = 0$$

Use $\alpha = e^2/4\pi = 1/137$,

$$\langle |\mathcal{M}|^2 \rangle = \frac{128\pi^2 \alpha^2}{s^2} [(p_1 \cdot p_3)(p_2 \cdot p_4) + (p_1 \cdot p_4)(p_2 \cdot p_3)]$$

Go to c.m. frame

Consider the c.m. frame:



Let $E = E_{\text{beam}} = E_{\text{cm}}/2$:

$$p_1 = (E, 0, 0, E)$$

$$p_2 = (E, 0, 0, -E)$$

$$p_3 = (E, 0, E \sin \theta, E \cos \theta) \quad (\text{suppose } \mu^+\mu^- \text{ in } (y,z) \text{ plane})$$

$$p_4 = (E, 0, -E \sin \theta, -E \cos \theta)$$

Differential cross section for $e^+e^- \rightarrow \mu^+\mu^-$

The required four-vector products are:

$$p_1 \cdot p_3 = p_2 \cdot p_4 = E^2(1 + \cos \theta) ,$$

$$p_1 \cdot p_4 = p_2 \cdot p_3 = E^2(1 - \cos \theta) ,$$

$$p_1 \cdot p_2 = p_3 \cdot p_4 = 2E^2 .$$

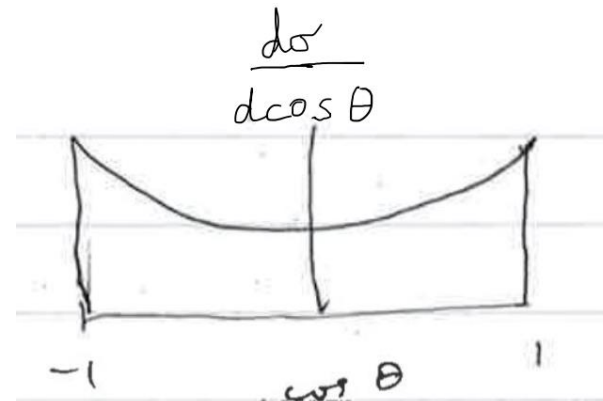
giving $\langle |\mathcal{M}|^2 \rangle = 16\pi^2\alpha^2(1 + \cos^2 \theta)$

Use in the equation for the differential cross section:

$$\frac{d\sigma}{d\Omega} = \frac{\langle |\mathcal{M}|^2 \rangle}{64\pi^2 E_{\text{cm}}^2} = \frac{\alpha^2}{4E_{\text{cm}}^2}(1 + \cos^2 \theta)$$

Integrate over φ :

$$\frac{d\sigma}{d\cos \theta} = \frac{\alpha^2\pi}{2E_{\text{cm}}^2}(1 + \cos^2 \theta)$$



Total cross section for $e^+e^- \rightarrow \mu^+\mu^-$

Integrate differential cross section over $\cos \theta$:

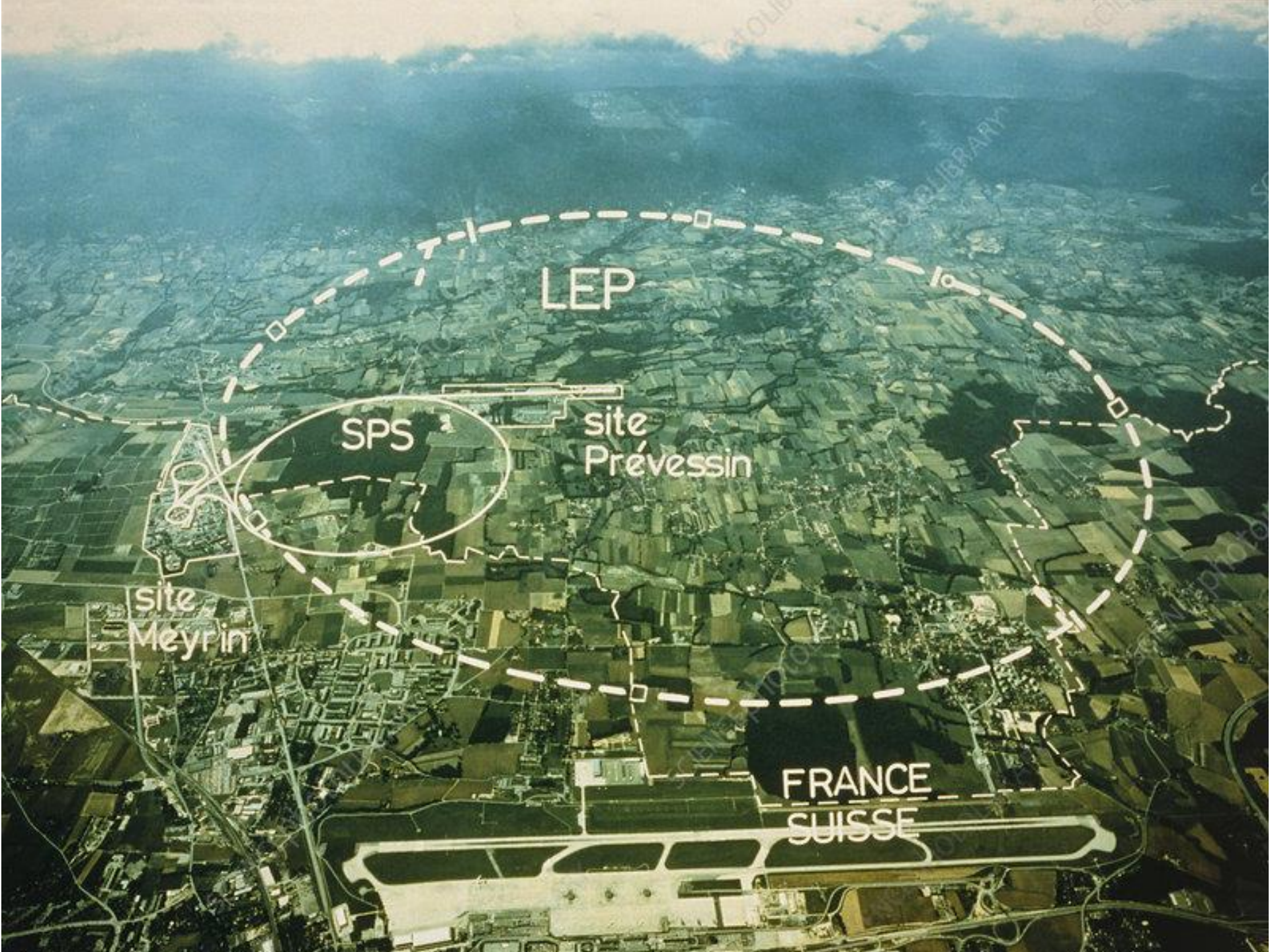
$$\sigma = \int_{-1}^1 \frac{d\sigma}{d\cos\theta} d\cos\theta = \frac{\pi\alpha^2}{2E_{\text{cm}}^2} \int_{-1}^1 (1+x^2) dx = \frac{4\pi\alpha^2}{3E_{\text{cm}}^2}$$

In experimentalists units (energy⁻² \rightarrow area): $\sigma = \frac{4\pi\alpha^2\hbar^2c^2}{3E_{\text{cm}}^2}$

TPC/2 γ Experiment at SLAC (1984-86) collected data at $E_{\text{cm}} = 29$ GeV with integrated luminosity $\int \mathcal{L} dt = 60 \text{ pb}^{-1}$

$$\sigma = 0.103 \text{ nb}$$

$$N = \sigma \int \mathcal{L} dt = 0.103 \text{ nb} \times 60 \text{ pb}^{-1} \times \frac{1000 \text{ nb}^{-1}}{1 \text{ pb}^{-1}} = 6180 \text{ events}$$



LEP

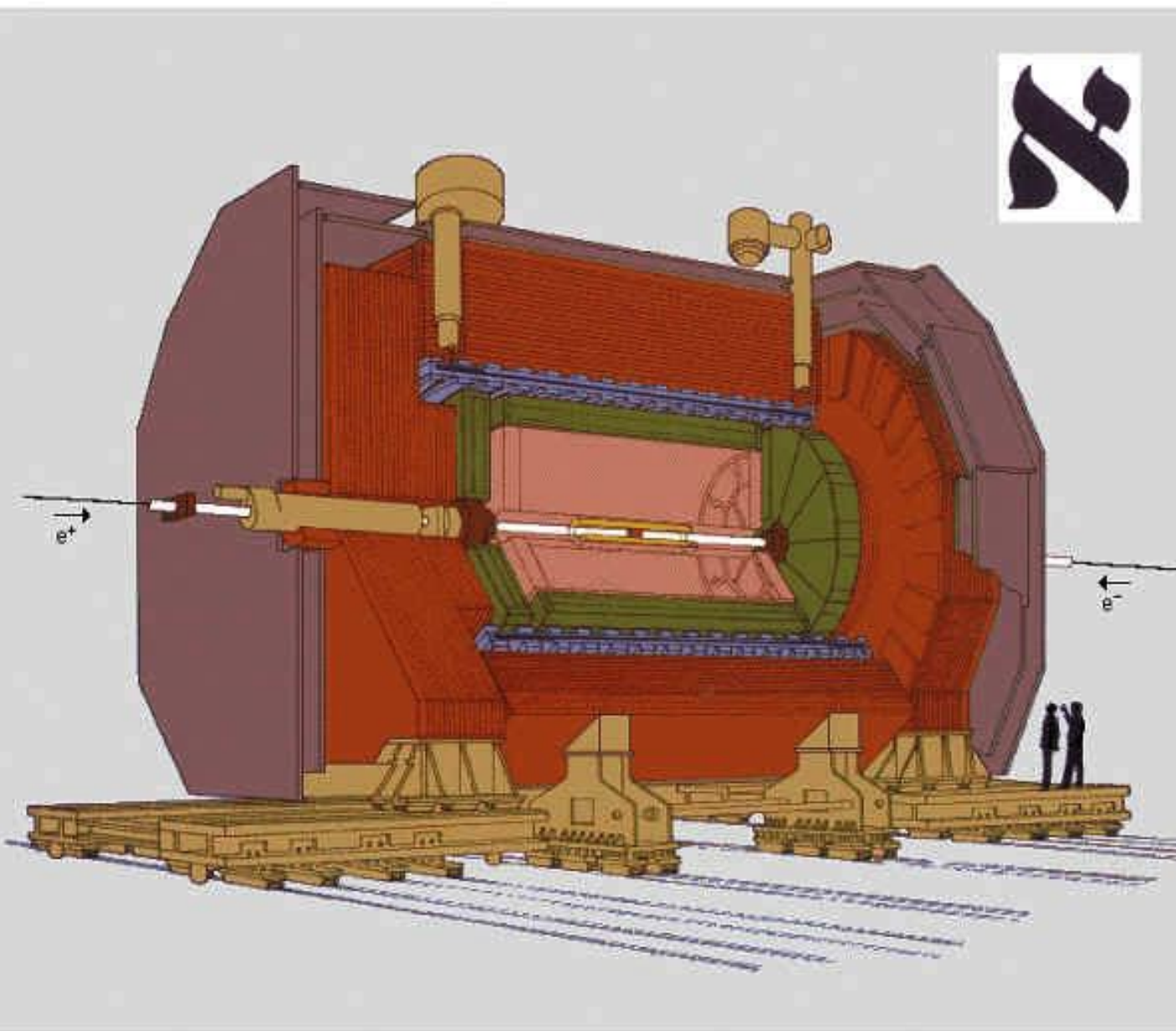
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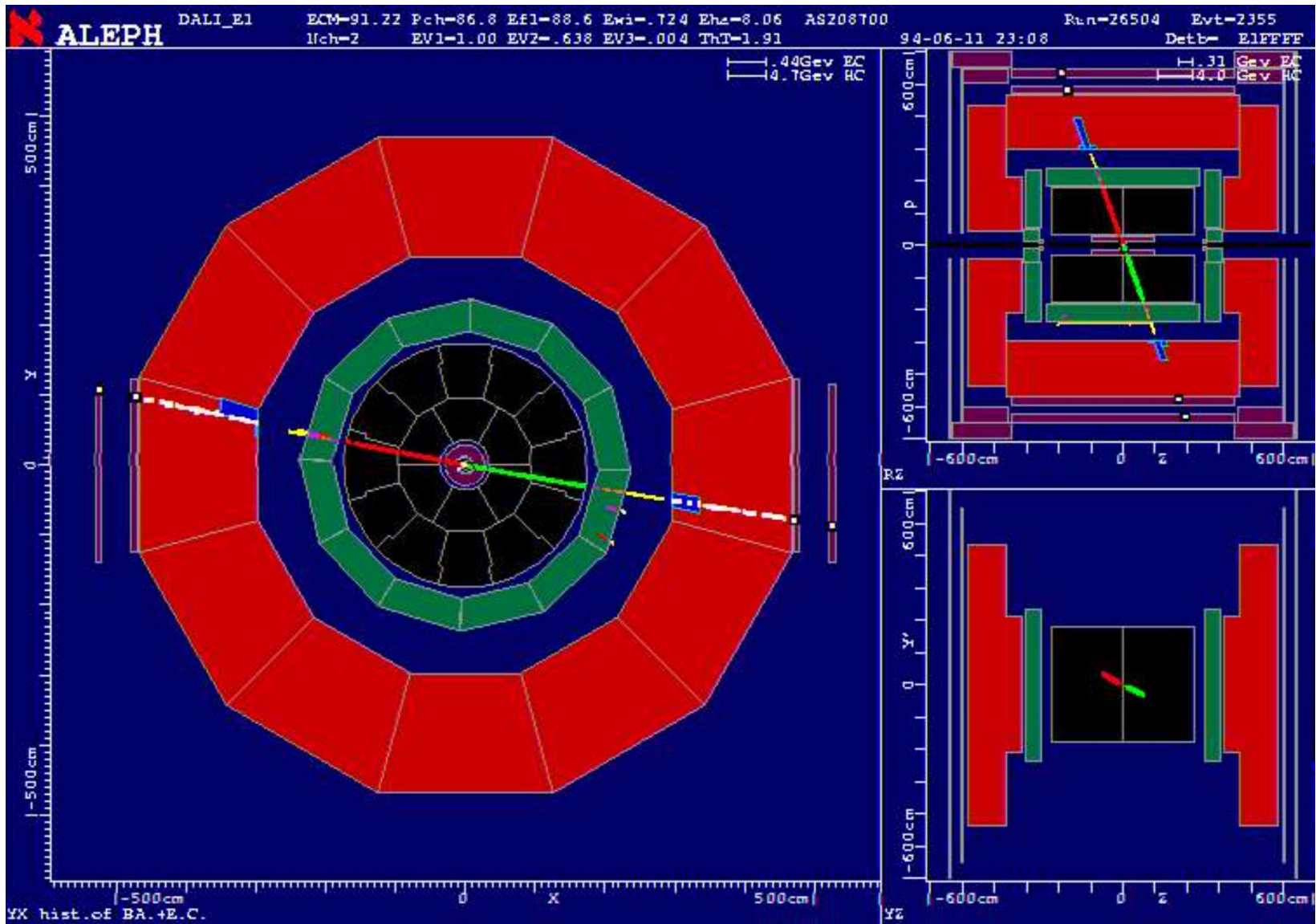
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The ALEPH detector at LEP



- Vertex Detector
- Inner Tracking Chamber
- Time Projection Chamber
- Electromagnetic Calorimeter
- Superconducting Magnet Coil
- Hadron Calorimeter
- Muon Chambers
- Luminosity Monitors

$e^+e^- \rightarrow \mu^+\mu^-$ in the ALEPH detector at LEP



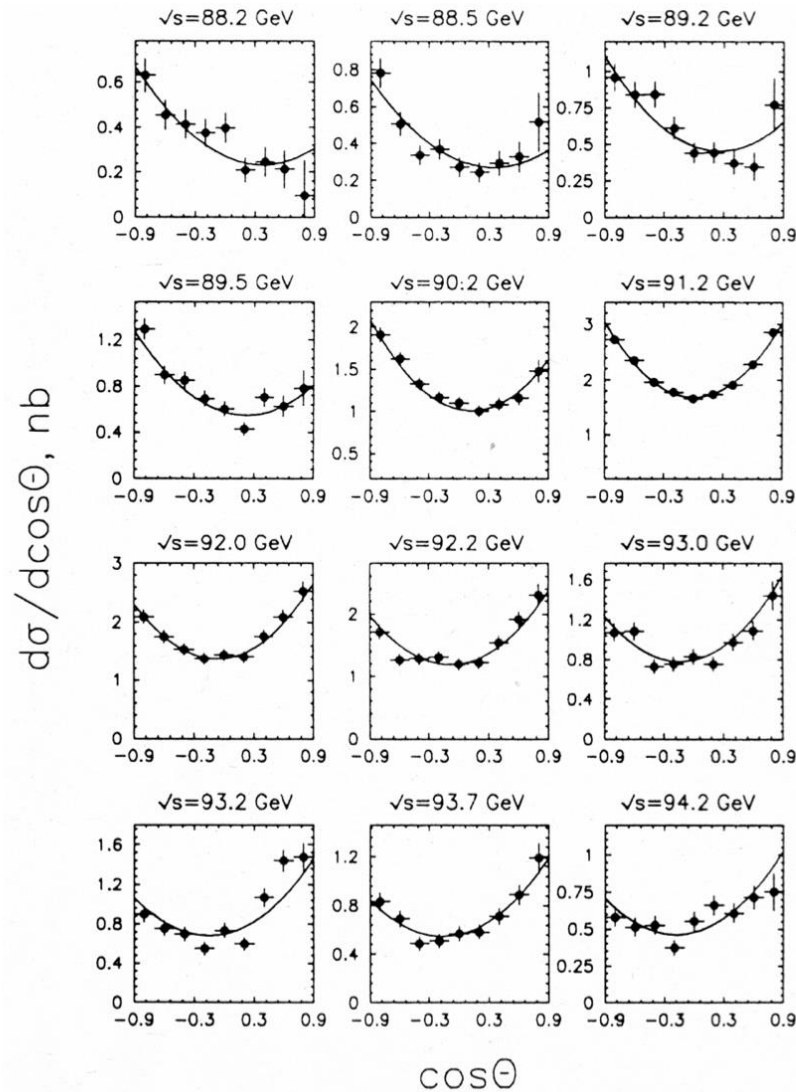


Figure 9.14: Differential cross sections $d\sigma/d\cos\theta$ for $e^+e^- \rightarrow l^+l^-$ for all charged leptons combined (from [80]).

$e^+e^- \rightarrow \mu^+\mu^-$ in the JADE detector at PETRA

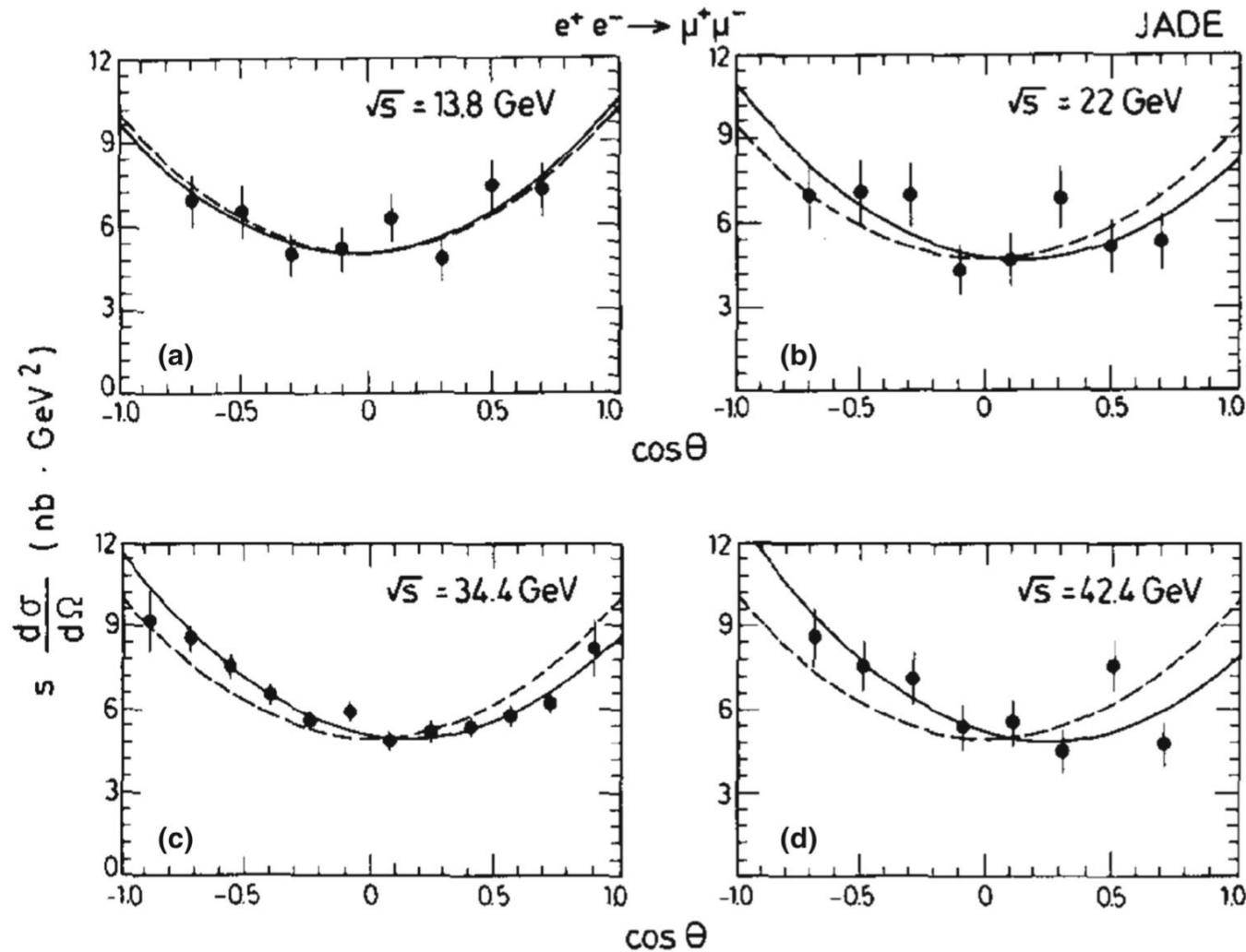


Fig. 19 Angular distributions of $e^+e^- \rightarrow \mu^+\mu^-$ for four c.m. energies. The dashed lines are symmetric fits $f(\theta) \propto (1 + \cos^2 \theta)$, the full lines are fits allowing an additional asymmetry $f(\theta) \propto (1 + \cos^2 \theta + B \cos \theta)$ [95]